## Numerical Linear Algebra Lecture 10



# We will construct a lower triangular matrix using Given's rotation from the matrix

$$A = \begin{bmatrix} 5 & 4 & 3 \\ 4 & 6 & 1 \\ 3 & 1 & 7 \end{bmatrix}.$$

Larisa Beilina, http://www.math.chalmers.se/~larisa/

# Given's matrix for j < k

function [G] = GivensMatrixLow(A, j,k)

$$a = A(k, k)$$
  

$$b = A(j, k)$$
  

$$r = sqrt(a^{2} + b^{2});$$
  

$$c = a/r;$$
  

$$s = -b/r;$$
  

$$G = eye(length(A));$$
  

$$G(j, j) = c;$$
  

$$G(k, k) = c;$$
  

$$G(j, k) = s;$$
  

$$G(k, j) = -s;$$

$$>>$$
G1up = GivensMatrixLow(A,2,3)

	1.0000000000000000000000000000000000000	0	0 ]
G1 =	0	0.989949493661166	-0.141421356237310
	0	0.141421356237310	0.989949493661166

>> A1 =G1\*A

 $A1 = \begin{bmatrix} 5.0000000000000 & 4.0000000000 & 3.00000000000 \\ 3.535533905932737 & 5.798275605729690 & -0.0000000000000 \\ 3.535533905932738 & 1.838477631085023 & 7.071067811865475 \end{bmatrix}$ 

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$$>>$$
G2 = GivensMatrixLow(A1,1,3)

 $G2 = \begin{bmatrix} 0.920574617898323 & 0 & -0.390566732942472 \\ 0 & 1.00000000000 & 0 \\ 0.390566732942472 & 0 & 0.920574617898323 \end{bmatrix}$ 

>> A2 = G2\*A1

 $A2 = \begin{bmatrix} 3.222011162644131 & 2.964250269632601 & -0.0000000000000000\\ 3.535533905932737 & 5.798275605729690 & -0.00000000000000\\ 5.207556439232954 & 3.254722774520597 & 7.681145747868607 \end{bmatrix}$ 

>>G3 = GivensMatrixLow(A2,1,2)

>> A3=G3\*A2

 $A3 = \begin{bmatrix} 1.259496302198541 & 0 & -0.00000000000000\\ 4.614653291088246 & 6.512048806713364 & -0.00000000000000\\ 5.207556439232954 & 3.254722774520597 & 7.681145747868607 \end{bmatrix}$ 



# We will construct an upper triangular matrix using Given's rotation from the matrix

$$A = \begin{bmatrix} 5 & 4 & 3 \\ 4 & 6 & 1 \\ 3 & 1 & 7 \end{bmatrix}.$$

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```
» G1=GivensMatrixUpper(A,2,1)
G1 =
0.7809 0.6247 0
-0.6247 0.7809 0
0 0 1.0000
» A1=G1*A
A1 =
6.4031 6.8716 2.9673
0 2.1864 -1.0932
3,0000 1,0000 7,0000
» G2=GivensMatrixUpper(A1,3,1)
G_{2} =
0.9055 0 0.4243
0 1.0000 0
-0.4243 0 0.9055
```

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» A2=G2\*A1 A2 =7.0711 6.6468 5.6569 0 2.1864 -1.0932 0.0000 -2.0099 5.0799 » G3=GivensMatrixUpper(A2,3,2) G3 =1,0000 0 0 0 0.7362 -0.6768 0 0.6768 0.7362 » A3=G3\*A2 Δ3 = 7.0711 6.6468 5.6569 -0.0000 2.9698 -4.24260.0000 0.0000 3.0000

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Transform the given matrix A to a tridiagonal by a single transformation based on a Given's rotation.

$$\mathbf{A} = \begin{bmatrix} 7 & 4 & 3 \\ 4 & 5 & 2 \\ 3 & 2 & 2 \end{bmatrix}$$

To obtain the tridiagonal matrix using Given's rotation we have to zero out (1,3), (3,1) elements of the matrix A.

To do that we compute elements of Givens matrix G and then compute  $GAG^{T}$ . Values of c, s are computed from the known a = 4 and b = 3 as

$$\begin{bmatrix} c & -s \\ s & c \end{bmatrix} \cdot \begin{bmatrix} a \\ b \end{bmatrix} = \begin{bmatrix} r \\ 0 \end{bmatrix}$$

We get formulas:

$$r = \sqrt{a^2 + b^2} = \sqrt{4^2 + 3^2} = 5,$$
  

$$c = \frac{a}{r} = 4/5 = 0.8,$$
  

$$s = \frac{-b}{r} = -3/5 = -0.6.$$

The Given's matrix will be

$$\mathbf{G} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c & -s \\ 0 & s & c \end{bmatrix}$$
$$\mathbf{G} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0.8 & 0.6 \\ 0 & -0.6 & 0.8 \end{bmatrix}$$

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or

Then the tridiagonal matrix will be computed as

$$\mathbf{GAG}^{\mathsf{T}} = \begin{bmatrix} 7 & 5 & 0 \\ 5 & 5.84 & -0.88 \\ 0 & -0.88 & 1.16 \end{bmatrix}$$

Another way to construct tridiagonal matrix is to compute the second Givens matrix  $G_2$  and apply it for  $A_1 = GA$  where

$$\mathbf{A_1} = \mathbf{G}\mathbf{A} = \begin{bmatrix} 7 & 4 & 3 \\ 5 & 5.2 & 2.8 \\ 0 & -1.4 & 0.4 \end{bmatrix}$$

Values of c, s are computed from the known a = 2.8 and b = 3 as

$$r = \sqrt{a^2 + b^2} = \sqrt{2.8^2 + 3^2} \approx 4.1037$$
  

$$c = \frac{a}{r} = 4/5 = 0.6823,$$
  

$$s = \frac{-b}{r} = -3/5 = -0.7311.$$

The second Given's matrix will be

$$\mathbf{G_2} = \begin{bmatrix} c & s & 0 \\ -s & c & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

The tridiagonal matrix is computed as:

$$\mathbf{A_2} = \mathbf{G_2}\mathbf{A_1} = \mathbf{G_2}\mathbf{G}\mathbf{A} = \begin{bmatrix} 1.121 & -1.0722 & 0\\ 8.529 & 6.4723 & 4.1037\\ 0 & -1.4 & 0.4 \end{bmatrix}$$

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### We will construct tridiagonal matrix from the matrix

$$A = \begin{bmatrix} 5 & 4 & 3 \\ 3 & 6 & 1 \\ 4 & 1 & 7 \end{bmatrix}$$

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using Hauseholder transformations.

To perform tridiagonal reduction for the matrix A we use Hauseholder transformation in following steps:

• Choose 
$$x = (3, 4)^T$$
 and compute

 $u = x + \alpha e_1$ , where  $\alpha = -sign(3) \cdot ||x||, ||x|| = \sqrt{25} = 5$ , and thus  $\alpha = -5$ .

- Construct  $u = x + \alpha e_1 = (3,4)^T (5,0)^T = (-2,4)^T$ .
- Construct

$$v = \frac{u}{\|u\|}$$

with  $||u|| = \sqrt{20}$ . Therefore  $v = (-2/\sqrt{(20)}, 4/\sqrt{(20)})^T$ .

Compute

$$Q' = I - 2vv^{T} = \begin{pmatrix} 0.6 & 0.8 \\ 0.8 & -0.6 \end{pmatrix}.$$

• Construct the matrix of the Householder transformation as:

$$Q_1 = egin{pmatrix} 1 & 0 & 0 \ 0 & 0.6 & 0.8 \ 0 & 0.8 & -0.6 \end{pmatrix}$$

• Then compute

$$A_1 = Q_1 A = egin{pmatrix} 5 & 4 & 3 \ 5 & 4.4 & 6.2 \ 0 & 4.2 & -3.4 \end{pmatrix}.$$

such that  $Q_1$  leaves the first row of  $Q_1A$  unchanged.

• Choose new vector  $x = (4,3)^T$  for  $A_1^T$  and compute

 $u = x + \alpha e_1$ ,

where  $\alpha = -sign(4) \cdot ||x||, ||x|| = \sqrt{25} = 5$ , and thus  $\alpha = -5$ .

• Construct 
$$u = x + \alpha e_1 = (4,3)^T - (5,0)^T = (-1,3)^T$$
.

• Construct 
$$v = \frac{u}{\|u\|}$$

with 
$$||u|| = \sqrt{10}$$
.  
Therefore  $v = (-1/\sqrt{10}, 3/\sqrt{10})^{T}$ .

Compute

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$$V' = I - 2vv^{T} = \begin{pmatrix} 0.8 & 0.6 \\ 0.6 & -0.8 \end{pmatrix}.$$

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• Construct the second matrix of the Householder transformation  $V_1$  as:

$$V_1 = \left\lfloor \begin{array}{c|c} 1 & 0 \\ \hline 0 & V' \end{array} \right\rfloor$$

to get

$$V_1 = egin{pmatrix} 1 & 0 & 0 \ 0 & 0.8 & 0.6 \ 0 & 0.6 & -0.8 \end{pmatrix}$$

and then compute

$$Q_1 A V_1 = \begin{pmatrix} 5 & 5 & 0 \\ 5 & 7.24 & -2.32 \\ 0 & 1.32 & 5.24 \end{pmatrix}.$$

such that  $V_1$  leaves the first column of  $A_1$  unchanged.

# **Canonical Forms**

DEFINITION. The polynomial  $p(\lambda) = det(A - \lambda I)$  is called the *characteristic polynomial of A*. The roots of  $p(\lambda) = 0$  are the eigenvalues of A.

Since the degree of the characteristic polynomial  $p(\lambda)$  equals *n*, the dimension of *A*, it has *n* roots, so *A* has *n* eigenvalues.

DEFINITION. A nonzero vector x satisfying  $Ax = \lambda x$  is a *(right)* eigenvector for the eigenvalue  $\lambda$ . A nonzero vector y such that  $y^*A = \lambda y^*$  is a *left eigenvector*. (Recall that  $y^* = (\bar{y})^T$  is the conjugate transpose of y.) DEFINITION. Let S be any nonsingular matrix. Then A and  $B = S^{-1}AS$  are called *similar* matrices, and S is a similarity transformation. PROPOSITION. Let  $B = S^{-1}AS$ , so A and B are similar. Then A and B have the same eigenvalues, and x (or y) is a right (or left) eigenvector of A if and only if  $S^{-1}x$  (or  $S^*y$ ) is a right (or left) eigenvector of B.

*Proof.* Using the fact that  $det(X \cdot Y) = det(X) \cdot det(Y)$  for any square matrices X and Y, we can write

$$\det(A - \lambda I) = \det(S^{-1}(A - \lambda I)S) = \det(B - \lambda I).$$

So A and B have the same characteristic polynomials.  $Ax = \lambda x$  holds if and only if  $\underbrace{S^{-1}AS}_{B} \underbrace{S^{-1}x}_{x^*} = \lambda \underbrace{S^{-1}x}_{x^*}$  or  $B(S^{-1}x) = \lambda(S^{-1}x)$ . Similarly,  $y^*A = \lambda y^*$  if and only if  $y^*SS^{-1}AS = \lambda y^*S$  or  $(S^*y)^*B = \lambda(S^*y)^*$ .  $\Box$  THEOREM. Jordan canonical form. Given A, there exists a nonsingular S such that  $S^{-1}AS = J$ , where J is in Jordan canonical form. This means that J is block diagonal, with  $J = diag(J_{n_1}(\lambda_1), J_{n_2}(\lambda_2), \dots, J_{n_k}(\lambda_k))$  and

$$J_{n_i}(\lambda_i) = \left[egin{array}{cccc} \lambda_i & 1 & & 0 \ & \ddots & \ddots & \ & & \ddots & 1 \ 0 & & & \lambda_i \end{array}
ight]^{n_i imes n_i}$$

J is unique, up to permutations of its diagonal blocks.

For a proof of this theorem, see a book on linear algebra such as [F. Gantmacher. The Theory of Matrices, vol. II (translation). Chelsea, New York, 1959] or [P. Halmos. Finite Dimensional Vector Spaces. Van Nostrand, New York, 1958].

- Each J<sub>m</sub>(λ) is called a Jordan block with eigenvalue λ of algebraic multiplicity m.
- If some n<sub>i</sub> = 1, and λ<sub>i</sub> is an eigenvalue of only that one Jordan block, then λ<sub>i</sub> is called a *simple eigenvalue*.
- If all n<sub>i</sub> = 1, so that J is diagonal, A is called *diagonalizable*; otherwise it is called *defective*.
- An n-by-n defective matrix does not have *n* eigenvectors. Although defective matrices are "rare" in a certain well-defined sense, the fact that some matrices do not have *n* eigenvectors is a fundamental fact confronting anyone designing algorithms to compute eigenvectors and eigenvalues.
- Symmetric matrices are never defective.

### PROPOSITION.

- A Jordan block has one right eigenvector, e<sub>1</sub> = [1,0,...,0]<sup>T</sup>, and one left eigenvector, e<sub>n</sub> = [0,...,0,1]<sup>T</sup>.
- Therefore, a matrix has *n* eigenvectors matching its *n* eigenvalues if and only if it is diagonalizable.
- In this case,  $S^{-1}AS = diag(\lambda_i)$ . This is equivalent to  $AS = S diag(\lambda_i)$ , so the i-th column of S is a right eigenvector for  $\lambda_i$ .
- It is also equivalent to  $S^{-1}A = diag(\lambda_i)S^{-1}$ , so the conjugate transpose of the ith row of  $S^{-1}$  is a left eigenvector for  $\lambda_i$ .
- If all *n* eigenvalues of a matrix *A* are distinct, then *A* is diagonalizable.

*Proof.* Let  $J = J_m(\lambda)$  for ease of notation. It is easy to see  $Je_1 = \lambda e_1$ and  $e_n^T J = \lambda e_n^T$ , so  $e_1$  and  $e_n$  are right and left eigenvectors of J, respectively. To see that J has only one right eigenvector (up to scalar multiples), note that any eigenvector x must satisfy  $(J - \lambda I)x = 0$ , so xis in the null space of

$$J - \lambda I = \begin{bmatrix} 0 & 1 & & \\ & \ddots & \ddots & \\ & & \ddots & 1 \\ & & & 0 \end{bmatrix}$$

But the null space of  $J - \lambda I$  is clearly span $(e_1)$ , so there is just one eigenvector. If all eigenvalues of A are distinct, then all its Jordan blocks must be 1-by-1, so  $J = diag(\lambda_1, \ldots, \lambda_n)$  is diagonal.  $\Box$ 

We illustrate the concepts of eigenvalue and eigenvector with a problem of *mechanical vibrations*. We will see a defective matrix arise in a natural physical context.

Newton's law F = ma applied to this system yields

$$egin{aligned} m_i \ddot{x}_i(t) &= k_i (x_{i-1}(t) - x_i(t)) \ & ext{force on mass from spring } i \ &+ k_{i+1} (x_{i+1}(t) - x_i(t)) \ & ext{force on mass from spring } i+1 \ &- b_i \dot{x}_i(t) \ & ext{force on mass from damper } i \end{aligned}$$

or

$$M\ddot{x}(t) = -B\dot{x}(t) - Kx(t),$$

where  $M = diag(m_1, \ldots, m_n)$ ,  $B = diag(b_1, \ldots, b_n)$ , and

$$\mathcal{K} = \begin{bmatrix}
 k_1 + k_2 & -k_2 & & \\
 -k_2 & k_2 + k_3 & -k_3 & & \\
 & \ddots & \ddots & \ddots & \\
 & -k_{n-1} & k_{n-1} + k_n & -k_n \\
 & & -k_n & k_n
 \end{bmatrix}$$

We assume that all the masses  $m_i$  are positive. M is called the mass matrix, B is the damping matrix, and K is the stiffness matrix.

Electrical engineers analyzing linear circuits arrive at an analogous equation by applying Kirchoff's and related laws instead of Newton's law. In this case x represents branch currents, M represent inductances, B represents resistances, and K represents admittances (reciprocal capacitances).

We will use a standard trick to change this second-order differential equation to a first-order differential equation, changing variables to

$$y(t) = \left[ egin{array}{c} \dot{x}(t) \ x(t) \end{array} 
ight]$$

## This yields

$$\dot{y}(t) = \begin{bmatrix} \ddot{x}(t) \\ \dot{x}(t) \end{bmatrix} = \begin{bmatrix} -M^{-1}B\dot{x}(t) - M^{-1}Kx(t) \\ \dot{x}(t) \end{bmatrix}$$
$$= \begin{bmatrix} -M^{-1}B & -M^{-1}K \\ I & 0 \end{bmatrix} \cdot \begin{bmatrix} \dot{x}(t) \\ x(t) \end{bmatrix}$$
$$= \begin{bmatrix} -M^{-1}B & -M^{-1}K \\ I & 0 \end{bmatrix} \cdot y(t) \equiv Ay(t).$$

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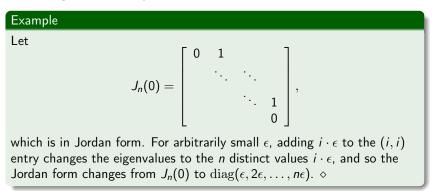
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- To solve y(t) = Ay(t), we assume that y(0) is given (i.e., the initial positions x(0) and velocities x(0) are given).
- One way to write down the solution of this differential equation is  $y(t) = e^{At}y(0)$ , where  $e^{At}$  is the matrix exponential. We will give another more elementary solution in the special case where A is diagonalizable; this will be true for almost all choices of  $m_i$ ,  $k_i$ , and  $b_i$ .
- When A is diagonalizable, we can write  $A = S\Lambda S^{-1}$ , where  $\Lambda = \text{diag}(\lambda_1, \dots, \lambda_n)$ .
- Then  $\dot{y}(t) = Ay(t)$  is equivalent to  $\dot{y}(t) = S\Lambda S^{-1}y(t)$  or  $S^{-1}\dot{y}(t) = \Lambda S^{-1}y(t)$  or  $\dot{z}(t) = \Lambda z(t)$ , where  $z(t) \equiv S^{-1}y(t)$ .
- This diagonal system of differential equations  $\dot{z}_i(t) = \lambda_i z_i(t)$  has solutions  $z_i(t) = e^{\lambda_i t} z_i(0)$ . Since  $z(t) \equiv S^{-1} y(t)$  and  $z(0) \equiv S^{-1} y(0)$  then  $Sz(t) \equiv y(t)$  and so  $y(t) = S \operatorname{diag}(e^{\lambda_1 t}, \dots, e^{\lambda_n t}) S^{-1} y(0) = S e^{\Lambda t} S^{-1} y(0)$ .

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# Disadvantages of Jordan form

First reason: It is a discontinuous function of A, so any rounding error can change it completely.



# Disadvantages of Jordan form

Second reason: It cannot be computed stably in general. In other words, when we have finished computing S and J, we cannot guarantee that  $S^{-1}(A + \delta A)S = J$  for some small  $\delta A$ .

#### Example

Suppose  $S^{-1}AS = J$  exactly, where S is very ill-conditioned. ( $\kappa(S) = ||S|| \cdot ||S^{-1}||$  is very large.) Suppose that we are extremely lucky and manage to compute S exactly and J with just a tiny error  $\delta J$  with  $||\delta J|| = O(\varepsilon)||A||$ . How big is the backward error? In other words, how big must  $\delta A$  be so that  $S^{-1}(A + \delta A)S = J + \delta J$ ? We get

$$\underbrace{S^{-1}AS}_{J} + \underbrace{S^{-1}\delta AS}_{\delta J} = J + \delta J \to \delta A = S\delta JS^{-1},$$

and all that we can conclude is that  $||\delta A|| \leq ||S|| \cdot ||\delta J|| \cdot ||S^{-1}|| = \kappa(S) \cdot ||\delta J|| = O(\varepsilon)\kappa(S)||A||$ . Thus  $||\delta A||$ may be much larger than  $\varepsilon ||A||$  because of  $\kappa(S)$ , which prevents backward stability.  $\diamond$ 

# From Jordan to Schur canonical form

- Instead of computing  $S^{-1}AS = J$ , where S can be an arbitrarily ill-conditioned matrix, we will restrict S to be orthogonal (so  $\kappa_2(S) = 1$ ) to guarantee stability.
- We cannot get a canonical form as simple as the Jordan form any more.

THEOREM. Schur canonical form. Given A, there exists a unitary matrix Q ( $Q^*Q = QQ^* = I$ ) and an upper triangular matrix T such that  $Q^*AQ = T$ . The eigenvalues of A are the diagonal entries of T.

*Proof.* We use induction on *n*. It is obviously true if *A* is 1 by 1. Now let  $\lambda$  be any eigenvalue and *u* a corresponding eigenvector normalized so  $||u||_2 = 1$ . Choose  $\tilde{U}$  so  $U = [u, \tilde{U}]$  is a square unitary matrix. (Note that  $\lambda$  and *u* may be complex even if *A* is real.) Then

$$U^* \cdot A \cdot U = \begin{bmatrix} u^* \\ \tilde{U}^* \end{bmatrix} \cdot A \cdot [u, \tilde{U}] = \begin{bmatrix} u^* A u & u^* A \tilde{U} \\ \tilde{U}^* A u & \tilde{U}^* A \tilde{U} \end{bmatrix}.$$

Now we can write 
$$u^*Au = \lambda u^*u = \lambda$$
, and  $\tilde{U}^*Au = \lambda \tilde{U}^*u = 0$  so  
 $U^*AU \equiv \begin{bmatrix} \lambda & \tilde{a}_{12} \\ 0 & \tilde{A}_{22} \end{bmatrix}$ . By induction, there is a unitary  $P$ , so  
 $P^*\tilde{A}_{22}P = \tilde{T}$  is upper triangular. Then

$$U^*AU = \begin{bmatrix} \lambda & \tilde{a}_{12} \\ 0 & P\tilde{T}P^* \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & P \end{bmatrix} \begin{bmatrix} \lambda & \tilde{a}_{12}P \\ 0 & \tilde{T} \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & P^* \end{bmatrix},$$

SO

$$\underbrace{\begin{bmatrix} 1 & 0 \\ 0 & P^* \end{bmatrix} U^*}_{Q^*} A \underbrace{U \begin{bmatrix} 1 & 0 \\ 0 & P \end{bmatrix}}_{Q} = \begin{bmatrix} \lambda & \tilde{a}_{12}P \\ 0 & \tilde{T} \end{bmatrix} = T$$

is upper triangular and  $Q = U \begin{bmatrix} 1 & 0 \\ 0 & P \end{bmatrix}$  is unitary as desired.  $\diamond$ 

# Remarks on Schur canonical form

- Notice that the Schur form is not unique, because the eigenvalues may appear on the diagonal of *T* in any order.
- This introduces complex numbers even when A is real. When A is real, we prefer a canonical form that uses only real numbers, because it will be cheaper to compute.
- This means that we will have to sacrifice a triangular canonical form and settle for a block-triangular canonical form.

# Real Schur canonical form

THEOREM. *Real Schur canonical form.* If A is real, there exists a real orthogonal matrix V such that  $V^T AV = T$  is quasi-upper triangular. This means that T is block upper triangular with 1-by-1 and 2-by-2 blocks on the diagonal. Its eigenvalues are the eigenvalues of its diagonal blocks. The 1-by-1 blocks correspond to real eigenvalues, and the 2-by-2 blocks to complex conjugate pairs of eigenvalues.

## Computing Eigenvectors from the Schur Form

Let  $Q^*AQ = T$  be the Schur form. Then if  $Tx = \lambda x$ , we have

$$Tx = \lambda x \rightarrow (Q^*AQ)x = Tx = \lambda x \rightarrow AQx = QTx = \lambda Qx$$

, so Qx is an eigenvector of A. So to find eigenvectors Qx of A, it suffices to find eigenvectors x of T.

Suppose that  $\lambda = t_{ii}$  has multiplicity 1 (i.e., it is simple). Write  $(T - \lambda I)x = 0$  as

$$0 = \begin{bmatrix} T_{11} - \lambda I & T_{12} & T_{13} \\ 0 & 0 & T_{23} \\ 0 & 0 & T_{33} - \lambda I \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix}$$
$$= \begin{bmatrix} (T_{11} - \lambda I)x_1 + T_{12}x_2 + T_{13}x_3 \\ T_{23}x_3 \\ (T_{33} - \lambda I)x_3 \end{bmatrix},$$

where  $T_{11}$  is (i - 1)-by-(i - 1),  $T_{22} = \lambda$  is 1-by-1,  $T_{33}$  is (n - i)-by-(n - i), and x is partitioned correspondingly.

Since  $\lambda$  is simple, both  $T_{11} - \lambda I$  and  $T_{33} - \lambda I$  are nonsingular, so  $(T_{33} - \lambda I)x_3 = 0$  implies  $x_3 = 0$ . Therefore  $(T_{11} - \lambda I)x_1 = -T_{12}x_2$ . Choosing (arbitrary)  $x_2 = 1$  means  $x_1 = -(T_{11} - \lambda I)^{-1}T_{12}$ , so

$$x = \begin{bmatrix} (\lambda I - T_{11})^{-1} T_{12} \\ 1 \\ 0 \end{bmatrix},$$

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# Condition number of the eigenvalue $\lambda$

THEOREM. Let  $\lambda$  be a simple eigenvalue of A with right eigenvector x and left eigenvector y, normalized so that  $||x||_2 = ||y||_2 = 1$ . Let  $\lambda + \delta \lambda$  be the corresponding eigenvalue of  $A + \delta A$ . Then

$$\begin{split} \delta \lambda &= \frac{y^* \delta A x}{y^* x} + O(||\delta A||^2) \text{ or } \\ |\delta \lambda| &\leq \frac{||\delta A||}{|y^* x|} + O(||\delta A||^2) = \sec \Theta(y, x) ||\delta A|| + O(||\delta A||^2), \end{split}$$

where  $\Theta(y, x)$  is the acute angle between y and x. In other words,  $\sec \Theta(y, x) = 1/|y^*x|$  is the condition number of the eigenvalue  $\lambda$ .

Proof. Subtract  $Ax = \lambda x$  from  $(A + \delta A)(x + \delta x) = (\lambda + \delta \lambda)(x + \delta x)$  to get

$$A\delta x + \delta A x + \underbrace{\delta A \delta x}_{0} = \lambda \delta x + \delta \lambda x + \underbrace{\delta \lambda \delta x}_{0}.$$

Ignore the second-order terms (those with two " $\delta$  terms" as factors:  $\delta A \delta x$  and  $\delta \lambda \delta x$ ) and multiply by  $y^*$  to get

$$\underbrace{y^* A \delta x}_{cancels} + y^* \delta A x = \underbrace{y^* \lambda \delta x}_{cancels} + y^* \delta \lambda x.$$

Now  $y^*A\delta x$  cancels  $y^*\lambda\delta x$ , so we can solve for  $\delta\lambda = (y^*\delta Ax)/(y^*x)$  as desired.  $\Box$ 

COROLLARY. Let A be symmetric (or normal:  $AA^* = A^*A$ ). Then  $|\delta\lambda| \le ||\delta A|| + O(||\delta A||^2)$ .

*Proof.* If A is symmetric or normal, then its eigenvectors are all orthogonal, i.e.,  $Q^*AQ = \Lambda$  with  $QQ^* = I$ . So the right eigenvectors x (columns of Q) and left eigenvectors y (conjugate transposes of the rows of  $Q^*$ ) are identical, and  $1/|y^*x| = 1$ .  $\Box$ 

## Bauer-Fike Theorem.

THEOREM. Bauer-Fike. Let A have all simple eigenvalues (i.e., be diagonalizable). Call them  $\lambda_i$ , with right and left eigenvectors  $x_i$  and  $y_i$ , normalized so  $||x_i||_2 = ||y_i||_2 = 1$ . Then the eigenvalues of  $A + \delta A$  lie in disks  $B_i$ , where  $B_i$  has center  $\lambda_i$  and radius  $n \frac{||\delta A||_2}{|y_i^* x_i|}$ . Our proof will use Gershgorin's theorem.

# GERSHGORIN'S THEOREM

GERSHGORIN'S THEOREM. Let *B* be an arbitrary matrix. Then the eigenvalues  $\lambda$  of *B* are located in the union of the *n* disks defined by  $|\lambda - b_{ii}| \leq \sum_{j \neq i} |b_{ij}|$  for i = 1 to *n*. We will also need two lemmas which we present on the next slides.

LEMMA. Let  $S = [x_1, ..., x_n]$  the nonsingular matrix of right eigenvectors. Then

$$S^{-1} = \begin{bmatrix} y_1^* / y_1^* x_1 \\ y_2^* / y_2^* x_2 \\ \vdots \\ y_n^* / y_n^* x_n \end{bmatrix}$$

Proof of Lemma. We know that A is diagonalizible and thus  $S^{-1}AS = \Lambda$ , or  $AS = S\Lambda$ , where  $\Lambda = \text{diag}(\lambda_1, \dots, \lambda_n)$ , since the columns  $x_i$  of S are eigenvectors. This is equivalent to  $S^{-1}ASS^{-1} = \Lambda S^{-1}$  or

 $S^{-1}A = \Lambda S^{-1}$ , so the rows of  $S^{-1}$  are conjugate transposes of the left eigenvectors  $y_i$ . So

$$S^{-1} = \begin{bmatrix} y_1^* \cdot c_1 \\ \vdots \\ y_n^* \cdot c_n \end{bmatrix}$$

for some constants  $c_i$ . But  $I = S^{-1}S$ , so  $1 = (S^{-1}S)_{ii} = y_i^* x_i \cdot c_i$ , and  $c_i = \frac{1}{y_i^* x_i}$  as desired.  $\Box$ 

LEMMA. If each column of (any matrix) *S* has two-norm equal to 1,  $||S||_2 \le \sqrt{n}$ . Similarly, if each row of a matrix has two-norm equal to 1, its two-norm is at most  $\sqrt{n}$ . *Proof of Lemma.*  $||S||_2 = ||S^T||_2 = \max_{||x||_2=1} ||S^Tx||_2$ . Each component of  $S^Tx$  is bounded by 1 by the Cauchy-Schwartz inequality, so  $||S^Tx||_2 \le ||[\underbrace{1,\ldots,1}_n]^T||_2 = \sqrt{n}$ .  $\Box$ 

#### Proof of the Bauer-Fike theorem.

We will apply Gershgorin's theorem to  $S^{-1}(A + \delta A)S = \Lambda + F$ , where  $\Lambda = S^{-1}AS = \operatorname{diag}(\lambda_1, \ldots, \lambda_n)$  and  $F = S^{-1}\delta AS$ . The idea is to show that the eigenvalues of  $A + \delta A$  lie in balls centered at the  $\lambda_i$  with the given radii. To do this, we take the disks containing the eigenvalues of  $\Lambda + F$  that are defined by Gershgorin's theorem,

$$|\lambda - (\lambda_i + f_{ii})| \le \sum_{j \ne i}^n |f_{ij}|,$$

and enlarge them slightly to get the disks

$$\begin{aligned} |\lambda - \lambda_i| &\leq \sum_{j}^{n} 1 \cdot |f_{ij}| \leq (\sum_{j}^{n} 1)^{1/2} \left( \sum_{j}^{n} |f_{ij}|^2 \right)^{1/2} \text{ by Cauchy - Schwarz} \\ &\leq n^{1/2} \left( \sum_{j}^{n} |f_{ij}|^2 \right)^{1/2} \text{ by Cauchy - Schwarz} \\ &= n^{1/2} \cdot ||F(i,:)||_2. \end{aligned}$$

Now we need to bound the two-norm of the *i*-th row F(i,:) of  $F = S^{-1} \delta AS$  :

$$\begin{aligned} ||F(i,:)||_2 &= ||(S^{-1}\delta AS)(i,:)||_2 \\ &\leq ||(S^{-1})(i,:)||_2 \cdot ||\delta A||_2 \cdot \underbrace{||S||_2}_{\leq \sqrt{n}} \\ &\leq \frac{n^{1/2}}{|y_i^* x_i|} \cdot ||\delta A||_2 \quad \text{by Lemmas.} \end{aligned}$$

Combined this equation with equation above, this proves the theorem.

$$\begin{aligned} |\lambda - \underbrace{\lambda_i}_{centers}| &\leq n^{1/2} \cdot ||F(i,:)||_2 \\ &\leq n^{1/2} \cdot \frac{n^{1/2}}{|y_i^* x_i|} \cdot ||\delta A||_2 = \underbrace{\frac{n}{|y_i^* x_i|} \cdot ||\delta A||_2}_{radius} \end{aligned}$$

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THEOREM. Let  $\lambda$  be a simple eigenvalue of A, with unit right and left eigenvectors x and y and condition number  $c = 1/|y^*x|$ . Then there is a  $\delta A$  such that  $A + \delta A$  has a multiple eigenvalue at  $\lambda$ , and

$$\frac{||\delta A||_2}{||A||_2} \le \frac{1}{\sqrt{c^2 - 1}}.$$

When  $c \gg 1$ , i.e., the eigenvalue is ill-conditioned, then the upper bound on the distance is  $1/\sqrt{c^2 - 1} \approx \frac{1}{c}$ , the reciprocal of the condition number.

#### We assume that A is real.

Power method

This method can find only the largest eigenvalue for A and the corresponding eigenvector.

Inverse iteration

We find all other eigenvalues and eigenvectors applying method for  $(A - \sigma I)^{-1}$  for some shift  $\sigma$ .

Orthogonal iteration

Lets compute entire invariant subspace.

QR iteration

reorganized orthogonal iteration, ultimate algorithm.

- Hessenberg reduction

### Power Method

ALGORITHM. Power method: Given  $x_0$ , we iterate

i = 0repeat  $y_{i+1} = Ax_i$   $x_{i+1} = y_{i+1}/||y_{i+1}||_2$  (approximate eigenvector)  $\tilde{\lambda}_{i+1} = x_{i+1}^T Ax_{i+1}$  (approximate eigenvalue) i = i + 1until convergence Let us first apply this algorithm in the very simple case when  $A = \operatorname{diag}(\lambda_1, ..., \lambda_n)$ , with  $|\lambda_1| > |\lambda_2| \ge \cdots \ge |\lambda_n|$ ). In this case the eigenvectors are just the columns  $e_i$  of the identity matrix. Note that  $x_i$  can also be written  $x_i = A^i x_0 / ||A^i x_0||_2$ , since the factors  $1 / ||y_{i+1}||_2$  only scale  $x_{i+1}$  to be a unit vector and do not change its direction. Taking S = I lets us write  $x_0 = S(S^{-1}x_0) = [\xi_1, \ldots, \xi_n]^T$ , or

$$\mathcal{A}^{i} x_{0} \equiv \mathcal{A}^{i} \begin{bmatrix} \xi_{1} \\ \xi_{2} \\ \vdots \\ \xi_{n} \end{bmatrix} = \begin{bmatrix} \xi_{1} \lambda_{1}^{i} \\ \xi_{2} \lambda_{2}^{i} \\ \vdots \\ \xi_{n} \lambda_{n}^{i} \end{bmatrix} = \xi_{1} \lambda_{1}^{i} \begin{bmatrix} 1 \\ \frac{\xi_{2}}{\xi_{1}} (\frac{\lambda_{2}}{\lambda_{1}})^{i} \\ \vdots \\ \frac{\xi_{n}}{\xi_{1}} (\frac{\lambda_{n}}{\lambda_{1}})^{i} \end{bmatrix},$$

where we have assumed  $\xi_1 \neq 0$ . Since all the fractions  $\lambda_j/\lambda_1$  are less than 1 in absolute value,  $A^i x_0$  becomes more and more nearly parallel to  $e_1$ , so  $x_i = A^i x_0/||A^i x_0||_2$  becomes closer and closer to  $\pm e_i$ , the eigenvector corresponding to the largest eigenvalue  $\lambda_1$ . The rate of convergence depends on how much smaller than 1 the ratios  $|\lambda_2/\lambda_1| \geq \cdots \geq |\lambda_n/\lambda_1|$ ) are, the smaller the faster. Since  $x_i$  converges to  $\pm e_1$ ,  $\tilde{\lambda}_i = x_i^T A x_i$  converges to  $\lambda_1$ , the largest eigenvalue.

### Assumptions on power method

- In the simplest case we applied algorithm when A is diagonal.
- Consider now general case: assume that  $A = S\Lambda S^{-1}$  is diagonalizable, with  $\Lambda = \operatorname{diag}(\lambda_1, \ldots, \lambda_n)$  and the eigenvalues sorted so that  $|\lambda_1| > |\lambda_2| \ge \cdots \ge |\lambda_n|$ . Write  $S = [s_1, \ldots, s_n]$ , where the columns  $s_i$  are the corresponding eigenvectors and also satisfy  $||s_i||_2 = 1$ ; in the previous slide we had S = I. This lets us write  $\mathbf{x_0} = \mathbf{S}(\mathbf{S}^{-1}\mathbf{x_0}) \equiv \mathbf{S}([\xi_1, \ldots, \xi_n]^T)$ . Also, since  $A = S\Lambda S^{-1}$ , we can write

$$A^{i} = \underbrace{(S\Lambda S^{-1})\cdots(S\Lambda S^{-1})}_{i \text{ times}} = S\Lambda^{i}S^{-1}$$

since all the  $S^{-1} \cdot S$  pairs cancel.

#### This finally lets us write

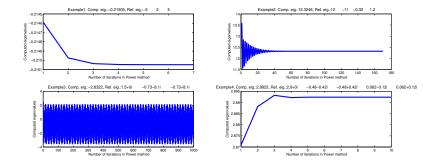
$$A^{i}x_{0} = \underbrace{S\Lambda^{i}S^{-1}}_{A^{i}} \cdot \underbrace{S\left[\begin{array}{c} \xi_{1} \\ \xi_{2} \\ \vdots \\ \xi_{n} \end{array}\right]}_{x_{0}} = S\left[\begin{array}{c} \xi_{1}\lambda_{1}^{i} \\ \xi_{2}\lambda_{2}^{i} \\ \vdots \\ \xi_{n}\lambda_{n}^{i} \end{array}\right] = \xi_{1}\lambda_{1}^{i}S\left[\begin{array}{c} 1 \\ \frac{\xi_{2}}{\xi_{1}}\left(\frac{\lambda_{2}}{\lambda_{1}}\right)^{i} \\ \vdots \\ \frac{\xi_{n}}{\xi_{1}}\left(\frac{\lambda_{n}}{\lambda_{1}}\right)^{i} \end{array}\right].$$

As before, the vector j in brackets converges to  $e_1$ , so  $A^i x_0$  gets closer and closer to a multiple of  $Se_1 = s_1$ , the eigenvector corresponding to  $\lambda_1$ . Therefore,  $\tilde{\lambda}_i = x_i^T A x_i$  converges to  $s_1^T A s_1 = s_1^T \lambda_1 s_1 = \lambda_1$ .

- A minor drawback of this method is the assumption that ξ<sub>1</sub> ≠ 0, i.e., that x<sub>0</sub> is not the invariant subspace span{s<sub>2</sub>,..., s<sub>n</sub>} this is true with very high probability if x<sub>0</sub> is chosen at random.
- A major drawback is that it converges to the eigenvalue/eigenvector pair only for the eigenvalue of largest absolute magnitude, and its convergence rate depends on |λ<sub>2</sub>/λ<sub>1</sub>|, a quantity which may be close to 1 and thus cause very slow convergence. Indeed, if A is real and the largest eigenvalue is complex, there are two complex conjugate eigenvalues of largest absolute value |λ<sub>1</sub>| = |λ<sub>2</sub>|, and so the above analysis does not work at all. In the extreme case of an orthogonal matrix, all the eigenvalues have the same absolute value, namely, 1.
- To plot the convergence of the power method, see HOMEPAGE/Matlab/powerplot.m.

Nonsymmetric Eigenvalue Problem

## Examples of running of Power method in Matlab



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Results are described in the next two slides.

Larisa Beilina, http://www.math.chalmers.se/~larisa/

Example 1. In this example we test the matrix

$$A = \left[ \begin{array}{rrrr} 5 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & -5 \end{array} \right]$$

with exact eigenvalues (5, 2, -5). The Power method can converge as to the exact first eigenvalue 5 as well as to the completely erroneous eigenvalue, see Figure. This is because two eigenvalues of this matrix, 5 and -5, have the same absolute values: |5| = |-5|. Thus, assumption 2 about convergence of the Power method is not fulfilled.

Example 2. In this example the matrix A is given by

$$A = \begin{bmatrix} 3 & 7 & 8 & 9 \\ 5 & -7 & 4 & -7 \\ 1 & -1 & 1 & -1 \\ 9 & 3 & 2 & 5 \end{bmatrix}$$

This matrix has four different reference real eigenvalues  $\lambda = (\lambda_1, ..., \lambda_4)$  given by

 $\lambda = (12.3246, -11.1644, -0.3246, 1.1644).$ 

Now all assumptions about matrix A are fulfilled and the Power method has converged to the reference eigenvalue 12.3246, see Figure. We note that reference eigenvalues on this figure are rounded.

Example 3. Now we take the matrix

$$A = \left[ \begin{array}{rrrr} 0 & 5 & -2 \\ 6 & 0 & -12 \\ 1 & 3 & 0 \end{array} \right]$$

with one real and two complex eigenvalues:

 $\lambda = (1.4522, -0.7261 + 8.0982i, -0.7261 - 8.0982i).$ 

We observe at Figure that Power method does not converge in this case again. This is because assumption 4 on the convergence of the power method is not fulfilled.

Example 4. In this example the matrix A has size 5 × 5. Elements of this matrix are uniformly distributed pseudorandom numbers on the open interval (0, 1). Using the Figure we observe that in the first round of our computations we have obtained good approximation to the eigenvalue 2.9 thought not all assumptions about Power method are fulfilled. This means that on the second round of computations we can get completely different eroneous eigenvalue. This example 1 where convergence was not achieved.

### Inverse Iteration

We will overcome the drawbacks of the power method just described by applying the power method to  $(A - \sigma I)^{-1}$  instead of A, where  $\sigma$  is called a *shift*. This will let us converge to the eigenvalue closest to  $\sigma$ , rather than just  $\lambda_1$ . This method is called *inverse iteration* or the *inverse power method*.

ALGORITHM. Inverse iteration: Given  $x_0$ , we iterate

$$\begin{split} & i = 0 \\ & repeat \\ & y_{i+1} = (A - \sigma I)^{-1} x_i \\ & x_{i+1} = y_{i+1} / ||y_{i+1}||_2 \quad (approximate \ eigenvector) \\ & \tilde{\lambda}_{i+1} = x_{i+1}^T A x_{i+1} \quad (approximate \ eigenvalue) \\ & i = i+1 \\ & until \ convergence \end{split}$$