

MATHEMATICAL SCIENCES  
Chalmers and Göteborg University

Date: 22 October 2005  
Material allowed: only the three  
attached pages with formulas  
Time: 8.30 – 13.30  
Christoffer Cromvik,  
ph. 0762-721860, will come  
at about 9.30 and 12.30

## Exam in Fourier Analysis

MAN 530, TMA362

1. State and prove the theorem about the best approximation of a function  $f \in PC[a, b]$  by means of the functions of an orthonormal system  $\{\phi_1, \dots, \phi_N\}$ .
2. Deduce by means of the Fourier transform the solution of the heat equation in the half-plane  $\{(x, t) : t > 0\}$ , with given initial values  $f(x)$  for  $t = 0$ .
3. Find the Fourier series of the function  $f(x) = x^3$  in the interval  $[-\ell, \ell]$ , for example in the following way: Choose a constant  $c$  in such a way that the Fourier series of  $g(x) = f(x) - cx$  can be differentiated term by term, and determine first the Fourier series of  $g'$ .
4. Solve the initial value problem

$$\begin{aligned}u_{tt}(x, t) &= c^2 u_{xx}(x, t), & 0 < x < \pi, & \quad t > 0 \\u(0, t) &= 0, & u(\pi, t) &= 0 \\u(x, 0) &= 0 \\u_t(x, 0) &= x \sin x.\end{aligned}$$

Here  $c > 0$  is a constant.

5. Solve the following Dirichlet problem in the rectangle  $0 < x < \ell$ ,  $0 < y < L$

$$\begin{aligned}\Delta u &= 0 \\u(0, y) &= 0, & u(\ell, y) &= 1 \\u(x, 0) &= 1, & u(x, L) &= 1.\end{aligned}$$

6. How many eigenvalues  $\lambda$  with  $\lambda < 50$  does the following Sturm-Liouville problem have?

$$f'' + \lambda f = 0, \quad f(0) = 0, \quad f'(1) = -f(1)$$

The grading will be finished by 8 November. You may then get your exam paper at the reception (mottagningsrummet) weekdays from 12.30 to 13.00.

**Solutions to  
exam in Fourier Analysis**

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3.

The function  $g(x) = x^3 - cx$  is defined and odd in the interval  $[-\ell, \ell]$ . In order to allow differentiation term by term of its Fourier series, we must choose the value of  $c$  so as to make the  $2\ell$ -periodic extension of  $g$  continuous on the line. This will be true if  $g(\ell) = 0$ , which implies  $c = \ell^2$ . With  $c$  determined in this way,  $g$  will extend to a piecewise smooth and continuous function on the line, and the theorem about differentiation of Fourier series applies. So if the Fourier series of  $g$  is

$$g(x) = \sum_1^{\infty} b_n \sin \frac{n\pi x}{\ell},$$

that of  $g'$  is

$$g'(x) = \sum_1^{\infty} \frac{n\pi}{\ell} b_n \cos \frac{n\pi x}{\ell}.$$

But  $g'(x) = 3x^2 - \ell^2$ , and the Fourier series of this even function can be obtained from entry 16 of the Fourier series table, with  $\theta = \pi x/\ell$ . One finds

$$3x^2 - \ell^2 = \frac{12\ell^2}{\pi^2} \sum_1^{\infty} \frac{(-1)^n}{n^2} \cos \frac{n\pi x}{\ell}.$$

Identifying coefficients, we have

$$b_n = \frac{12\ell^3}{\pi^3} \frac{(-1)^n}{n^3}.$$

Thus we have the Fourier series of  $g$ , and we need also that of the term  $\ell^2 x$ . Entry 1 of the table shows that

$$\ell^2 x = \frac{2\ell^3}{\pi} \sum_1^{\infty} \frac{(-1)^{n+1}}{n} \sin \frac{n\pi x}{\ell}.$$

Summing up, we get

$$x^3 = \ell^3 \sum_1^{\infty} (-1)^n \left( \frac{12}{\pi^3 n^3} - \frac{2}{\pi n} \right) \sin \frac{n\pi x}{\ell}.$$

It is of course also possible to compute the coefficients by straightforward integration.

4.

This is a standard problem for the wave equation, and one knows that the solution is given by an expression of type

$$u(x, t) = \sum_1^{\infty} b_n \sin nx \sin nct.$$

Notice that there will be no cosine term in  $nct$ , since the initial values  $u(x, 0)$  are 0. The initial values of  $u_t$  lead to

$$\sum_1^{\infty} ncb_n \sin nx = x \sin x, \quad 0 < x < \pi.$$

To find the Fourier sine series expansion of  $x \sin x$ , one can compute the coefficients directly. But we choose another method, and start by expanding  $x$  in a cosine series. From entry 2 in the table, we have

$$x = \frac{\pi}{2} - \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{\cos(2n-1)x}{(2n-1)^2}, \quad 0 < x < \pi.$$

Now  $\cos(2n-1)x \sin x = (\sin 2nx - \sin(2n-2)x)/2$ , and so

$$x \sin x = \frac{\pi}{2} \sin x - \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{\sin 2nx}{(2n-1)^2} + \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{\sin(2n-2)x}{(2n-1)^2}.$$

In the last sum here, we replace  $n$  by  $n+1$ . Since the first term is 0, the sum equals

$$\sum_{n=1}^{\infty} \frac{\sin 2nx}{(2n+1)^2}.$$

Inserting this above, we get

$$\begin{aligned} x \sin x &= \frac{\pi}{2} \sin x - \frac{2}{\pi} \sum_{n=1}^{\infty} \left( \frac{1}{(2n-1)^2} - \frac{1}{(2n+1)^2} \right) \sin 2nx \\ &= \frac{\pi}{2} \sin x - \frac{16}{\pi} \sum_{n=1}^{\infty} \frac{n}{(4n^2-1)^2} \sin 2nx. \end{aligned}$$

We conclude that  $b_1 = \pi/2c$ , that  $b_n = 0$  for odd  $n > 1$  and that

$$b_{2n} = -\frac{8}{\pi c(4n^2 - 1)^2}.$$

Finally, we get

$$u(x, t) = \frac{\pi}{2c} \sin x \sin ct - \frac{8}{\pi c} \sum_{n=1}^{\infty} \frac{1}{(4n^2 - 1)^2} \sin 2nx \sin 2nct.$$

5.

Consider first the problem with the boundary condition  $u(\ell, y) = 1$  replaced by the homogeneous condition  $u(\ell, y) = 0$ . By separation of variables, one finds that the solution  $u_1$  of this problem is of the form

$$u_1(x, y) = \sum_{n=1}^{\infty} \sin \frac{n\pi x}{\ell} \left( a_n \sinh \frac{n\pi y}{\ell} + b_n \sinh \frac{n\pi(L-y)}{\ell} \right).$$

The coefficients  $a_n$  and  $b_n$  are determined by the boundary conditions for  $y = 0$  and  $y = L$ . This leads to

$$\sum_{n=1}^{\infty} b_n \sinh \frac{n\pi L}{\ell} \sin \frac{n\pi x}{\ell} = 1$$

and the same equation for  $a_n$ . The expansion in a sine series of the constant 1 is (entry 6 in the table)

$$1 = \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{2n-1} \sin \frac{(2n-1)\pi x}{\ell}, \quad 0 < x < \ell.$$

Thus

$$a_{2n-1} = b_{2n-1} = \frac{4}{\pi} \frac{1}{(2n-1) \sinh (2n-1)\pi L/\ell}, \quad n = 1, 2, \dots$$

and

$$a_{2n} = b_{2n} = 0, \quad n = 1, 2, \dots,$$

so that

$$u_1(x, y) = \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{(2n-1) \sinh (2n-1)\pi L/\ell} \cdot \sin \frac{(2n-1)\pi x}{\ell} \left( \sinh \frac{(2n-1)\pi y}{\ell} + \sinh \frac{(2n-1)\pi(L-y)}{\ell} \right).$$

Next we consider the Dirichlet problem obtained by replacing in the given problem the boundary conditions  $u(x, 0) = u(x, L) = 1$  by  $u(x, 0) = u(x, L) = 0$ . To find the solution  $u_2$ , we can apply the preceding method after swapping the variables  $x$  and  $y$  and also the lengths  $\ell$  and  $L$ . Thus we have

$$u_2(x, y) = \sum_{n=1}^{\infty} \sin \frac{n\pi y}{L} \left( a'_n \sinh \frac{n\pi x}{\ell} + b'_n \sinh \frac{n\pi(\ell - x)}{L} \right).$$

Because of the conditions for  $x = 0$  and  $x = \ell$ , we get expressions for  $a'_n$  analogous to those of  $a_n$  above, whereas all  $b'_n = 0$ . This means that

$$u_2(x, y) = \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{(2n-1) \sinh (2n-1)\pi\ell/L} \sin \frac{(2n-1)\pi y}{L} \sinh \frac{(2n-1)\pi x}{L}.$$

Since the solution  $u$  of the given problem is  $u = u_1 + u_2$ , we finally find

$$\begin{aligned} u(x, y) &= \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{(2n-1) \sinh (2n-1)\pi L/\ell} \sin \frac{(2n-1)\pi x}{\ell} \\ &\quad \cdot \left( \sinh \frac{(2n-1)\pi y}{\ell} + \sinh \frac{(2n-1)\pi(L-y)}{\ell} \right) \\ &+ \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{(2n-1) \sinh (2n-1)\pi\ell/L} \sin \frac{(2n-1)\pi y}{L} \sinh \frac{(2n-1)\pi x}{L}. \end{aligned}$$

We remark that a somewhat simpler method is to observe that the boundary values 1 on all four sides of the rectangle would trivially give a solution which is identically 1 in the rectangle. Thus  $1 - u$  solves a Dirichlet problem with value 1 on one rectangle side and 0 on the other three sides. In this way, one can write expressions for  $1 - u$  and  $u$  of the above type but shorter.

6.

In the case  $\lambda < 0$  we let  $\lambda = -\mu^2$  with  $\mu > 0$ . Then the general solution of the differential equation is  $f(x) = a \cosh \mu x + b \sinh \mu x$ . The first boundary condition implies  $a = 0$ , and the second leads to  $b\mu \cosh \mu = -b \sinh \mu$ . We get  $\tanh \mu = -\mu$ , an equation which has no solution  $\mu > 0$ . Thus there are no negative eigenvalues.

For  $\lambda = 0$  we get solutions  $f(x) = ax + b$ , and the boundary conditions imply  $b = 0$  and  $a = -a$ , so that  $a = b = 0$ . This means that 0 is not an eigenvalue.

For  $0 < \lambda < 50$  we let  $\lambda = \nu^2$  with  $0 < \nu < \sqrt{50}$ . The solutions of the differential equation are  $f(x) = a \cos \nu x + b \sin \nu x$ , and the boundary

conditions imply  $a = 0$  and  $b\nu \cos \nu = -\sin \nu$ , so that  $\tan \nu = -\nu$ . We must thus determine the number of solution of the equation  $\tan \nu = -\nu$  with  $0 < \nu < \sqrt{50}$ . In the  $\nu, y$  plane, this means the number of intersections between the line  $y = -\nu$  and the curve  $y = \tan \nu$ , with  $\nu$  in the indicated interval. It is immediately seen that the line has precisely one intersection with each branch of the tan curve, and in particular one intersection in each interval  $(n - 1/2)\pi < \nu < n\pi$ ,  $n = 1, 2, \dots$ . We must thus determine how many of these intersections that fall in  $0 < \nu < \sqrt{50}$ . But  $\sqrt{50} > 2\pi$ , since (equivalently)  $\pi^2 < 12.5$ . And  $\sqrt{50} < 2.5\pi$ , since  $\pi^2 > 8$ . This means that we get the first two intersections, and the final answer is two.