

# Quantized Differential Operators, Quantized Weyl Algebras, and Quantized Generalized Verma Modules

In memory of Leonid Vaksman. Kristineberg 2-8 June 2019

Hans Plesner Jakobsen  
University of Copenhagen

June 5 2019

# The Harmonic Representation (Segal-Shale-Weil, Metaplectic,...)

# The Harmonic Representation (Segal-Shale-Weil, Metaplectic,...)

A (very) short history:

# The Harmonic Representation (Segal-Shale-Weil, Metaplectic,...)

A (very) short history:

Van Hove, Léon Sur certaines représentations unitaires d'un groupe infini de transformations, Acad. Roy. Belgique. Cl. Sci. Mém. Coll. in 8° **26**, 5-97 (1951). no. 6

# The Harmonic Representation (Segal-Shale-Weil, Metaplectic,...)

A (very) short history:

Van Hove, Léon Sur certaines représentations unitaires d'un groupe infini de transformations, Acad. Roy. Belgique. Cl. Sci. Mém.

Coll. in 8° **26**, 5-97 (1951). no. 6

Irving E. Segal wrote the entry in Math. Rew.

# The Harmonic Representation (Segal-Shale-Weil, Metaplectic,...)

A (very) short history:

Van Hove, Léon Sur certaines représentations unitaires d'un groupe infini de transformations, Acad. Roy. Belgique. Cl. Sci. Mém.

Coll. in 8° **26**, 5-97 (1951). no. 6

Irving E. Segal wrote the entry in Math. Rew.

D. Shale, A. Weil, K. Gross and R. Kunze, M. Kashiwara and M. Vergne.

# The Harmonic Representation (Segal-Shale-Weil, Metaplectic,...)

A (very) short history:

Van Hove, Léon Sur certaines représentations unitaires d'un groupe infini de transformations, Acad. Roy. Belgique. Cl. Sci. Mém.

Coll. in 8° **26**, 5-97 (1951). no. 6

Irving E. Segal wrote the entry in Math. Rew.

D. Shale, A. Weil, K. Gross and R. Kunze, M. Kashiwara and M. Vergne.

**The Stone-von Neumann Theorem:** There is a unitary irreducible representation of the Heisenberg Group  $H_n$ . It is essentially unique.

# The Harmonic Representation (Segal-Shale-Weil, Metaplectic,...)

A (very) short history:

Van Hove, Léon Sur certaines représentations unitaires d'un groupe infini de transformations, Acad. Roy. Belgique. Cl. Sci. Mém.

Coll. in 8° **26**, 5-97 (1951). no. 6

Irving E. Segal wrote the entry in Math. Rew.

D. Shale, A. Weil, K. Gross and R. Kunze, M. Kashiwara and M. Vergne.

**The Stone-von Neumann Theorem:** There is a unitary irreducible representation of the Heisenberg Group  $H_n$ . It is essentially unique.

The Symplectic Group  $Sp(n, \mathbb{R})$  acts by automorphisms on  $H_n$ .

Due to the uniqueness we get a projective unitary representation of  $Sp(n, \mathbb{R})$ .

# The Harmonic Representation (Segal-Shale-Weil, Metaplectic,...)

A (very) short history:

Van Hove, Léon Sur certaines représentations unitaires d'un groupe infini de transformations, Acad. Roy. Belgique. Cl. Sci. Mém.

Coll. in 8° **26**, 5-97 (1951). no. 6

Irving E. Segal wrote the entry in Math. Rew.

D. Shale, A. Weil, K. Gross and R. Kunze, M. Kashiwara and M. Vergne.

**The Stone-von Neumann Theorem:** There is a unitary irreducible representation of the Heisenberg Group  $H_n$ . It is essentially unique.

The Symplectic Group  $Sp(n, \mathbb{R})$  acts by automorphisms on  $H_n$ .

Due to the uniqueness we get a projective unitary representation of  $Sp(n, \mathbb{R})$ .

Can be lifted to a **genuine unitary representation of  $Mp(n, \mathbb{R})$**  (in the same space).

# The Harmonic Representation (Segal-Shale-Weil, Metaplectic,...)

A (very) short history:

Van Hove, Léon Sur certaines représentations unitaires d'un groupe infini de transformations, Acad. Roy. Belgique. Cl. Sci. Mém.

Coll. in 8° **26**, 5-97 (1951). no. 6

Irving E. Segal wrote the entry in Math. Rew.

D. Shale, A. Weil, K. Gross and R. Kunze, M. Kashiwara and M. Vergne.

**The Stone-von Neumann Theorem:** There is a unitary irreducible representation of the Heisenberg Group  $H_n$ . It is essentially unique.  
The Symplectic Group  $Sp(n, \mathbb{R})$  acts by automorphisms on  $H_n$ .

Due to the uniqueness we get a projective unitary representation of  $Sp(n, \mathbb{R})$ .

Can be lifted to a **genuine unitary representation of  $Mp(n, \mathbb{R})$  (in the same space).**

$Su(n, n) \subseteq Sp(n, \mathbb{R})$  (look at  $Im(\cdot, \cdot)_{n,n}$ )

# The Kashiwara-Vergne Conjecture (KVC)

**KVC:** One gets all unitary highest weight representations of  $SU(n, n)$  from tensor products of the (not irreducible) Harmonic representation.

# The Kashiwara-Vergne Conjecture (KVC)

**KVC:** One gets all unitary highest weight representations of  $SU(n, n)$  from tensor products of the (not irreducible) Harmonic representation.

M. Kashiwara and M. Vergne, On the Segal-Shale-Weil representation and harmonic polynomials, Invent. Math. **44**, 1-47 (1978)

# The Kashiwara-Vergne Conjecture (KVC)

**KVC:** One gets all unitary highest weight representations of  $SU(n, n)$  from tensor products of the (not irreducible) Harmonic representation.

M. Kashiwara and M. Vergne, On the Segal-Shale-Weil representation and harmonic polynomials, Invent. Math. **44**, 1-47 (1978)

This was all global! (on the group level)

# The Kashiwara-Vergne Conjecture (KVC)

**KVC:** One gets all unitary highest weight representations of  $SU(n, n)$  from tensor products of the (not irreducible) Harmonic representation.

M. Kashiwara and M. Vergne, On the Segal-Shale-Weil representation and harmonic polynomials, Invent. Math. **44**, 1-47 (1978)

This was all global! (on the group level)

Idea of the proof of the KVC and ideas therefrom: The Harmonic representation contains “fewer”  $K$ -types than the generic unitary h.w. representations. So there must be a proper quotient. On the most singular  $K$ -types, the Hermitian form is a first order polynomial in  $\lambda$  (from the center of  $K$ )

# The classical situation for $SU(n, n)$

# The classical situation for $SU(n, n)$

$$\mathfrak{g}^{\mathbb{C}} = \mathfrak{p}^- \oplus \mathfrak{k}^{\mathbb{C}} \oplus \mathfrak{p}^+$$

# The classical situation for $SU(n, n)$

$$\begin{aligned}\mathfrak{g}^{\mathbb{C}} &= \mathfrak{p}^- \oplus \mathfrak{k}^{\mathbb{C}} \oplus \mathfrak{p}^+ \\ \mathfrak{p}^\pm &\text{ abelian } \mathfrak{k} \text{ modules}\end{aligned}$$

# The classical situation for $SU(n, n)$

$$\mathfrak{g}^{\mathbb{C}} = \mathfrak{p}^- \oplus \mathfrak{k}^{\mathbb{C}} \oplus \mathfrak{p}^+$$

$\mathfrak{p}^\pm$  abelian  $\mathfrak{k}$  modules

$$\mathcal{U}(\mathfrak{g}^{\mathbb{C}}) = \mathcal{P}(\mathfrak{p}^-) \cdot \mathcal{U}(\mathfrak{k}^{\mathbb{C}}) \cdot \mathcal{P}(\mathfrak{p}^+)$$

# The classical situation for $SU(n, n)$

$$\mathfrak{g}^{\mathbb{C}} = \mathfrak{p}^- \oplus \mathfrak{k}^{\mathbb{C}} \oplus \mathfrak{p}^+$$

$\mathfrak{p}^\pm$  abelian  $\mathfrak{k}$  modules

$$\mathcal{U}(\mathfrak{g}^{\mathbb{C}}) = \mathcal{P}(\mathfrak{p}^-) \cdot \mathcal{U}(\mathfrak{k}^{\mathbb{C}}) \cdot \mathcal{P}(\mathfrak{p}^+)$$

The Killing form

# The classical situation for $SU(n, n)$

$$\mathfrak{g}^{\mathbb{C}} = \mathfrak{p}^- \oplus \mathfrak{k}^{\mathbb{C}} \oplus \mathfrak{p}^+$$

$\mathfrak{p}^\pm$  abelian  $\mathfrak{k}$  modules

$$\mathcal{U}(\mathfrak{g}^{\mathbb{C}}) = \mathcal{P}(\mathfrak{p}^-) \cdot \mathcal{U}(\mathfrak{k}^{\mathbb{C}}) \cdot \mathcal{P}(\mathfrak{p}^+)$$

The Killing form

$$\mathcal{P}(\mathfrak{p}^\pm) = \mathcal{P}^\pm = \text{holomorphic polynomials on } \mathfrak{p}^\mp$$

# The classical situation for $SU(n, n)$

$$\mathfrak{g}^{\mathbb{C}} = \mathfrak{p}^- \oplus \mathfrak{k}^{\mathbb{C}} \oplus \mathfrak{p}^+$$

$\mathfrak{p}^\pm$  abelian  $\mathfrak{k}$  modules

$$\mathcal{U}(\mathfrak{g}^{\mathbb{C}}) = \mathcal{P}(\mathfrak{p}^-) \cdot \mathcal{U}(\mathfrak{k}^{\mathbb{C}}) \cdot \mathcal{P}(\mathfrak{p}^+)$$

The Killing form

$$\mathcal{P}(\mathfrak{p}^\pm) = \mathcal{P}^\pm = \text{holomorphic polynomials on } \mathfrak{p}^\mp$$

$$\mathcal{P}(\mathfrak{p}^\pm) = \mathcal{P}^\pm = \text{holomorphic differential operators on } \mathfrak{p}^\pm$$

$V_\Lambda$  a finite-dimensional (unitary) representation of  $\mathfrak{k}^{\mathbb{C}}$ ).

# The classical situation for $SU(n, n)$

$$\mathfrak{g}^{\mathbb{C}} = \mathfrak{p}^- \oplus \mathfrak{k}^{\mathbb{C}} \oplus \mathfrak{p}^+$$

$\mathfrak{p}^\pm$  abelian  $\mathfrak{k}$  modules

$$\mathcal{U}(\mathfrak{g}^{\mathbb{C}}) = \mathcal{P}(\mathfrak{p}^-) \cdot \mathcal{U}(\mathfrak{k}^{\mathbb{C}}) \cdot \mathcal{P}(\mathfrak{p}^+)$$

The Killing form

$$\mathcal{P}(\mathfrak{p}^\pm) = \mathcal{P}^\pm = \text{holomorphic polynomials on } \mathfrak{p}^\mp$$

$$\mathcal{P}(\mathfrak{p}^\pm) = \mathcal{P}^\pm = \text{holomorphic differential operators on } \mathfrak{p}^\pm$$

$V_\Lambda$  a finite-dimensional (unitary) representation of  $\mathfrak{k}^{\mathbb{C}}$ ).

# The classical situation for $SU(n, n)$

$$\mathfrak{g}^{\mathbb{C}} = \mathfrak{p}^- \oplus \mathfrak{k}^{\mathbb{C}} \oplus \mathfrak{p}^+$$

$\mathfrak{p}^\pm$  abelian  $\mathfrak{k}$  modules

$$\mathcal{U}(\mathfrak{g}^{\mathbb{C}}) = \mathcal{P}(\mathfrak{p}^-) \cdot \mathcal{U}(\mathfrak{k}^{\mathbb{C}}) \cdot \mathcal{P}(\mathfrak{p}^+)$$

The Killing form

$$\mathcal{P}(\mathfrak{p}^\pm) = \mathcal{P}^\pm = \text{holomorphic polynomials on } \mathfrak{p}^\mp$$

$$\mathcal{P}(\mathfrak{p}^\pm) = \mathcal{P}^\pm = \text{holomorphic differential operators on } \mathfrak{p}^\pm$$

$V_\Lambda$  a finite-dimensional (unitary) representation of  $\mathfrak{k}^{\mathbb{C}}$ ).

$$\Lambda \leftrightarrow (\Lambda_0, \lambda).$$

The generalized Verma module  $M(V_\Lambda)$ :

$$M(V_\Lambda) = \mathcal{U}(\mathfrak{g}^{\mathbb{C}}) \bigotimes_{\mathcal{U}(\mathfrak{k}^{\mathbb{C}} \oplus \mathfrak{p}^+)} V_\Lambda \quad (1)$$

# Details of holomorphically induced representations

## Details of holomorphically induced representations

$\mathcal{P}(V_\tau) :=$  Holomorphic polynomials with values in  $V_\tau$ .

## Details of holomorphically induced representations

$\mathcal{P}(V_\tau) :=$  Holomorphic polynomials with values in  $V_\tau$ .

$$\begin{aligned}\mathcal{P}(V_\tau) &\equiv \\ \{f : G \rightarrow V_\tau \mid & \forall W \in \mathfrak{p}^+ : r(W)f = 0 \\ & \forall g \in G \forall k \in K \quad f(gk) = \tau(k^{-1})f(g)\}.\end{aligned}$$

## Details of holomorphically induced representations

$\mathcal{P}(V_\tau) :=$  Holomorphic polynomials with values in  $V_\tau$ .

$$\begin{aligned}\mathcal{P}(V_\tau) &\equiv \\ \{f : G \rightarrow V_\tau \mid & \forall W \in \mathfrak{p}^+ : r(W)f = 0 \\ & \forall g \in G \forall k \in K \quad f(gk) = \tau(k^{-1})f(g)\}.\end{aligned}$$

**Theorem** There is a nondegenerate pairing between the modules  $\mathcal{P}(V_\tau)$  and  $M(V_{\tau'})$ .

## Details of holomorphically induced representations

$\mathcal{P}(V_\tau) :=$  Holomorphic polynomials with values in  $V_\tau$ .

$$\begin{aligned}\mathcal{P}(V_\tau) &\equiv \\ \{f : G \rightarrow V_\tau \mid & \forall W \in \mathfrak{p}^+ : r(W)f = 0 \\ & \forall g \in G \forall k \in K \quad f(gk) = \tau(k^{-1})f(g)\}.\end{aligned}$$

**Theorem** There is a nondegenerate pairing between the modules  $\mathcal{P}(V_\tau)$  and  $M(V_{\tau'})$ .

Specifically, as a vector space,  $M(V_{\tau'})$  can be seen as the space of (constant coefficient)  $V_{\tau'}$ -valued holomorphic differential operators on  $\mathfrak{p}^-$ . (Domain  $G/K = \mathcal{D} \subseteq \mathfrak{p}^-$ .)

$$(q, p) = (q(\frac{\partial}{\partial Z})p)(0).$$

## Homomorphisms

$\phi : M(V_{\tau'_1}) \rightarrow M(V_{\tau'})$  homomorphism

determined as follows:  $\{f_1, f_2, \dots, f_{N_1}\}$  and  $\{e_1, e_2, \dots, e_N\}$  bases of  $V_{\tau'_1}$  and  $V_{\tau'}$ , respectively. Define the  $N \times N_1$  matrix  $\Phi$  as follows:

$$\Phi_{ij} = \phi(f)_{ith \text{ coord.}} \in \mathcal{U}(\mathfrak{p}^-).$$

$$M(V_{\tau'_1}) \ni q \rightarrow \Phi q \in M(V_{\tau}).$$

## Homomorphisms

$\phi : M(V_{\tau'_1}) \rightarrow M(V_{\tau'})$  homomorphism

determined as follows:  $\{f_1, f_2, \dots, f_{N_1}\}$  and  $\{e_1, e_2, \dots, e_N\}$  bases of  $V_{\tau'_1}$  and  $V_{\tau'}$ , respectively. Define the  $N \times N_1$  matrix  $\Phi$  as follows:

$$\Phi_{ij} = \phi(f)_{ith \text{ coord.}} \in \mathcal{U}(\mathfrak{p}^-).$$

$$M(V_{\tau'_1}) \ni q \rightarrow \Phi q \in M(V_{\tau}).$$

$$M(V_{\tau'_1}) \ni \begin{pmatrix} q_1 \\ \vdots \\ q \\ \vdots \\ q_{N_1} \end{pmatrix} \rightarrow \begin{pmatrix} \Phi_{11} & \dots & \Phi_{1j} & \dots & \Phi_{1N_1} \\ \vdots & \dots & \vdots & \dots & \vdots \\ \Phi_{i1} & \dots & \Phi_{ij} & \dots & \Phi_{iN_1} \\ \vdots & \dots & \vdots & \dots & \vdots \\ \Phi_{N1} & \dots & \Phi_{Nj} & \dots & \Phi_{NN_1} \end{pmatrix} \begin{pmatrix} q_1 \\ \vdots \\ q \\ \vdots \\ q_{N_1} \end{pmatrix} \in M(V_{\tau'})$$

# Covariant differential operators

# Covariant differential operators

By duality we get a covariant differential operator

$$\mathbb{D}_{\Psi} : \mathcal{P}(V_{\tau}) \rightarrow \mathcal{P}(V_{\tau_1}).$$

## Covariant differential operators

By duality we get a covariant differential operator

$$\mathbb{D}_{\Psi} : \mathcal{P}(V_{\tau}) \rightarrow \mathcal{P}(V_{\tau_1}).$$

$$\mathbb{D}_{\Psi} = \begin{pmatrix} \hat{\Phi}_{11} & \dots & \hat{\Phi}_{i1} & \dots & \hat{\Phi}_{N1} \\ \vdots & \dots & \vdots & \dots & \vdots \\ \hat{\Phi}_{1j} & \dots & \hat{\Phi}_{ij} & \dots & \hat{\Phi}_{Nj} \\ \vdots & \dots & \vdots & \dots & \vdots \\ \hat{\Phi}_{1N_1} & \dots & \hat{\Phi}_{iN_1} & \dots & \hat{\Phi}_{NN_1} \end{pmatrix}$$

## Covariant differential operators

By duality we get a covariant differential operator

$$\mathbb{D}_\Psi : \mathcal{P}(V_\tau) \rightarrow \mathcal{P}(V_{\tau_1}).$$

$$\mathbb{D}_\Psi = \begin{pmatrix} \hat{\Phi}_{11} & \dots & \hat{\Phi}_{i1} & \dots & \hat{\Phi}_{N1} \\ \vdots & \dots & \vdots & \dots & \vdots \\ \hat{\Phi}_{1j} & \dots & \hat{\Phi}_{ij} & \dots & \hat{\Phi}_{Nj} \\ \vdots & \dots & \vdots & \dots & \vdots \\ \hat{\Phi}_{1N_1} & \dots & \hat{\Phi}_{iN_1} & \dots & \hat{\Phi}_{NN_1} \end{pmatrix}$$

**Theorem.** There is a simple bijective correspondence between covariant differential operators and homomorphisms between generalized Verma modules.

## Examples from $su(2, 2)$ (Dirac, Maxwell, ...)

## Examples from $su(2, 2)$ (Dirac, Maxwell, ...)

$$Z = \begin{pmatrix} z_0 + z_1 & z_2 + iz_3 \\ z_2 - iz_3 & z_0 - z_1 \end{pmatrix}; \quad c(Z) = \begin{pmatrix} z_0 - z_1 & -z_2 - iz_3 \\ -z_2 + iz_3 & z_0 + z_1 \end{pmatrix}.$$

## Examples from $su(2, 2)$ (Dirac, Maxwell, ...)

$$Z = \begin{pmatrix} z_0 + z_1 & z_2 + iz_3 \\ z_2 - iz_3 & z_0 - z_1 \end{pmatrix}; \quad c(Z) = \begin{pmatrix} z_0 - z_1 & -z_2 - iz_3 \\ -z_2 + iz_3 & z_0 + z_1 \end{pmatrix}.$$

$$\det Z$$

## Examples from $su(2, 2)$ (Dirac, Maxwell, ...)

$$Z = \begin{pmatrix} z_0 + z_1 & z_2 + iz_3 \\ z_2 - iz_3 & z_0 - z_1 \end{pmatrix}; \quad c(Z) = \begin{pmatrix} z_0 - z_1 & -z_2 - iz_3 \\ -z_2 + iz_3 & z_0 + z_1 \end{pmatrix}.$$

$$\det Z = z_0^2 - z_2^2 - z_2^2 - z_3^2$$

## Examples from $su(2, 2)$ (Dirac, Maxwell, ...)

$$Z = \begin{pmatrix} z_0 + z_1 & z_2 + iz_3 \\ z_2 - iz_3 & z_0 - z_1 \end{pmatrix}; \quad c(Z) = \begin{pmatrix} z_0 - z_1 & -z_2 - iz_3 \\ -z_2 + iz_3 & z_0 + z_1 \end{pmatrix}.$$

$$\det Z = z_0^2 - z_2^2 - z_2^2 - z_3^2 \rightarrow$$

## Examples from $su(2, 2)$ (Dirac, Maxwell, ...)

$$Z = \begin{pmatrix} z_0 + z_1 & z_2 + iz_3 \\ z_2 - iz_3 & z_0 - z_1 \end{pmatrix}; \quad c(Z) = \begin{pmatrix} z_0 - z_1 & -z_2 - iz_3 \\ -z_2 + iz_3 & z_0 + z_1 \end{pmatrix}.$$

$$\det Z = z_0^2 - z_2^2 - z_2^2 - z_3^2 \rightarrow \square$$

## Examples from $su(2, 2)$ (Dirac, Maxwell, ...)

$$Z = \begin{pmatrix} z_0 + z_1 & z_2 + iz_3 \\ z_2 - iz_3 & z_0 - z_1 \end{pmatrix}; \quad c(Z) = \begin{pmatrix} z_0 - z_1 & -z_2 - iz_3 \\ -z_2 + iz_3 & z_0 + z_1 \end{pmatrix}.$$

$$\det Z = z_0^2 - z_1^2 - z_2^2 - z_3^2 \rightarrow \square = \frac{\partial^2}{\partial z_0^2} - \frac{\partial^2}{\partial z_1^2} - \frac{\partial^2}{\partial z_2^2} - \frac{\partial^2}{\partial z_3^2}$$

## Examples from $su(2, 2)$ (Dirac, Maxwell, ...)

$$Z = \begin{pmatrix} z_0 + z_1 & z_2 + iz_3 \\ z_2 - iz_3 & z_0 - z_1 \end{pmatrix}; \quad c(Z) = \begin{pmatrix} z_0 - z_1 & -z_2 - iz_3 \\ -z_2 + iz_3 & z_0 + z_1 \end{pmatrix}.$$

$$\det Z = z_0^2 - z_1^2 - z_2^2 - z_3^2 \rightarrow \square = \frac{\partial^2}{\partial z_0^2} - \frac{\partial^2}{\partial z_1^2} - \frac{\partial^2}{\partial z_2^2} - \frac{\partial^2}{\partial z_3^2}$$

$$\nabla = \begin{pmatrix} 0 & 0 & \frac{\partial}{\partial z_0} + \frac{\partial}{\partial z_1} & \frac{\partial}{\partial z_2} + i \frac{\partial}{\partial z_3} \\ 0 & 0 & \frac{\partial}{\partial z_2} - i \frac{\partial}{\partial z_3} & \frac{\partial}{\partial z_0} - \frac{\partial}{\partial z_1} \\ \frac{\partial}{\partial z_0} - \frac{\partial}{\partial z_1} & -\frac{\partial}{\partial z_2} - i \frac{\partial}{\partial z_3} & 0 & 0 \\ -\frac{\partial}{\partial z_2} + i \frac{\partial}{\partial z_3} & \frac{\partial}{\partial z_0} + \frac{\partial}{\partial z_1} & 0 & 0 \end{pmatrix}.$$

## Examples from $su(2, 2)$ (Dirac, Maxwell, ...)

$$Z = \begin{pmatrix} z_0 + z_1 & z_2 + iz_3 \\ z_2 - iz_3 & z_0 - z_1 \end{pmatrix}; \quad c(Z) = \begin{pmatrix} z_0 - z_1 & -z_2 - iz_3 \\ -z_2 + iz_3 & z_0 + z_1 \end{pmatrix}.$$

$$\det Z = z_0^2 - z_1^2 - z_2^2 - z_3^2 \rightarrow \square = \frac{\partial^2}{\partial z_0^2} - \frac{\partial^2}{\partial z_1^2} - \frac{\partial^2}{\partial z_2^2} - \frac{\partial^2}{\partial z_3^2}$$

$$\nabla = \begin{pmatrix} 0 & 0 & \frac{\partial}{\partial z_0} + \frac{\partial}{\partial z_1} & \frac{\partial}{\partial z_2} + i \frac{\partial}{\partial z_3} \\ 0 & 0 & \frac{\partial}{\partial z_2} - i \frac{\partial}{\partial z_3} & \frac{\partial}{\partial z_0} - \frac{\partial}{\partial z_1} \\ \frac{\partial}{\partial z_0} - \frac{\partial}{\partial z_1} & -\frac{\partial}{\partial z_2} - i \frac{\partial}{\partial z_3} & 0 & 0 \\ -\frac{\partial}{\partial z_2} + i \frac{\partial}{\partial z_3} & \frac{\partial}{\partial z_0} + \frac{\partial}{\partial z_1} & 0 & 0 \end{pmatrix}.$$

$1 \otimes \nabla + \nabla \otimes 1$  (on symmetric tensors).

## Examples from $su(2, 2)$ (Dirac, Maxwell, ...)

$$Z = \begin{pmatrix} z_0 + z_1 & z_2 + iz_3 \\ z_2 - iz_3 & z_0 - z_1 \end{pmatrix}; \quad c(Z) = \begin{pmatrix} z_0 - z_1 & -z_2 - iz_3 \\ -z_2 + iz_3 & z_0 + z_1 \end{pmatrix}.$$

$$\det Z = z_0^2 - z_1^2 - z_2^2 - z_3^2 \rightarrow \square = \frac{\partial^2}{\partial z_0^2} - \frac{\partial^2}{\partial z_1^2} - \frac{\partial^2}{\partial z_2^2} - \frac{\partial^2}{\partial z_3^2}$$

$$\nabla = \begin{pmatrix} 0 & 0 & \frac{\partial}{\partial z_0} + \frac{\partial}{\partial z_1} & \frac{\partial}{\partial z_2} + i \frac{\partial}{\partial z_3} \\ 0 & 0 & \frac{\partial}{\partial z_2} - i \frac{\partial}{\partial z_3} & \frac{\partial}{\partial z_0} - \frac{\partial}{\partial z_1} \\ \frac{\partial}{\partial z_0} - \frac{\partial}{\partial z_1} & -\frac{\partial}{\partial z_2} - i \frac{\partial}{\partial z_3} & 0 & 0 \\ -\frac{\partial}{\partial z_2} + i \frac{\partial}{\partial z_3} & \frac{\partial}{\partial z_0} + \frac{\partial}{\partial z_1} & 0 & 0 \end{pmatrix}.$$

$1 \otimes \nabla + \nabla \otimes 1$  (on symmetric tensors).

$1 \otimes 1 \otimes 1 \otimes \cdots \otimes 1 \otimes \nabla + 1 \otimes 1 \otimes \cdots \otimes \nabla \otimes 1 + \cdots + \nabla \otimes 1 \otimes 1 \otimes \cdots \otimes 1 \otimes (\text{sym}).$

# The Harmonic Representation Returns

# The Harmonic Representation Returns

The infinitesimal version:

# The Harmonic Representation Returns

The infinitesimal version:

The Heisenberg Lie algebra  $\mathfrak{h}_{2n}$ :

# The Harmonic Representation Returns

The infinitesimal version:

The Heisenberg Lie algebra  $\mathfrak{h}_{2n}$ :

$$[Q_k, P_\ell] = \delta_{k,\ell} C.$$

# The Harmonic Representation Returns

The infinitesimal version:

The Heisenberg Lie algebra  $\mathfrak{h}_{2n}$ :

$$[Q_k, P_\ell] = \delta_{k,\ell} C.$$

$$Q_k = x_k, \text{ and } P_\ell = -i \frac{\partial}{\partial x_\ell}; \quad [x_k, -i \frac{\partial}{\partial x_\ell}] = i \delta_{k,\ell}.$$

# The Harmonic Representation Returns

The infinitesimal version:

The Heisenberg Lie algebra  $\mathfrak{h}_{2n}$ :

$$[Q_k, P_\ell] = \delta_{k,\ell} C.$$

$$Q_k = x_k, \text{ and } P_\ell = -i \frac{\partial}{\partial x_\ell}; \quad [x_k, -i \frac{\partial}{\partial x_\ell}] = i \delta_{k,\ell}.$$

Annihilation operators:  $A_k = \frac{1}{\sqrt{2}}(Q_k + iP_k)$

Creation operators:  $A_\ell^\dagger = \frac{1}{\sqrt{2}}(Q_\ell + iP_\ell)$

$$[A_k, A_\ell^\dagger] = \delta_{k,\ell}.$$

# The Harmonic Representation Returns

The infinitesimal version:

The Heisenberg Lie algebra  $\mathfrak{h}_{2n}$ :

$$[Q_k, P_\ell] = \delta_{k,\ell} C.$$

$$Q_k = x_k, \text{ and } P_\ell = -i \frac{\partial}{\partial x_\ell}; \quad [x_k, -i \frac{\partial}{\partial x_\ell}] = i \delta_{k,\ell}.$$

Annihilation operators:  $A_k = \frac{1}{\sqrt{2}}(Q_k + iP_k)$

Creation operators:  $A_\ell^\dagger = \frac{1}{\sqrt{2}}(Q_\ell + iP_\ell)$

$$[A_k, A_\ell^\dagger] = \delta_{k,\ell}.$$

# The (infinitesimal) Stone - von Neumann representation

# The (infinitesimal) Stone - von Neumann representation

$v_0$  is the “vacuum vector”:

# The (infinitesimal) Stone - von Neumann representation

$v_0$  is the “vacuum vector”:

$$A_k v_0 = 0 \text{ for all } k = 1, 2, \dots, 2n$$

# The (infinitesimal) Stone - von Neumann representation

$v_0$  is the “vacuum vector”:

$$A_k v_0 = 0 \text{ for all } k = 1, 2, \dots, 2n$$

$$\mathcal{H} = \text{Span}_{\mathbb{C}} \{(A^\dagger)_1^{r_1} (A^\dagger)_2^{r_2} \cdots (A^\dagger)_{2n}^{r_{2n}} v_0 \mid r_1, \dots, r_{2n} \in \mathbb{N}_0\}.$$

# The (infinitesimal) Stone - von Neumann representation

$v_0$  is the “vacuum vector”:

$$A_k v_0 = 0 \text{ for all } k = 1, 2, \dots, 2n$$

$$\mathcal{H} = \text{Span}_{\mathbb{C}} \{(A^\dagger)_1^{r_1} (A^\dagger)_2^{r_2} \cdots (A^\dagger)_{2n}^{r_{2n}} v_0 \mid r_1, \dots, r_{2n} \in \mathbb{N}_0\}.$$

This is a unitary representation.

# The (infinitesimal) Stone - von Neumann representation

$v_0$  is the “vacuum vector”:

$$A_k v_0 = 0 \text{ for all } k = 1, 2, \dots, 2n$$

$$\mathcal{H} = \text{Span}_{\mathbb{C}} \{(A^\dagger)_1^{r_1} (A^\dagger)_2^{r_2} \cdots (A^\dagger)_{2n}^{r_{2n}} v_0 \mid r_1, \dots, r_{2n} \in \mathbb{N}_0\}.$$

This is a unitary representation.

A simple but **key** Lemma:

$$(A^\dagger)^2 A - 2(A^\dagger A A^\dagger) + A(A^\dagger)^2 = 0$$
$$A^2 A^\dagger - 2(A A^\dagger A) + A^\dagger A^2 = 0$$

## Constructing $su(n, n)$

Set

$$\begin{aligned} e &= A_n^\dagger A_{n+1}^\dagger \quad ; \quad f = A_n A_{n+1} \\ H_k &= A_k A_k^\dagger \quad ; \quad h_e = H_n + H_{n+1} - 1 \\ \mu_k^+ &= A_k^\dagger A_{k+1} \quad ; \quad \mu_k^- = A_k A_{k+1}^\dagger \quad (k < n) \\ \nu_k^+ &= A_{k+1}^\dagger A_k \quad ; \quad \nu_k^- = A_{k+1} A_k^\dagger \quad (k > n). \end{aligned}$$

## Constructing $su(n, n)$

Set

$$\begin{aligned} e &= A_n^\dagger A_{n+1}^\dagger \quad ; \quad f = A_n A_{n+1} \\ H_k &= A_k A_k^\dagger \quad ; \quad h_e = H_n + H_{n+1} - 1 \\ \mu_k^+ &= A_k^\dagger A_{k+1} \quad ; \quad \mu_k^- = A_k A_{k+1}^\dagger \quad (k < n) \\ \nu_k^+ &= A_{k+1}^\dagger A_k \quad ; \quad \nu_k^- = A_{k+1} A_k^\dagger \quad (k > n). \end{aligned}$$

Then these satisfy the defining relations for  $su(n, n)$ .

## Constructing $su(n, n)$

Set

$$\begin{aligned} e &= A_n^\dagger A_{n+1}^\dagger \quad ; \quad f = A_n A_{n+1} \\ H_k &= A_k A_k^\dagger \quad ; \quad h_e = H_n + H_{n+1} - 1 \\ \mu_k^+ &= A_k^\dagger A_{k+1} \quad ; \quad \mu_k^- = A_k A_{k+1}^\dagger \quad (k < n) \\ \nu_k^+ &= A_{k+1}^\dagger A_k \quad ; \quad \nu_k^- = A_{k+1} A_k^\dagger \quad (k > n). \end{aligned}$$

Then these satisfy the defining relations for  $su(n, n)$ .

Indeed, we get the Serre relations, e.g.

$$e^2 \mu_{n-1}^+ - 2e \mu_{n-1}^+ e + \mu_{n-1}^+ e^2 = 0 \quad \text{etc.}$$

## Constructing $su(n, n)$

Set

$$\begin{aligned} e &= A_n^\dagger A_{n+1}^\dagger \quad ; \quad f = A_n A_{n+1} \\ H_k &= A_k A_k^\dagger \quad ; \quad h_e = H_n + H_{n+1} - 1 \\ \mu_k^+ &= A_k^\dagger A_{k+1} \quad ; \quad \mu_k^- = A_k A_{k+1}^\dagger \quad (k < n) \\ \nu_k^+ &= A_{k+1}^\dagger A_k \quad ; \quad \nu_k^- = A_{k+1} A_k^\dagger \quad (k > n). \end{aligned}$$

Then these satisfy the defining relations for  $su(n, n)$ .

Indeed, we get the Serre relations, e.g.

$$e^2 \mu_{n-1}^+ - 2e \mu_{n-1}^+ e + \mu_{n-1}^+ e^2 = 0 \quad \text{etc.}$$

$$Z_{k,\ell} = iA_{n-k+1}^\dagger A_{n+\ell}^\dagger \quad \text{and} \quad W_{k,\ell} = iA_{n-k+1} A_{n+\ell}.$$

# Minors

Set

$$X = \begin{pmatrix} A_1 \\ A_2 \\ \dots \\ A_n \end{pmatrix} \text{ and } Y = \begin{pmatrix} A_{n+1} \\ A_{n+2} \\ \dots \\ A_{2n} \end{pmatrix}$$

Then  $W = X \cdot Y^T$

# Minors

Set

$$X = \begin{pmatrix} A_1 \\ A_2 \\ \dots \\ A_n \end{pmatrix} \text{ and } Y = \begin{pmatrix} A_{n+1} \\ A_{n+2} \\ \dots \\ A_{2n} \end{pmatrix}$$

Then  $W = X \cdot Y^T$  So  $W$  is rank 1.

# Minors

Set

$$X = \begin{pmatrix} A_1 \\ A_2 \\ \dots \\ A_n \end{pmatrix} \text{ and } Y = \begin{pmatrix} A_{n+1} \\ A_{n+2} \\ \dots \\ A_{2n} \end{pmatrix}$$

Then  $W = X \cdot Y^T$  So  $W$  is rank 1.

**Proposition.** All  $2 \times 2$  minors in the variables  $W_{k,\ell}$  vanish.

# Minors

Set

$$X = \begin{pmatrix} A_1 \\ A_2 \\ \dots \\ A_n \end{pmatrix} \text{ and } Y = \begin{pmatrix} A_{n+1} \\ A_{n+2} \\ \dots \\ A_{2n} \end{pmatrix}$$

Then  $W = X \cdot Y^T$  So  $W$  is rank 1.

**Proposition.** All  $2 \times 2$  minors in the variables  $W_{k,\ell}$  vanish.

More generally, set  $X_k = (X^1, X^2, \dots, X^k)$  and  
 $Y = (Y^1, Y^2, \dots, Y^k)$  (these are then  $n \times k$  matrices).

Then

# Minors

Set

$$X = \begin{pmatrix} A_1 \\ A_2 \\ \dots \\ A_n \end{pmatrix} \text{ and } Y = \begin{pmatrix} A_{n+1} \\ A_{n+2} \\ \dots \\ A_{2n} \end{pmatrix}$$

Then  $W = X \cdot Y^T$  So  $W$  is rank 1.

**Proposition.** All  $2 \times 2$  minors in the variables  $W_{k,\ell}$  vanish.

More generally, set  $X_k = (X^1, X^2, \dots, X^k)$  and  
 $Y = (Y^1, Y^2, \dots, Y^k)$  (these are then  $n \times k$  matrices).

Then

$$W_k = X_k \cdot Y_k^T$$

is of rank  $k$ , hence all  $(k+1) \times (k+1)$  minors vanish.

# Minors

# Minors

**Corollary.** The ideal generated by all  $k \times k$  minors is prime

## Minors

**Corollary.** The ideal generated by all  $k \times k$  minors is prime  
Mount, K.R., A remark on determinantal loci, J. London Math.  
Soc. 42 (1967), 595–598.

## Minors

**Corollary.** The ideal generated by all  $k \times k$  minors is prime  
Mount, K.R., A remark on determinantal loci, J. London Math.  
Soc. 42 (1967), 595–598.

See also discussion in the end of Chapter 2 in Bruns and Vetter,  
Determinantal Rings [1988]

## Minors

**Corollary.** The ideal generated by all  $k \times k$  minors is prime  
Mount, K.R., A remark on determinantal loci, J. London Math.  
Soc. 42 (1967), 595–598.

See also discussion in the end of Chapter 2 in Bruns and Vetter,  
Determinantal Rings [1988]

So, we immediately get interesting “scalar” modules, but of course, there are many more.

## Minors

**Corollary.** The ideal generated by all  $k \times k$  minors is prime  
Mount, K.R., A remark on determinantal loci, J. London Math.  
Soc. 42 (1967), 595–598.

See also discussion in the end of Chapter 2 in Bruns and Vetter,  
Determinantal Rings [1988]

So, we immediately get interesting “scalar” modules, but of course, there are many more.

Even in the simplest case ( $k = 1$ ), the vectors

$$(A_n^\dagger)^i v_0 \text{ and } (A_{n+1}^\dagger)_0^j$$

are highest weight vectors.

## Minors

**Corollary.** The ideal generated by all  $k \times k$  minors is prime  
Mount, K.R., A remark on determinantal loci, J. London Math.  
Soc. 42 (1967), 595–598.

See also discussion in the end of Chapter 2 in Bruns and Vetter,  
Determinantal Rings [1988]

So, we immediately get interesting “scalar” modules, but of course, there are many more.

Even in the simplest case ( $k = 1$ ), the vectors

$$(A_n^\dagger)^i v_0 \text{ and } (A_{n+1}^\dagger)_0^j$$

are highest weight vectors.

The **ladder series** of mass 0 and spin  $\frac{i}{2}$ , resp.  $\frac{j}{2}$ .

Now the  $q$  version

## Now the $q$ version

This is really Vaksman's domain.

## Now the $q$ version

This is really Vaksman's domain.

Structures also related to e.g. those found in

## Now the $q$ version

This is really Vaksman's domain.

Structures also related to e.g. those found in

Olga Bershtein, Olof Giselsson, and Lyudmila Turowska, Maximum modulus principle for "holomorphic functions" on the quantum matrix ball arXiv:1711.05981 [math.OA]

## Now the $q$ version

This is really Vaksman's domain.

Structures also related to e.g. those found in

Olga Bershtein, Olof Giselsson, and Lyudmila Turowska, Maximum modulus principle for "holomorphic functions" on the quantum matrix ball arXiv:1711.05981 [math.OA]

S.Z. Levendorskii and Ya. S. Soibelman, Some applications of the quantum Weyl Groups, J. Geom. Phys. **7** (1990), 241-254.

S. Levendorskii and Y. Soibelman, Algebras of functions on compact quantum groups, Schubert Cells and quantum tori, Comm. Math. Phys **139** (1991), 141-170.

## Now the $q$ version

This is really Vaksman's domain.

Structures also related to e.g. those found in

Olga Bershtein, Olof Giselsson, and Lyudmila Turowska, Maximum modulus principle for "holomorphic functions" on the quantum matrix ball arXiv:1711.05981 [math.OA]

S.Z. Levendorskii and Ya. S. Soibelman, Some applications of the quantum Weyl Groups, J. Geom. Phys. **7** (1990), 241-254.

S. Levendorskii and Y. Soibelman, Algebras of functions on compact quantum groups, Schubert Cells and quantum tori, Comm. Math. Phys **139** (1991), 141-170.

A. Kamita, Y. Morita, T. Tanisaki, Quantum deformations of certain prehomogeneous vector spaces I, Hiroshima Math. J., 28 (1998), No 3, 527 – 540.

## Now the $q$ version

This is really Vaksman's domain.

Structures also related to e.g. those found in

Olga Bershtein, Olof Giselsson, and Lyudmila Turowska, Maximum modulus principle for "holomorphic functions" on the quantum matrix ball arXiv:1711.05981 [math.OA]

S.Z. Levendorskii and Ya. S. Soibelman, Some applications of the quantum Weyl Groups, J. Geom. Phys. **7** (1990), 241-254.

S. Levendorskii and Y. Soibelman, Algebras of functions on compact quantum groups, Schubert Cells and quantum tori, Comm. Math. Phys **139** (1991), 141-170.

A. Kamita, Y. Morita, T. Tanisaki, Quantum deformations of certain prehomogeneous vector spaces I, Hiroshima Math. J., 28 (1998), No 3, 527 – 540.

S. Sinel'shchikov, L. Vaksman, Quantum Groups and Bounded Symmetric Domains, arXiv:math/0410530

## Now the $q$ version

My own history:

## Now the $q$ version

My own history:

Quantized hermitian symmetric spaces. In: "Lie Theory and its Applications in Physics", 105-116, World Scientific, Singapore (1996)

## Now the $q$ version

My own history:

Quantized hermitian symmetric spaces. In: "Lie Theory and its Applications in Physics", 105-116, World Scientific, Singapore (1996)

Q-differential operators. Preprint (19pp) 1999  
(<http://xxx.lanl.gov/abs/math.QA/ 9907009>)

## Now the $q$ version

My own history:

Quantized hermitian symmetric spaces. In: "Lie Theory and its Applications in Physics", 105-116, World Scientific, Singapore (1996)

Q-differential operators. Preprint (19pp) 1999  
(<http://xxx.lanl.gov/abs/math.QA/9907009>)

Quantized Dirac operators. Czech. J. Phys. 50, 1265–1270 (2000)

## The $q$ version: Structure

## The $q$ version: Structure

$$\mathcal{U}_q(\mathfrak{g}^{\mathbb{C}}) = \mathcal{A}_q^- \cdot \mathcal{U}_q(\mathfrak{k}^{\mathbb{C}}) \cdot \mathcal{A}_q^+.$$

$\mathcal{A}_q^{\mp}$  quadratic algebras.

## The $q$ version: Structure

$$\mathcal{U}_q(\mathfrak{g}^{\mathbb{C}}) = \mathcal{A}_q^- \cdot \mathcal{U}_q(\mathfrak{k}^{\mathbb{C}}) \cdot \mathcal{A}_q^+.$$

$\mathcal{A}_q^{\mp}$  quadratic algebras.

$$Z_{ij} Z_{ik} = q^{-1} Z_{ik} Z_{ij} \text{ if } j < k;$$

$$Z_{ij} Z_{kj} = q^{-1} Z_{kj} Z_{ij} \text{ if } i < k;$$

$$Z_{ij} Z_{st} = Z_{st} Z_{ij} \text{ if } i < s \text{ and } t < j;$$

$$Z_{ij} Z_{st} = Z_{st} Z_{ij} - (q - q^{-1}) Z_{it} Z_{sj} = \text{ if } i < s \text{ and } j < t.$$

# Quantized modules

# Quantized modules

$$M(V_{\tau'})$$

# Quantized modules

$$M(V_{\tau'}) \rightarrow M_q(V_{\tau'})$$

# Quantized modules

$M(V_{\tau'}) \rightarrow M_q(V_{\tau'})$  (straightforward)

# Quantized modules

$M(V_{\tau'}) \rightarrow M_q(V_{\tau'})$  (straightforward)

$\mathcal{P}(V_{\tau})$

# Quantized modules

$M(V_{\tau'}) \rightarrow M_q(V_{\tau'})$  (straightforward)

$\mathcal{P}(V_{\tau}) \rightarrow \mathcal{A}_q^-(V_{\tau})$

# Quantized modules

$M(V_{\tau'}) \rightarrow M_q(V_{\tau'})$  (straightforward)

$\mathcal{P}(V_{\tau}) \rightarrow \mathcal{A}_q^-(V_{\tau})$  (also straightforward)

# Quantized modules

$$M(V_{\tau'}) \rightarrow M_q(V_{\tau'}) \text{ (straightforward)}$$

$$\mathcal{P}(V_{\tau}) \rightarrow \mathcal{A}_q^-(V_{\tau}) \text{ (also straightforward)}$$

Duality “essentially” the  $q$  Killing form

## $q$ homomorphisms

$$M_q(V_{\tau'_1}) \xrightarrow{\phi_q} M_q(V_{\tau'}).$$

The resulting map on the whole space  $\mathcal{A}^- \otimes V_{\Lambda_1}$  may then be written as

## $q$ homomorphisms

$$M_q(V_{\tau'_1}) \xrightarrow{\phi_q} M_q(V_{\tau'}).$$

The resulting map on the whole space  $\mathcal{A}^- \otimes V_{\Lambda_1}$  may then be written as

$$\mathcal{A}^- \otimes V_{\Lambda_1} \ni (W_1, W_2, \dots, W_{N_1}) \xrightarrow{\phi_{\Lambda, \Lambda_1}}$$

## $q$ homomorphisms

$$M_q(V_{\tau'_1}) \xrightarrow{\phi_q} M_q(V_{\tau'}).$$

The resulting map on the whole space  $\mathcal{A}^- \otimes V_{\Lambda_1}$  may then be written as

$$\mathcal{A}^- \otimes V_{\Lambda_1} \ni (W_1, W_2, \dots, W_{N_1}) \xrightarrow{\phi_{\Lambda, \Lambda_1}}$$

$$(W_1, W_2, \dots, W_{N_1}) \begin{pmatrix} p_{11} & p_{21} & \cdots & p_{N1} \\ p_{12} & p_{22} & \cdots & p_{N2} \\ \vdots & \vdots & \ddots & \vdots \\ p_{1N_1} & p_{2N_1} & \cdots & L_{p_{NN_1}} \end{pmatrix} \in \mathcal{A}^- \otimes V_{\Lambda}.$$

## $q$ homomorphisms

We set  $\underline{\underline{\phi}}_{\Lambda, \Lambda_1} = \begin{pmatrix} p_{11} & p_{21} & \cdots & p_{N1} \\ p_{12} & p_{22} & \cdots & p_{N2} \\ \vdots & \vdots & \ddots & \vdots \\ p_{1N_1} & p_{2N_1} & \cdots & L_{p_{NN_1}} \end{pmatrix}$ . Notice that if  
 $\phi_{\Lambda, \Lambda_1} \circ \phi_{\Lambda_1, \Lambda_2} = \phi_{\Lambda, \Lambda_2}$  then  $\underline{\underline{\phi}}_{\Lambda, \Lambda_2} = \underline{\underline{\phi}}_{\Lambda_1, \Lambda_2} \cdot \underline{\underline{\phi}}_{\Lambda, \Lambda_1}$ .

## $q$ Examples (also higher dimensional)

## $q$ Examples (also higher dimensional)

$$\square_q \leftrightarrow \det_q$$

## $q$ Examples (also higher dimensional)

$$\square_q \leftrightarrow \det_q$$

$$\nabla_q = \begin{pmatrix} 0 & 0 & \frac{\partial}{\partial z_1} & \frac{\partial}{\partial z_2} \\ 0 & 0 & q \frac{\partial}{\partial z_3} & q \frac{\partial}{\partial z_4} \\ \frac{\partial}{\partial z_4} & -\frac{\partial}{\partial z_2} & 0 & 0 \\ -q^{-1} \frac{\partial}{\partial z_3} & q^{-1} \frac{\partial}{\partial z_1} & 0 & 0 \end{pmatrix}.$$

## $q$ Examples (also higher dimensional)

$$\square_q \leftrightarrow \det_q$$

$$\nabla_q = \begin{pmatrix} 0 & 0 & \frac{\partial}{\partial z_1} & \frac{\partial}{\partial z_2} \\ 0 & 0 & q \frac{\partial}{\partial z_3} & q \frac{\partial}{\partial z_4} \\ \frac{\partial}{\partial z_4} & -\frac{\partial}{\partial z_2} & 0 & 0 \\ -q^{-1} \frac{\partial}{\partial z_3} & q^{-1} \frac{\partial}{\partial z_1} & 0 & 0 \end{pmatrix}.$$

$$\nabla_q = \begin{pmatrix} z_{n,1} & z_{n,2} & z_{n,3} & \cdots & z_{n,n} \\ (-q)^{-1} z_{n-1,1} & (-q)^{-1} z_{n-1,2} & (-q)^{-1} z_{n-1,3} & \cdots & (-q)^{-1} z_{n-1,n} \\ (-q)^{-2} z_{n-2,1} & (-q)^{-2} z_{n-2,2} & (-q)^{-2} z_{n-2,3} & \cdots & (-q)^{-2} z_{n-2,n} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ (-q)^{-n+1} z_{1,1} & (-q)^{-n+1} z_{1,2} & (-q)^{-n+1} z_{1,3} & \cdots & (-q)^{-n+1} z_{1,n} \end{pmatrix}.$$

## $q$ Examples (also higher dimensional)

$$\square_q \leftrightarrow \det_q$$

$$\nabla_q = \begin{pmatrix} 0 & 0 & \frac{\partial}{\partial z_1} & \frac{\partial}{\partial z_2} \\ 0 & 0 & q \frac{\partial}{\partial z_3} & q \frac{\partial}{\partial z_4} \\ \frac{\partial}{\partial z_4} & -\frac{\partial}{\partial z_2} & 0 & 0 \\ -q^{-1} \frac{\partial}{\partial z_3} & q^{-1} \frac{\partial}{\partial z_1} & 0 & 0 \end{pmatrix}.$$

$$\nabla_q = \begin{pmatrix} z_{n,1} & z_{n,2} & z_{n,3} & \cdots & z_{n,n} \\ (-q)^{-1} z_{n-1,1} & (-q)^{-1} z_{n-1,2} & (-q)^{-1} z_{n-1,3} & \cdots & (-q)^{-1} z_{n-1,n} \\ (-q)^{-2} z_{n-2,1} & (-q)^{-2} z_{n-2,2} & (-q)^{-2} z_{n-2,3} & \cdots & (-q)^{-2} z_{n-2,n} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ (-q)^{-n+1} z_{1,1} & (-q)^{-n+1} z_{1,2} & (-q)^{-n+1} z_{1,3} & \cdots & (-q)^{-n+1} z_{1,n} \end{pmatrix}.$$

Dmitry Shklyarov, Genkai Zhang Covariant  $q$ -differential operators and unitary highest weight representations for  $U_q(su(n, n))$ . J. Math Phys **46**, 062307 (2005); <https://doi.org/10.1063/1.1927077>

## $q$ Examples (also higher dimensional)

$$\square_q \leftrightarrow \det_q$$

$$\nabla_q = \begin{pmatrix} 0 & 0 & \frac{\partial}{\partial z_1} & \frac{\partial}{\partial z_2} \\ 0 & 0 & q \frac{\partial}{\partial z_3} & q \frac{\partial}{\partial z_4} \\ \frac{\partial}{\partial z_4} & -\frac{\partial}{\partial z_2} & 0 & 0 \\ -q^{-1} \frac{\partial}{\partial z_3} & q^{-1} \frac{\partial}{\partial z_1} & 0 & 0 \end{pmatrix}.$$

$$\nabla_q = \begin{pmatrix} z_{n,1} & z_{n,2} & z_{n,3} & \cdots & z_{n,n} \\ (-q)^{-1}z_{n-1,1} & (-q)^{-1}z_{n-1,2} & (-q)^{-1}z_{n-1,3} & \cdots & (-q)^{-1}z_{n-1,n} \\ (-q)^{-2}z_{n-2,1} & (-q)^{-2}z_{n-2,2} & (-q)^{-2}z_{n-2,3} & \cdots & (-q)^{-2}z_{n-2,n} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ (-q)^{-n+1}z_{1,1} & (-q)^{-n+1}z_{1,2} & (-q)^{-n+1}z_{1,3} & \cdots & (-q)^{-n+1}z_{1,n} \end{pmatrix}.$$

Dmitry Shklyarov, Genkai Zhang Covariant  $q$ -differential operators and unitary highest weight representations for  $U_q(su(n, n))$ . J. Math Phys **46**, 062307 (2005); <https://doi.org/10.1063/1.1927077>  
 See also J: Special Classes of Homomorphisms between generalized Verma modules for  $U_q(su(n, n))$  (10 pp) 2019 J. Phys.: Conf. Ser. 1194 012055



But what is a  $q$  differential operator actually?

## But what is a $q$ differential operator actually?

$$Z^{\mathbf{a}} := Z_{11}^{a_{11}} Z_{21}^{a_{21}} \cdots Z_{n1}^{a_{n1}} Z_{12}^{a_{12}} \cdots Z_{nn}^{a_{nn}}, \quad (2)$$

$$W^{\mathbf{a}} := W_{11}^{a_{11}} W_{21}^{a_{21}} \cdots W_{n1}^{a_{n1}} W_{12}^{a_{12}} \cdots W_{nn}^{a_{nn}}. \quad (3)$$

## But what is a $q$ differential operator actually?

$$Z^{\mathbf{a}} := Z_{11}^{a_{11}} Z_{21}^{a_{21}} \cdots Z_{n1}^{a_{n1}} Z_{12}^{a_{12}} \cdots Z_{nn}^{a_{nn}}, \quad (2)$$

$$W^{\mathbf{a}} := W_{11}^{a_{11}} W_{21}^{a_{21}} \cdots W_{n1}^{a_{n1}} W_{12}^{a_{12}} \cdots W_{nn}^{a_{nn}}. \quad (3)$$

$$\begin{aligned} [[a]]_q &= 1 + q^2 + \cdots + q^{2a-2} (J), \\ \{\{a\}\}_q &= 1 + q^{-2} + \cdots + q^{-(2a-2)} (K), \text{ and} \\ [a]_q &= q^{-a+1} + \cdots + q^{a-1} (L). \end{aligned}$$

$$(Z^{\mathbf{a}}, W^{\mathbf{b}})_J = \left( \frac{-1}{q - q^{-1}} \right)^{|\mathbf{a}|} \delta_{\mathbf{a}, \mathbf{b}} [[\mathbf{a}]]_q!$$

Is left right?

Is left right?

$$\left( \widehat{\partial_{xy}^R} Z^a, W^b \right) = \left( Z^a, W^b W_{xy} \right)$$

# Is left right?

$$\begin{aligned} \left( \widehat{\partial_{xy}^R} Z^a, W^b \right) &= \left( Z^a, W^b W_{xy} \right) \\ \left( \widehat{\partial_{xy}^L} Z^a, W^b \right) &= \left( Z^a, W_{xy} W^b \right) \end{aligned}$$

# Is left right?

$$\left( \widehat{\partial_{xy}^R} Z^a, W^b \right) = \left( Z^a, W^b W_{xy} \right)$$

$$\left( \widehat{\partial_{xy}^L} Z^a, W^b \right) = \left( Z^a, W_{xy} W^b \right)$$

$$\left( \textcolor{blue}{Z_{xy}} Z^a, W^b \right) ; \quad \left( Z^a \textcolor{blue}{Z_{xy}}, W^b \right)$$

# The quantized Weyl algebra

# The quantized Weyl algebra

$$M_{ij}(Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}} \cdots Z_{nn}^{a_{nn}}) = Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}+1} \cdots Z_{nn}^{a_{nn}}$$

$$D_{ij}(Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}} \cdots Z_{nn}^{a_{nn}}) = [a_{ij}]_q Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}-1} \cdots Z_{nn}^{a_{nn}}$$

$$H_{ij}^{\pm 1}(Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}} \cdots Z_{nn}^{a_{nn}}) = q^{\pm a_{ij}} Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}+1} \cdots Z_{nn}^{a_{nn}}$$

# The quantized Weyl algebra

$$M_{ij}(Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}} \cdots Z_{nn}^{a_{nn}}) = Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}+1} \cdots Z_{nn}^{a_{nn}}$$

$$D_{ij}(Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}} \cdots Z_{nn}^{a_{nn}}) = [a_{ij}]_q Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}-1} \cdots Z_{nn}^{a_{nn}}$$

$$H_{ij}^{\pm 1}(Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}} \cdots Z_{nn}^{a_{nn}}) = q^{\pm a_{ij}} Z_{11}^{a_{11}} \cdots Z_{ij}^{a_{ij}+1} \cdots Z_{nn}^{a_{nn}}$$

$$D_{ij} M_{ij} - q M_{ij} D_{ij} = H_{ij}^{-1}$$

$$D_{ij} M_{ij}^o - q^{-1} M_{ij} D_{ij} = H_{ij}$$

$$H_{ij} D_{ij} = q^{-1} D_{ij} H_{ij}$$

$$H_{ij} M_{ij} = q M_{ij} H_{ij}.$$

# A Result

## A Result

**Proposition[arXiv:1905.04478v1 [math.QA]]** In the pairing indexed by  $L$ , the algebra of covariant differential operators is equal to the quantum Weyl algebra

Where do the (unitary!) ( $q$ -) covariant differential operators come from?

# Where do the (unitary!) ( $q$ -) covariant differential operators come from?

**Proposition.** In  $\mathcal{U}_q(su(N, N))$  there is a (unitary) scalar highest weight module  $\Pi_k$  in which all  $k \times k$  minors vanish.

# Where do the (unitary!) ( $q$ -) covariant differential operators come from?

**Proposition.** In  $\mathcal{U}_q(su(N, N))$  there is a (unitary) scalar highest weight module  $\Pi_k$  in which all  $k \times k$  minors vanish.

Consider

$$\mathcal{U}_q(su(n, n))$$

# Where do the (unitary!) ( $q$ -) covariant differential operators come from?

**Proposition.** In  $\mathcal{U}_q(su(N, N))$  there is a (unitary) scalar highest weight module  $\Pi_k$  in which all  $k \times k$  minors vanish.

Consider

$$\mathcal{U}_q(su(n, n)) \subset \mathcal{U}_q(su(n, 2n))$$

# Where do the (unitary!) ( $q$ -) covariant differential operators come from?

**Proposition.** In  $\mathcal{U}_q(su(N, N))$  there is a (unitary) scalar highest weight module  $\Pi_k$  in which all  $k \times k$  minors vanish.

Consider

$$\mathcal{U}_q(su(n, n)) \subset \mathcal{U}_q(su(n, 2n)) \subset \mathcal{U}_q(su(N, N))$$

# Where do the (unitary!) ( $q$ -) covariant differential operators come from?

**Proposition.** In  $\mathcal{U}_q(su(N, N))$  there is a (unitary) scalar highest weight module  $\Pi_k$  in which all  $k \times k$  minors vanish.

Consider

$$\mathcal{U}_q(su(n, n)) \subset \mathcal{U}_q(su(n, 2n)) \subset \mathcal{U}_q(su(N, N))$$

Where do the (unitary!) ( $q$ -) covariant differential operators come from?

## Where do the (unitary!) ( $q$ -) covariant differential operators come from?

The (auxiliary) variables  $W_{a,b}$ ;  $a \leq n$  and  $n+1 \leq b \leq 2n$  can be used to define highest weight modules for  $\mathcal{U}_q(su(n,n))$ .

## Where do the (unitary!) ( $q$ -) covariant differential operators come from?

The (auxiliary) variables  $W_{a,b}$ ;  $a \leq n$  and  $n+1 \leq b \leq 2n$  can be used to define highest weight modules for  $\mathcal{U}_q(su(n,n))$ .

Thus we have a representation with highest weight vector  $W_{1,n+1}^R$  ( $R = 1, 2, \dots$ ). Consider

## Where do the (unitary!) ( $q$ -) covariant differential operators come from?

The (auxiliary) variables  $W_{a,b}$ ;  $a \leq n$  and  $n+1 \leq b \leq 2n$  can be used to define highest weight modules for  $\mathcal{U}_q(su(n,n))$ .

Thus we have a representation with highest weight vector  $W_{1,n+1}^R$  ( $R = 1, 2, \dots$ ). Consider (clearly a h.w. vector!)

## Where do the (unitary!) ( $q$ -) covariant differential operators come from?

The (auxiliary) variables  $W_{a,b}$ ;  $a \leq n$  and  $n+1 \leq b \leq 2n$  can be used to define highest weight modules for  $\mathcal{U}_q(su(n,n))$ .

Thus we have a representation with highest weight vector  $W_{1,n+1}^R$  ( $R = 1, 2, \dots$ ). Consider (clearly a h.w. vector!)

$$(W_{11} W_{2,n+1} - q W_{21} W_{1,n+1}) W_{1,n+1}^{R-1} =$$

$$W_{11} W_{1,n+1}^{R-1} W_{2,n+1} - q W_{2,1} W_{1,n+1}^R.$$

## Where do the (unitary!) ( $q$ -) covariant differential operators come from?

The (auxiliary) variables  $W_{a,b}$ ;  $a \leq n$  and  $n+1 \leq b \leq 2n$  can be used to define highest weight modules for  $\mathcal{U}_q(\mathfrak{su}(n,n))$ .

Thus we have a representation with highest weight vector  $W_{1,n+1}^R$  ( $R = 1, 2, \dots$ ). Consider (clearly a h.w. vector!)

$$(W_{11} W_{2,n+1} - q W_{21} W_{1,n+1}) W_{1,n+1}^{R-1} =$$

$$W_{11} W_{1,n+1}^{R-1} W_{2,n+1} - q W_{2,1} W_{1,n+1}^R.$$

Both vectors  $W_{1,n+1}^{R-1} W_{2,n+1}$  and  $W_{1,n+1}^R$  are in the h.w. module generated by  $W_{1,n+1}^R$ , but the vector vanishes, i.e. there is a homomorphism!!! into  $M_q(W_{1,n+1}^R)$ .

# The support of a Representation

# The support of a Representation

**Definition** The support of  $M(V_\Lambda)$ ,  $\Lambda = (\Lambda_0, \lambda)$ , is the smallest  $r > 0$  for which there exists a  $M(V_{\Lambda_1})$  and a homomorphism  $\psi : M(V_{\Lambda_1}) \rightarrow M(V_\Lambda)$  such that

$$\text{Minor}_r \otimes (V_\Lambda) \subseteq \phi(M(V_{\Lambda_1}))$$

Here,  $\text{Minor}_r$  is the ideal generated by  $r \times r$  minors.

# The support of a Representation

**Definition** The support of  $M(V_\Lambda)$ ,  $\Lambda = (\Lambda_0, \lambda)$ , is the smallest  $r > 0$  for which there exists a  $M(V_{\Lambda_1})$  and a homomorphism  $\psi : M(V_{\Lambda_1}) \rightarrow M(V_\Lambda)$  such that

$$\text{Minor}_r \otimes (V_\Lambda) \subseteq \phi(M(V_{\Lambda_1}))$$

Here,  $\text{Minor}_r$  is the ideal generated by  $r \times r$  minors.(Look at dual...) But left or right ideal??

## The support of a Representation

**Definition** The support of  $M(V_\Lambda)$ ,  $\Lambda = (\Lambda_0, \lambda)$ , is the smallest  $r > 0$  for which there exists a  $M(V_{\Lambda_1})$  and a homomorphism  $\psi : M(V_{\Lambda_1}) \rightarrow M(V_\Lambda)$  such that

$$\text{Minor}_r \otimes (V_\Lambda) \subseteq \phi(M(V_{\Lambda_1}))$$

Here,  $\text{Minor}_r$  is the ideal generated by  $r \times r$  minors.(Look at dual...) But left or right ideal??

**Corollary** The left ideal  $\mathcal{I}_q^+(n-1)_L$  in  $\mathcal{A}_q^+$  generated by the set of all  $(n-1) \times (n-1)$  minors is equal to the right ideal  $\mathcal{I}_q^+(n-1)_R$  generated by the same.

# The support of a Representation

**Definition** The support of  $M(V_\Lambda)$ ,  $\Lambda = (\Lambda_0, \lambda)$ , is the smallest  $r > 0$  for which there exists a  $M(V_{\Lambda_1})$  and a homomorphism  $\psi : M(V_{\Lambda_1}) \rightarrow M(V_\Lambda)$  such that

$$\text{Minor}_r \otimes (V_\Lambda) \subseteq \phi(M(V_{\Lambda_1}))$$

Here,  $\text{Minor}_r$  is the ideal generated by  $r \times r$  minors. (Look at dual...) But left or right ideal??

**Corollary** The left ideal  $\mathcal{I}_q^+(n-1)_L$  in  $\mathcal{A}_q^+$  generated by the set of all  $(n-1) \times (n-1)$  minors is equal to the right ideal  $\mathcal{I}_q^+(n-1)_R$  generated by the same.

**Theorem** We get all **unitary** representations of support  $r = 1, \dots, n$  this way.

# The Hayashi-Weyl Algebra

# The Hayashi-Weyl Algebra

T. Hayashi, q-analogues of Clifford and Weyl Algebras–Spinor and Oscillator Representations of Quantum Enveloping Algebras,  
Comm. Math. Phys. **127**, Number 1 (1990), 129-144.

# The Hayashi-Weyl Algebra

T. Hayashi, q-analogues of Clifford and Weyl Algebras–Spinor and Oscillator Representations of Quantum Enveloping Algebras, Comm. Math. Phys. **127**, Number 1 (1990), 129-144.

We start with a quantum analog for the Weyl algebra

$\mathcal{A}_q = \mathcal{A}_q^-(N)$  introduced by Hayashi. Let  $q$  be a non-zero complex number such that  $q^4 \neq 1$ . Then the algebra is defined as an associative unital algebra with generators  $\psi_i, \psi_i^*, \omega_i^{\pm 1}$  and relations

$$\omega_i \omega_j = \omega_j \omega_i, \quad \omega_i \omega_i^{-1} = \omega_i^{-1} \omega_i = 1,$$

$$\psi_i \psi_j = \psi_j \psi_i, \quad \psi_i^* \psi_j^* = \psi_j^* \psi_i^*,$$

$$\psi_i \psi_j = \psi_j \psi_i \quad i \neq j,$$

$$\omega_i \psi_j \omega_i^{-1} = q^{-\delta_{ij}} \psi_j, \quad \omega_i \psi_j^* \omega_i^{-1} = q^{\delta_{ij}} \psi_j^*,$$

$$\psi_i \psi_i^* - q^2 \psi_i^* \psi_i = \omega_i^{-2}, \quad \psi_i \psi_i^* - q^{-2} \psi_i^* \psi_i = \omega_i^2.$$

The last pair of relations is equivalent to the following

$$\psi_i \psi_i^* = \frac{(q \omega_i)^2 - (q \omega_i)^{-2}}{q^2 - q^{-2}}, \quad \psi_i^* \psi_i = \frac{\omega_i^2 - \omega_i^{-2}}{q^2 - q^{-2}}.$$

# The Hayashi-Weyl Algebra

# The Hayashi-Weyl Algebra

If

$$H = \omega, D = \psi, M = \psi^*, \text{ and } q \rightarrow q^2$$

## The Hayashi-Weyl Algebra

If

$$H = \omega, D = \psi, M = \psi^*, \text{ and } q \rightarrow q^2$$

This is just the quantized Weyl algebra. The introduced algebra has a natural representation in infinite-dimensional vector space  $V^- = \{v(\mathbf{m}) | \mathbf{m} \in \mathbb{Z}_+^N\}$ . Namely,

$$\omega_i v(\mathbf{m}) = q^{m_i} v(\mathbf{m}), \quad \psi_i v(\mathbf{m}) = [m_i]_{q^2} v(\mathbf{m} - \mathbf{e}_i), \quad \psi_i^* v(\mathbf{m}) = v(\mathbf{m} + \mathbf{e}_i),$$

where  $\mathbf{m} = \{m_1, \dots, m_N\}$  and  $[m]_q = \frac{q^m - q^{-m}}{q - q^{-1}}$ .

Let us consider  $N = 2n$  and obtain explicit formulas for embedding of  $U_q \mathfrak{sl}_{2n}$  into  $\mathcal{A}_q$ . Formulas for simple roots remain the same

$$E_n = i\psi_n \psi_{n+1},$$

$$F_n = i\psi_n^* \psi_{n+1}^*,$$

$$E_j = \psi_j \psi_{j+1}^*,$$

$$F_j = \psi_j^* \psi_{j+1}, \quad j < n,$$

$$E_j = \psi_j^* \psi_{j+1},$$

$$F_j = \psi_j \psi_{j+1}^*, \quad j > n,$$

# The Hayashi-Weyl algebra and $\mathcal{U}_q(su(n, n))$

$$K_n = q^{-2} \omega_n^{-2} \omega_{n+1}^{-2}; \quad E_n F_n - F_n E_n = \frac{K_n - K_n^{-1}}{q^2 - q^{-2}}$$

$$K_j = \omega_{j+1}^2 \omega_j^{-2} \text{ for } j < n$$

$$K_j = \omega_j^2 \omega_{j+1}^{-2} \text{ for } j > n$$

# The Hayashi-Weyl algebra and $\mathcal{U}_q(su(n, n))$

$$K_n = q^{-2} \omega_n^{-2} \omega_{n+1}^{-2}; \quad E_n F_n - F_n E_n = \frac{K_n - K_n^{-1}}{q^2 - q^{-2}}$$

$$K_j = \omega_{j+1}^2 \omega_j^{-2} \text{ for } j < n$$

$$K_j = \omega_j^2 \omega_{j+1}^{-2} \text{ for } j > n$$

$$\begin{pmatrix} (-q)^{-2(n-1)} \prod_2^n \omega_i^{-2} \psi_1 \psi_{n+1} & (-q)^{-2n} \prod_2^{n+1} \omega_i^{-2} \psi_1 \psi_{n+2} & \cdots \\ (-q)^{-2(n-2)} \prod_3^n \omega_i^{-2} \psi_2 \psi_{n+1} & (-q)^{-2(n-1)} \prod_3^{n+1} \omega_i^{-2} \psi_2 \psi_{n+2} & \cdots \\ \vdots & \vdots & \cdots \\ \psi_n \psi_{n+1} & (-q)^{-2} \omega_{n+1}^{-2} \psi_n \psi_{n+2} & \cdots \end{pmatrix} \quad (4)$$

# The Hayashi-Weyl algebra and $\mathcal{U}_q(su(n, n))$

$$Z_{ij} = (i\psi_{n+1-i}\psi_{n+j})(-q^{-2})^{i+j-2} \prod_{x=1}^{i-1} \prod_{y=1}^{j-1} \omega_{n+1-x}^{-2} \omega_{n+j}^{-2}.$$

$$Z_{11}Z_{22} - Z_{22}Z_{22} = (q^{-2} - q^2)Z_{21}Z_{12}; Z_{11}Z_{21} = q^{-2}Z_{21}Z_{11}, ETC$$

# The Hayashi-Weyl algebra and $\mathcal{U}_q(su(n, n))$

$$Z_{ij} = (i\psi_{n+1-i}\psi_{n+j})(-q^{-2})^{i+j-2} \prod_{x=1}^{i-1} \prod_{y=1}^{j-1} \omega_{n+1-x}^{-2} \omega_{n+j}^{-2}.$$

$$Z_{11}Z_{22} - Z_{22}Z_{22} = (q^{-2} - q^2)Z_{21}Z_{12}; Z_{11}Z_{21} = q^{-2}Z_{21}Z_{11}, ETC$$

**Proposition.** All  $2 \times 2$   $q$  minors in the variables  $Z_{k,\ell}$  vanish.

# The Hayashi-Weyl algebra and $\mathcal{U}_q(su(n, n))$

$$Z_{ij} = (i\psi_{n+1-i}\psi_{n+j})(-q^{-2})^{i+j-2} \prod_{x=1}^{i-1} \prod_{y=1}^{j-1} \omega_{n+1-x}^{-2} \omega_{n+j}^{-2}.$$

$$Z_{11}Z_{22} - Z_{22}Z_{22} = (q^{-2} - q^2)Z_{21}Z_{12}; Z_{11}Z_{21} = q^{-2}Z_{21}Z_{11}, \text{ ETC}$$

**Proposition.** All  $2 \times 2$   $q$  minors in the variables  $Z_{k,\ell}$  vanish.

D. Shklyarov, S. Sinel'shchikov, A. Stolin, and L. Vaksman, ON A  
q-ANALOGUE OF THE PENROSE TRANSFORM. In:  
LECTURES ON q-ANALOGUES OF CARTAN DOMAINS AND  
ASSOCIATED HARISH-CHANDRA MODULES L. Vaksman (Ed.)  
arXiv:math/0109198

## $\mathcal{U}_q(su(n, n))$ and higher order minors

More generally, for  $k = 1, 2, \dots$ :

## $\mathcal{U}_q(su(n, n))$ and higher order minors

More generally, for  $k = 1, 2, \dots$ :

$$\begin{aligned} Z_{1,1}^{(k)} &= \psi_n^{(1)} \psi_{n+1}^{(1)} + A_n^{(1)} A_{n+1}^{(1)} \psi_n^{(2)} \psi_{n+1}^{(2)} + A_n^{(1)} A_{n+1}^{(1)} A_n^{(2)} A_{n+1}^{(2)} \psi_n^{(3)} \psi_{n+1}^{(3)} \\ &\quad + \cdots + A_n^{(1)} A_{n+1}^{(1)} \cdots A_n^{(k-1)} A_{n+1}^{(k-1)} \psi_n^{(k)} \psi_{n+1}^{(k)} \end{aligned}$$

## $\mathcal{U}_q(su(n, n))$ and higher order minors

More generally, for  $k = 1, 2, \dots$ :

$$\begin{aligned} Z_{1,1}^{(k)} &= \psi_n^{(1)} \psi_{n+1}^{(1)} + A_n^{(1)} A_{n+1}^{(1)} \psi_n^{(2)} \psi_{n+1}^{(2)} + A_n^{(1)} A_{n+1}^{(1)} A_n^{(2)} A_{n+1}^{(2)} \psi_n^{(3)} \psi_{n+1}^{(3)} \\ &\quad + \cdots + A_n^{(1)} A_{n+1}^{(1)} \cdots A_n^{(k-1)} A_{n+1}^{(k-1)} \psi_n^{(k)} \psi_{n+1}^{(k)} \\ \mu_{n-1}^{(k)} &= \psi_{n-1}^{(1)} (\psi_n^{(1)})^* + q^{-2} (A_n^{(1)})^{-1} A_{n-1}^{(1)} \psi_{n-1}^{(2)} (\psi_n^{(2)})^* + \\ &\quad q^{-4} (A_n^{(1)})^{-1} A_{n-1}^{(1)} (A_n^{(2)})^{-1} A_n^{(2)} \psi_n^{(3)} \psi_{n+1}^{(3)} \\ &\quad \cdots + q^{-2(k-1)} (A_n^{(1)})^{-1} A_{n-1}^{(1)} \cdots (A_n^{(k-1)})^{-1} A_{n-1}^{(k-1)} \psi_{n-1}^{(k)} (\psi_n^{(k)})^* \end{aligned}$$

## $\mathcal{U}_q(su(n, n))$ and higher order minors

More generally, for  $k = 1, 2, \dots$ :

$$Z_{1,1}^{(k)} = \psi_n^{(1)} \psi_{n+1}^{(1)} + A_n^{(1)} A_{n+1}^{(1)} \psi_n^{(2)} \psi_{n+1}^{(2)} + A_n^{(1)} A_{n+1}^{(1)} A_n^{(2)} A_{n+1}^{(2)} \psi_n^{(3)} \psi_{n+1}^{(3)}$$
$$+ \cdots + A_n^{(1)} A_{n+1}^{(1)} \cdots A_n^{(k-1)} A_{n+1}^{(k-1)} \psi_n^{(k)} \psi_{n+1}^{(k)}$$

$$\mu_{n-1}^{(k)} = \psi_{n-1}^{(1)} (\psi_n^{(1)})^* + q^{-2} (A_n^{(1)})^{-1} A_{n-1}^{(1)} \psi_{n-1}^{(2)} (\psi_n^{(2)})^* +$$
$$q^{-4} (A_n^{(1)})^{-1} A_{n-1}^{(1)} (A_n^{(2)})^{-1} A_n^{(2)} \psi_n^{(3)} \psi_{n+1}^{(3)}$$
$$\cdots + q^{-2(k-1)} (A_n^{(1)})^{-1} A_{n-1}^{(1)} \cdots (A_n^{(k-1)})^{-1} A_{n-1}^{(k-1)} \psi_{n-1}^{(k)} (\psi_n^{(k)})^*$$

$$\nu_{n-1}^{(k)} = \text{similarly}$$

## $\mathcal{U}_q(su(n, n))$ and higher order minors

More generally, for  $k = 1, 2, \dots$ :

$$\begin{aligned} Z_{1,1}^{(k)} &= \psi_n^{(1)} \psi_{n+1}^{(1)} + A_n^{(1)} A_{n+1}^{(1)} \psi_n^{(2)} \psi_{n+1}^{(2)} + A_n^{(1)} A_{n+1}^{(1)} A_n^{(2)} A_{n+1}^{(2)} \psi_n^{(3)} \psi_{n+1}^{(3)} \\ &\quad + \cdots + A_n^{(1)} A_{n+1}^{(1)} \cdots A_n^{(k-1)} A_{n+1}^{(k-1)} \psi_n^{(k)} \psi_{n+1}^{(k)} \end{aligned}$$

$$\begin{aligned} \mu_{n-1}^{(k)} &= \psi_{n-1}^{(1)} (\psi_n^{(1)})^* + q^{-2} (A_n^{(1)})^{-1} A_{n-1}^{(1)} \psi_{n-1}^{(2)} (\psi_n^{(2)})^* + \\ &\quad q^{-4} (A_n^{(1)})^{-1} A_{n-1}^{(1)} (A_n^{(2)})^{-1} A_n^{(2)} \psi_n^{(3)} \psi_{n+1}^{(3)} \\ &\quad \cdots + q^{-2(k-1)} (A_n^{(1)})^{-1} A_{n-1}^{(1)} \cdots (A_n^{(k-1)})^{-1} A_{n-1}^{(k-1)} \psi_{n-1}^{(k)} (\psi_n^{(k)})^* \end{aligned}$$

$$\nu_{n-1}^{(k)} = \text{similarly}$$

$$\begin{aligned} Z_{s,\ell}^{(k)} &= \psi_{n+1-s}^{(1)} \psi_{n+\ell}^{(1)} + A_{n+1-s}^{(1)} A_{n+\ell}^{(1)} \psi_{n+1-s}^{(2)} \psi_{n+\ell}^{(2)} \\ &\quad + A_{n+1-s}^{(1)} A_{n+\ell}^{(1)} A_{n+1-s}^{(2)} A_{n+\ell}^{(2)} \psi_{n+1-s}^{(3)} \psi_{n+\ell}^{(3)} \\ &\quad + \cdots + A_{n+1-s}^{(1)} A_{n+\ell}^{(1)} \cdots A_{n+1-s}^{(k-1)} A_{n+\ell}^{(k-1)} \psi_{n+1-s}^{(k)} \psi_{n+\ell}^{(k)} \end{aligned}$$

## $\mathcal{U}_q(su(n, n))$ and higher order minors

More generally, for  $k = 1, 2, \dots$ :

$$\begin{aligned} Z_{1,1}^{(k)} &= \psi_n^{(1)} \psi_{n+1}^{(1)} + A_n^{(1)} A_{n+1}^{(1)} \psi_n^{(2)} \psi_{n+1}^{(2)} + A_n^{(1)} A_{n+1}^{(1)} A_n^{(2)} A_{n+1}^{(2)} \psi_n^{(3)} \psi_{n+1}^{(3)} \\ &\quad + \cdots + A_n^{(1)} A_{n+1}^{(1)} \cdots A_n^{(k-1)} A_{n+1}^{(k-1)} \psi_n^{(k)} \psi_{n+1}^{(k)} \end{aligned}$$

$$\begin{aligned} \mu_{n-1}^{(k)} &= \psi_{n-1}^{(1)} (\psi_n^{(1)})^* + q^{-2} (A_n^{(1)})^{-1} A_{n-1}^{(1)} \psi_{n-1}^{(2)} (\psi_n^{(2)})^* + \\ &\quad q^{-4} (A_n^{(1)})^{-1} A_{n-1}^{(1)} (A_n^{(2)})^{-1} A_n^{(2)} \psi_n^{(3)} \psi_{n+1}^{(3)} \\ &\quad \cdots + q^{-2(k-1)} (A_n^{(1)})^{-1} A_{n-1}^{(1)} \cdots (A_n^{(k-1)})^{-1} A_{n-1}^{(k-1)} \psi_{n-1}^{(k)} (\psi_n^{(k)})^* \end{aligned}$$

$$\nu_{n-1}^{(k)} = \text{similarly}$$

$$\begin{aligned} Z_{s,\ell}^{(k)} &= \psi_{n+1-s}^{(1)} \psi_{n+\ell}^{(1)} + A_{n+1-s}^{(1)} A_{n+\ell}^{(1)} \psi_{n+1-s}^{(2)} \psi_{n+\ell}^{(2)} \\ &\quad + A_{n+1-s}^{(1)} A_{n+\ell}^{(1)} A_{n+1-s}^{(2)} A_{n+\ell}^{(2)} \psi_{n+1-s}^{(3)} \psi_{n+\ell}^{(3)} \\ &\quad + \cdots + A_{n+1-s}^{(1)} A_{n+\ell}^{(1)} \cdots A_{n+1-s}^{(k-1)} A_{n+\ell}^{(k-1)} \psi_{n+1-s}^{(k)} \psi_{n+\ell}^{(k)} \\ &\quad + p_{s,\ell}(\psi, \psi^*)!!! \end{aligned}$$

## $\mathcal{U}_q(su(n, n))$ and higher order minors

More generally, for  $k = 1, 2, \dots$ :

$$\begin{aligned} Z_{1,1}^{(k)} &= \psi_n^{(1)} \psi_{n+1}^{(1)} + A_n^{(1)} A_{n+1}^{(1)} \psi_n^{(2)} \psi_{n+1}^{(2)} + A_n^{(1)} A_{n+1}^{(1)} A_n^{(2)} A_{n+1}^{(2)} \psi_n^{(3)} \psi_{n+1}^{(3)} \\ &\quad + \cdots + A_n^{(1)} A_{n+1}^{(1)} \cdots A_n^{(k-1)} A_{n+1}^{(k-1)} \psi_n^{(k)} \psi_{n+1}^{(k)} \end{aligned}$$

$$\begin{aligned} \mu_{n-1}^{(k)} &= \psi_{n-1}^{(1)} (\psi_n^{(1)})^* + q^{-2} (A_n^{(1)})^{-1} A_{n-1}^{(1)} \psi_{n-1}^{(2)} (\psi_n^{(2)})^* + \\ &\quad q^{-4} (A_n^{(1)})^{-1} A_{n-1}^{(1)} (A_n^{(2)})^{-1} A_n^{(2)} \psi_n^{(3)} \psi_{n+1}^{(3)} \\ &\quad \cdots + q^{-2(k-1)} (A_n^{(1)})^{-1} A_{n-1}^{(1)} \cdots (A_n^{(k-1)})^{-1} A_{n-1}^{(k-1)} \psi_{n-1}^{(k)} (\psi_n^{(k)})^* \end{aligned}$$

$$\nu_{n-1}^{(k)} = \text{similarly}$$

$$\begin{aligned} Z_{s,\ell}^{(k)} &= \psi_{n+1-s}^{(1)} \psi_{n+\ell}^{(1)} + A_{n+1-s}^{(1)} A_{n+\ell}^{(1)} \psi_{n+1-s}^{(2)} \psi_{n+\ell}^{(2)} \\ &\quad + A_{n+1-s}^{(1)} A_{n+\ell}^{(1)} A_{n+1-s}^{(2)} A_{n+\ell}^{(2)} \psi_{n+1-s}^{(3)} \psi_{n+\ell}^{(3)} \\ &\quad + \cdots + A_{n+1-s}^{(1)} A_{n+\ell}^{(1)} \cdots A_{n+1-s}^{(k-1)} A_{n+\ell}^{(k-1)} \psi_{n+1-s}^{(k)} \psi_{n+\ell}^{(k)} \\ &\quad + p_{s,\ell}(\psi, \psi^*)!!! \end{aligned}$$

# $\mathcal{U}_q(su(n, n))$ and higher order minors

# $\mathcal{U}_q(su(n, n))$ and higher order minors

**Proposition**  $\mathcal{W}eyl_q(1)$  is a UFD.

# $\mathcal{U}_q(su(n, n))$ and higher order minors

**Proposition**  $\mathcal{W}eyl_q(1)$  is a UFD.

**Theorem** The ideal of quantum  $(k + 1) \times (k + 1)$  minor is prime.

# $\mathcal{U}_q(su(n, n))$ and higher order minors

**Proposition**  $\mathcal{W}eyl_q(1)$  is a UFD.

**Theorem** The ideal of quantum  $(k + 1) \times (k + 1)$  minor is prime.  
(Since we are in the  $k$ th tensor product of rank 1.)

# $\mathcal{U}_q(su(n, n))$ and higher order minors

**Proposition**  $\mathcal{W}eyl_q(1)$  is a UFD.

**Theorem** The ideal of quantum  $(k + 1) \times (k + 1)$  minor is prime.  
(Since we are in the  $k$ th tensor product of rank 1.)

S. Launois, T. H. Lenagan, and L. Rigal, Quantum Unique Factorisation Domains, J. London Math. Soc. **74**, 321-340 (2006)

# The End

# The End

Thank you