

Numerical Lecture Notes

Thm

Let X and \hat{X} be solutions to SDE with initial values X_0 and \hat{X}_0 . Then for $C = C(L, T)$

$$E \left(\max_{0 \leq t \leq T} |\hat{X}(t) - X(t)|^2 \right) \leq C E(|\hat{X}_0 - X_0|^2)$$

Proof

$$\phi(t) = E \left(\max_{0 \leq s \leq t} |\hat{X}(s) - X(s)|^2 \right)$$

$$= E \left(\max_{0 \leq s \leq t} \left| \hat{X}_0 - X_0 + \int_0^s [\mu(\hat{X}(r), r) - \mu(X(r), r)] dr + \int_0^s [\sigma(\hat{X}(r), r) - \sigma(X(r), r)] dB(r) \right|^2 \right)$$

$(x+y+z)^2 \leq 3x^2 + 3y^2 + 3z^2$

$$\leq 3E(|\hat{X}_0 - X_0|^2) + 3E \left(\max_{0 \leq s \leq t} \left| \int_0^s (\hat{\mu} - \mu) dr \right|^2 \right) + 3E \left(\max_{0 \leq s \leq t} \left| \int_0^s (\hat{\sigma} - \sigma) dB \right|^2 \right)$$

Cauchy-Schwarz
Doob

$$\leq 3E(|\hat{X}_0 - X_0|^2) + 3E \left(\max_{0 \leq s \leq t} \int_0^s 1^2 dt \int_0^s (\hat{\mu} - \mu)^2 dr \right) + 12 E \left(\int_0^t (\hat{\sigma} - \sigma)^2 dB \right)$$

isometry

$$\leq 3E(|\hat{X}_0 - X_0|^2) + 3E \left(\int_0^t (\hat{\mu} - \mu)^2 dr \right) + 12 E \left(\int_0^t (\hat{\sigma} - \sigma)^2 dr \right)$$

Lipschitz

$$\leq 3E(|\hat{X}_0 - X_0|^2) + 3TL^2 \int_0^t E((\hat{X}(r) - X(r))^2) dr + 12L^2 \int_0^t E((\hat{X}(r) - X(r))^2) dr$$

$$\leq 3E(|\hat{X}_0 - X_0|^2) + (3TL^2 + 12L^2) \int_0^t \phi(r) dr \quad \xRightarrow{\text{Grönwall}}$$

$$\phi(t) \leq 3E(|\hat{X}_0 - X_0|^2) e^{(3TL^2 + 12L^2)t}$$

for $t \in [0, T]$.
finish by taking $t = T$.

Thm

For $C = C(L, T)$ we have

$$E(|X(t) - X(s)|^2) \leq C(1 + E(X_0^2)) |t - s| \quad \text{for } s, t \in [0, T]$$

Proof $E(|X(t) - X(s)|^2) = E\left(\left|\int_s^t \mu(X(r), r) dr + \int_s^t \sigma(X(r), r) dB(r)\right|^2\right)$

$[(x+y)^2 \leq 2x^2 + 2y^2] \leq 2E\left(\left(\int_s^t \mu dr\right)^2\right) + 2E\left(\left(\int_s^t \sigma dB\right)^2\right)$

$[\text{Cauchy-Schwarz isometry}] \leq 2 \int_s^t |\mu|^2 dr \int_s^t E(\mu^2) dr + 2E\left(\int_s^t \sigma^2 dr\right)$

$[\text{linear growth}] \leq 4(t-s) \int_s^t L^2 E(1 + |X(r)|^2) dr + 4 \int_s^t L^2 E(1 + |X(r)|^2) dr$

$[\text{Theorem 1 Stig's notes}] \leq 4(T+1)L^2(1 + C(1 + E(X_0^2)))(t-s) \neq$

Notation $0 = t_0 < t_1 < \dots < t_N = T, h_n = t_{n+1} - t_n, \Delta B_n = B(t_{n+1}) - B(t_n)$

$X(t_{n+1}) = X(t_n) + \int_{t_n}^{t_{n+1}} \mu(X(s), s) ds + \int_{t_n}^{t_{n+1}} \sigma(X(s), s) dB(s), X(0) = X_0$

$Y_{n+1} = Y_n + \mu(Y_n, t_n) \int_{t_n}^{t_{n+1}} ds + \sigma(Y_n, t_n) \int_{t_n}^{t_{n+1}} dB(s)$
 $= Y_n + \mu(Y_n, t_n) h_n + \sigma(Y_n, t_n) \Delta B_n, Y_0 \approx X_0$

$Y(t) = Y_0 + \int_0^t \bar{\mu}(s) ds + \int_0^t \bar{\sigma}(s) dB(s)$ satisfies $Y(t_n) = Y_n$ for

$\bar{\mu}(s) = \mu(Y_n, t_n)$ and $\bar{\sigma}(s) = \sigma(Y_n, t_n)$ for $s \in [t_n, t_{n+1}]$.

Thm

IF $E(|Y_0 - X_0|^2) \leq kh$ and $E(X_0^2) \leq M$ then

$E\left(\max_{0 \leq t \leq T} |Y(t) - X(t)|^2\right) \leq Ch$

Proof $\phi(t) = E(\max_{0 \leq s \leq t} |Y(s) - X(s)|^2)$

$$\left[\begin{array}{l} (x+y+z)^2 \\ \leq 3(x^2+y^2+z^2) \end{array} \right] \leq 3E((Y_0 - X_0)^2) + 3E\left(\max_{0 \leq s \leq t} \left(\int_0^s (\bar{u}(r) - u(r)) dr\right)^2\right) + 3E\left(\max_{0 \leq s \leq t} \left(\int_0^s (\bar{\sigma}(r) - \sigma(r)) dB(r)\right)^2\right)$$

$$\left[\begin{array}{l} \text{Cauchy-Schwarz} \\ \text{Doob} \end{array} \right] \leq 3Kh + 3 \int_0^t L^2 dr \int_0^s (\bar{u}(r) - u(r))^2 dr + 12 E\left(\int_0^t (\bar{\sigma}(r) - \sigma(r))^2 dB(r)\right)^2$$

$$\left[\text{isometry} \right] \leq 3Kh + 3T \int_0^t E\left(\frac{(\bar{u}(r) - u(r))^2}{dr}\right) + 12 \int_0^t E\left(\frac{(\bar{\sigma}(r) - \sigma(r))^2}{dr}\right) = (*)$$

$$E(\bar{u}(s) - u(s))^2 \leq 2E(|u(Y(t_n), t_n) - u(X(s), t_n)|^2) + 2E(|u(X(s), t_n) - u(X(t), s)|^2)$$

$$\leq 2L^2 E(|Y(t_n) - X(s)|^2) + 2L^2 |t_n - s|^2 \quad (L_n) = L$$

$$\leq 4L^2 E(|Y(t_n) - X(t_n)|^2) + 4L^2 E(|X(t_n) - X(s)|^2) + 2L^2 h_n^2$$

$$\leq 4L^2 E\left(\max_{0 \leq s \leq t} |Y(s) - X(s)|^2\right) + \underbrace{4L^2 C |t_n - s| + 2L^2 h_n^2}_{\leq Ch \leq Ch} \quad \text{for } s \in [t_n, t_{n+1}]$$

We get same with other constants for $E((\bar{\sigma}(s) - \sigma(s))^2)$ and so we have

$$\phi(t) = Ch + \int_0^t \phi(s) ds \quad \text{for } t \in [0, T] \text{ so that}$$

$$\phi(t) = Ch e^{Bt} \quad \text{for } t \in [0, T]. \quad \#$$

Ito-Taylor expansion

$$dX(t) = \mu(X(t), t) dt + \sigma(X(t), t) dB(t)$$

$$X(t) = X(t_0) + \int_{t_0}^t \mu(X(s), s) ds + \int_{t_0}^t \sigma(X(s), s) dB(s) \quad (*)$$

$$f(X(t), s) = f(X(t_0), t_0) + \int_{t_0}^s f'_x(X(r), r) [\mu(X(r), r) dr + \sigma(X(r), r) dB(r)] \\ + \int_{t_0}^s \frac{1}{2} f''_{xx}(X(r), r) \sigma(X(r), r)^2 dr + \int_{t_0}^s f'_t(X(r), r) dr \quad (**)$$

Using (**) with $f(x, t) = \mu(x, t)$ and $f(x, t) = \sigma(x, t)$ in (*) gives

$$\underline{X(t)} = \underline{X(t_0)} + \underline{(t-t_0)\mu(X(t_0), t_0)} + \int_{t_0}^t \int_{t_0}^s \underline{\mu'_x(X(r), r)\mu(X(r), r)} dr ds \\ + \int_{t_0}^t \int_{t_0}^s \underline{\mu'_x(X(r), r)\sigma(X(r), r)} dB(r) ds + \int_{t_0}^t \int_{t_0}^s \frac{1}{2} \underline{\mu''_{xx}(X(r), r)\sigma(X(r), r)^2} dr ds \\ + \int_{t_0}^t \int_{t_0}^s \underline{\mu'_t(X(r), r)} dr ds + \underline{(B(t)-B(t_0))\sigma(X(t_0), t_0)} \\ + \int_{t_0}^t \int_{t_0}^s \underline{\sigma'_x(X(r), r)\mu(X(r), r)} dr dB(s) + \int_{t_0}^t \int_{t_0}^s \underline{\sigma'_x(X(r), r)\sigma(X(r), r)} dB(r) dB(s) \\ + \int_{t_0}^t \int_{t_0}^s \frac{1}{2} \underline{\sigma''_{xx}(X(r), r)\sigma(X(r), r)^2} dr dB(s) + \int_{t_0}^t \int_{t_0}^s \underline{\sigma'_t(X(r), r)} dr dB(s)$$

The underlined — parts of above make up the motivation for the Euler (Maruyama) method for numerical approximative solution of the SDE. By adding the mm part similarly approximated by

$$(\sigma'_x \sigma)(X(t_0), t_0) \int_{t_0}^t \int_{t_0}^s dB(r) dB(s) = (\sigma'_x \sigma)(X(t_0), t_0) \int_{t_0}^t (B(s) - B(t_0)) dB(s) \\ = (\sigma'_x \sigma)(X(t_0), t_0) \left(\frac{1}{2} (B(t)^2 - B(t_0)^2) - (t-t_0) - B(t_0)(B(t) - B(t_0)) \right) = (\sigma'_x \sigma)(X(t_0), t_0) \left(\frac{1}{2} (B(t) - B(t_0))^2 - \frac{1}{2} (t-t_0) \right)$$

we get the most used higher order Milstein method.

Weak numerical solutions

N.5

In a strong solution one is given $\mu(x,t)$, $\sigma(x,t)$ and $B(t)$ and should use them to manufacture $X(t)$.

In a weak solution one is given $\mu(x,t)$, $\sigma(x,t)$ and should use them to manufacture $B(t)$ and $X(t)$.

In theory the difference is bigger than it may seem in practice. There exist SDE which has unique weak solution but no strong solution, e.g., Tanaka SDE $dX(t) = \text{sign}(X(t)) dB(t)$.

Uniqueness of strong solution means $P(X_1(t) = X_2(t)) = 1$ for all solutions $X_1(t)$ and $X_2(t)$

Uniqueness of weak solution means $F_{X_1(t_1), \dots, X_n(t_n)}(x_1, \dots, x_n) = F_{X_2(t_1), \dots, X_n(t_n)}(x_1, \dots, x_n)$.

The numerical methods we have seen so far are strong numerical methods in the sense that if we are given $\mu(x,t)$, $\sigma(x,t)$ and $B(t)$ we can plug in them in method and get $X(t)$ - example Euler, implicit Euler, Milstein.

When we use these methods in computer no one gives us $B(t)$ so we still create $B(t)$ ourselves to get going.

Example (Langevin equation) $dX(t) = -\alpha X(t) dt + \sigma dB(t)$

Euler method $Y_{n+1} = Y_n - \alpha Y_n (t_{n+1} - t_n) + \sigma (B(t_{n+1}) - B(t_n))$

is used in applied lecture. Exact solution is

$$X(t) = e^{-\alpha t} \left(X_0 + \int_0^t \sigma e^{\alpha r} dB(r) \right)$$

which gives

$$X(t) = e^{-\alpha(t-s)} X(s) + \underbrace{e^{-\alpha t} \int_s^t \sigma e^{\alpha r} dB(r)}_{N(0, \frac{\sigma^2}{2\alpha}(1 - e^{-2\alpha(t-s)}))}$$

$$N(0, \frac{\sigma^2}{2\alpha}(1 - e^{-2\alpha(t-s)}))$$

$$f_{X(t)|X(s)}(y|x) = \frac{1}{\sqrt{2\pi \frac{\sigma^2}{2\alpha}(1 - e^{-2\alpha(t-s)})}} e^{-\frac{y - e^{-\alpha(t-s)}x}{\frac{\sigma^2}{2\alpha}(1 - e^{-2\alpha(t-s)})}}$$

so we can create an exact weak solution at times $0 = t_0 < t_1 < \dots < t_N = T$ by doing

$$X(t_{n+1}) = e^{-\alpha(t_{n+1} - t_n)} X(t_n) + NID(0, \frac{\sigma^2}{2\alpha}(1 - e^{-2\alpha(t_{n+1} - t_n)}))$$

$$n = 0, \dots, N-1, \quad X(0) = x_0$$

We cannot do an exact theoretical solution in

reality because we cannot calculate ^{N.7}
 $\int_0^t e^{rt} dB(t)$ exactly!

Remember that solution to SDE are
Markov processes. Their transition
probability densities

$$p(y, t, x, s) = f_{X(t)|X(s)}(y|x)$$

satisfy the Kolmogorov backward PDE

$$u(x, s) \frac{\partial p}{\partial x} + \frac{\sigma(x, s)^2}{2} \frac{\partial^2 p}{\partial x^2} + \frac{\partial p}{\partial s} = 0$$

and by solving that PDE for a $p(y, t, x, s)$
that satisfies

$$\lim_{s \uparrow t} \int_{-\infty}^{+\infty} g(y) p(y, t, x, s) dy = g(x)$$

for smooth $g(x)$ one can find $p(y, t, x, s)$.

But given an OK family of transition PDF
there exists a Markov process that
has these transition PDF's and that
Markov process is weak solution to
SDE.

One can create an exact weak
solution by the recursive scheme

simulate $X(t_{n+1})$ with PDF $p(\cdot, t_{n+1}, X(t_n), t_n)$
for $n=0, \dots, N-1$,

$$X(0) = X_0$$

To create a exact weak solution
is in principle same as finding the
Markov transition density.

And a numerical weak scheme
same as approximating the
Markov transition density.

The Euler method

$$Y_{n+1} = Y_n + \mu(Y_n, t_n)(t_{n+1} - t_n) + \sigma(Y_n, t_n)(\sqrt{t_{n+1} - t_n})Z_{n+1}$$

we have seen has strong convergence

$$\sqrt{E\left(\max_{0 \leq n \leq N} |Y_n - X(t_n)|^2\right)} \leq Ch^{1/2} \text{ for } h = \max_{0 \leq n \leq N-1} t_{n+1} - t_n$$

so

$$E\left(\max_{0 \leq n \leq N} |Y_n - X(t_n)|\right) \leq \sqrt{E\left(\max_{0 \leq n \leq N} |Y_n - X(t_n)|^2\right)} \leq Ch^{1/2}$$

But if we use same Euler method
to just approximate $g(X(t))$ with $g(Y_N)$

$$\text{we have } |E(g(Y_N)) - E(g(X(t)))| \leq Ch$$

N.9

The proof of this uses some quite difficult PDE techniques.

If we want to find $E(g(X(T)))$ in reality we have to generate many independent observations $\{g(Y_N^{(i)})\}_{i=1}^n$ of $g(Y_N)$ and then use Monte Carlo

$$E(g(X(T))) \approx E(g(Y_N)) \approx \frac{1}{n} \sum_{i=1}^n g(Y_N^{(i)})$$