

1.

$$T = 2\pi, \quad \Omega = \frac{2\pi}{T} = 1$$

$$\begin{aligned} a_n &= \frac{2}{T} \int_{-T/2}^{T/2} f(t) \cos(n\Omega t) dt = \frac{1}{\pi} \int_{-\pi}^{\pi} f(t) \cos(nt) dt = \frac{1}{\pi} \int_0^{\pi} \sin(t) \cos(nt) dt \\ &= \frac{1}{2\pi} \int_0^{\pi} \left(\sin((1-n)t) + \sin((1+n)t) \right) dt = \frac{1}{2\pi} \left[-\frac{\cos((1-n)t)}{1-n} - \frac{\cos((1+n)t)}{1+n} \right]_0^{\pi} \\ &= \frac{1}{2\pi} \left[\frac{\cos((n-1)t)}{n-1} - \frac{\cos((n+1)t)}{n+1} \right]_0^{\pi} = \frac{1}{2\pi} \left(\frac{\cos((n-1)\pi) - 1}{n-1} - \frac{\cos((n+1)\pi) - 1}{n+1} \right) \\ &= \begin{cases} 0 & \text{för } n = 2k + 1 \text{ (udda),} \\ \frac{2}{\pi} \frac{-1}{(2k+1)(2k-1)} & \text{för } n = 2k \text{ (jämnt),} \end{cases} \end{aligned}$$

speciellt: $a_0 = \frac{2}{\pi}$.

$$\begin{aligned} b_n &= \frac{2}{T} \int_{-T/2}^{T/2} f(t) \sin(n\Omega t) dt = \frac{1}{\pi} \int_{-\pi}^{\pi} f(t) \sin(nt) dt = \frac{1}{\pi} \int_0^{\pi} \sin(t) \sin(nt) dt \\ &= \frac{1}{2\pi} \int_0^{\pi} \left(\cos((n-1)t) - \cos((n+1)t) \right) dt = \frac{1}{2\pi} \left[\frac{\sin((n-1)t)}{n-1} - \frac{\sin((n+1)t)}{n+1} \right]_0^{\pi} \\ &= 0 \quad \text{för } n > 1, \end{aligned}$$

$$\begin{aligned} b_1 &= \frac{2}{T} \int_{-T/2}^{T/2} f(t) \sin(\Omega t) dt = \frac{1}{\pi} \int_{-\pi}^{\pi} f(t) \sin(t) dt = \frac{1}{\pi} \int_0^{\pi} \sin(t) \sin(t) dt \\ &= \frac{1}{2\pi} \int_0^{\pi} (1 - \cos(2t)) dt = \frac{1}{2\pi} \left[t - \frac{\sin(2t)}{2} \right]_0^{\pi} = \frac{1}{2}. \end{aligned}$$

$$\begin{aligned} f(t) &= \frac{1}{2} a_0 + \sum_{n=1}^{\infty} \left(a_n \cos n\Omega t + b_n \sin n\Omega t \right) \\ &= \frac{1}{\pi} + \frac{1}{2} \sin(t) - \frac{2}{\pi} \sum_{k=1}^{\infty} \frac{1}{(2k+1)(2k-1)} \cos(2kt) \end{aligned}$$

2. (a) Se boken.

(b) $\lambda_n = -((2n+1)\frac{\pi}{2})^2$, $X_n(x) = \sin((2n+1)\frac{\pi}{2}x)$, $n = 0, 1, 2, \dots$

3. (a) Den karakteristiska ekvationen är $r^2 - 9 = 0$ med rötterna $r_1 = -3$, $r_2 = 3$. Den allmänna lösningen blir

$$\begin{aligned} u(t) &= Ae^{-3t} + Be^{3t} \\ u'(t) &= -3Ae^{-t} + 3Be^{3t}. \end{aligned}$$

Begynnelsevillkoren ger

$$\begin{aligned} u_0 &= u(0) = A + B, \\ u_1 &= u'(0) = -3A + 3B, \end{aligned}$$

dvs $A = \frac{1}{6}(3u_0 - u_1)$, $B = \frac{1}{6}(3u_0 + u_1)$. Alltså:

$$u(t) = \frac{1}{6}(3u_0 - u_1)e^{-3t} + \frac{1}{6}(3u_0 + u_1)e^{3t} = u_0 \cosh(3t) + \frac{1}{3}u_1 \sinh(3t).$$

(b) Med $x_1 = u$, $x_2 = u'$ får vi

$$\begin{aligned}x_1' &= u' = x_2, \\x_2' &= u'' = 9u = 9x_1,\end{aligned}$$

dvs, på matrisform,

$$\begin{bmatrix} x_1' \\ x_2' \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ 9 & 0 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}.$$

(c) Matrisen har egenvärdena $\lambda_1 = -3$, $\lambda_2 = 3$ och egenvektorer $g_1 = \begin{bmatrix} 1 \\ -3 \end{bmatrix}$, $g_2 = \begin{bmatrix} 1 \\ 3 \end{bmatrix}$. Lösningen blir

$$x(t) = Ae^{\lambda_1 t} g_1 + Be^{\lambda_2 t} g_2 = Ae^{-3t} \begin{bmatrix} 1 \\ -3 \end{bmatrix} + Be^{3t} \begin{bmatrix} 1 \\ 3 \end{bmatrix}.$$

Begynnelsevillkoret ger

$$\begin{bmatrix} u_0 \\ u_1 \end{bmatrix} = x(0) = A \begin{bmatrix} 1 \\ -3 \end{bmatrix} + B \begin{bmatrix} 1 \\ 3 \end{bmatrix},$$

dvs $A = \frac{1}{6}(3u_0 - u_1)$, $B = \frac{1}{6}(3u_0 + u_1)$. Alltså:

$$x(t) = \begin{bmatrix} u(t) \\ u'(t) \end{bmatrix} = \frac{1}{6}(3u_0 - u_1)e^{-3t} \begin{bmatrix} 1 \\ -3 \end{bmatrix} + \frac{1}{6}(3u_0 + u_1)e^{3t} \begin{bmatrix} 1 \\ 3 \end{bmatrix}.$$

(d) Man skriver en m-fil kallad `funk.m`:

```
function xprime=funk(t,x)
xprime=[0 1; 9 0]*x;
```

sedan exekverar man matlabkommandona

```
>> u0=1; u1=2; x0=[u0; u1]; T=5;
>> [t,x]=ode45('funk',[0;T],x0);
```

(e) Det linjära systemet har egenvärdena -3 och 3 , dvs ett är negativt. Det betyder att systemet är instabilt.

4. Den första ekvationen divideras med $q_f c_f$, den andra med $\rho c_p q_f T_f$. Med $\tau = V/q_f$ får vi

$$\begin{aligned}\tau \frac{d}{dt} \left(\frac{c}{c_f} \right) &= \frac{q}{q_f} \frac{c_f - c}{c_f} - \frac{c}{c_f} \tau k_0 \exp \left(-\frac{E}{RT_f} \frac{T_f}{T} \right), \\ \tau \frac{d}{dt} \left(\frac{T}{T_f} \right) &= \frac{q}{q_f} \frac{T_f - T}{T_f} + \frac{(-\Delta H)c_f}{\rho c_p T_f} \frac{c}{c_f} \tau k_0 \exp \left(-\frac{E}{RT_f} \frac{T_f}{T} \right) - \frac{\kappa A \tau}{\rho c_p V} \left(\frac{T - T_f}{T_f} - \frac{T_K - T_f}{T_f} \right).\end{aligned}$$

Med de dimensionslösa variablerna

$$\begin{aligned}s &= t/\tau, \\ X_1(s) &= \frac{c(s\tau)}{c_f}, \quad X_2(s) = \frac{T(s\tau)}{T_f}, \\ U_1(s) &= \frac{q(s\tau)}{q_f}, \quad U_2(s) = \frac{T_K(s\tau)}{T_f}, \\ \alpha &= \frac{(-\Delta H)c_f}{\rho c_p T_f}, \quad \beta = \frac{\kappa A \tau}{\rho c_p V}, \quad \gamma = \frac{E}{RT_f}, \quad \delta = \tau k_0 e^{-\gamma}, \quad f(X_2) = \delta \exp(\gamma - \gamma/X_2),\end{aligned}$$

får vi

$$\begin{aligned} \frac{c}{c_f} &= X_1, & \frac{T}{T_f} &= X_2, & \tau \frac{d}{dt} \left(\frac{c}{c_f} \right) &= \frac{dX_1}{ds}, & \tau \frac{d}{dt} \left(\frac{T}{T_f} \right) &= \frac{dX_2}{ds}, \\ \tau k_0 \exp \left(-\frac{E}{RT_f} \frac{T_f}{T} \right) &= \tau k_0 \exp \left(-\frac{E}{RT_f} \right) \exp \left(\frac{E}{RT_f} - \frac{E}{RT_f} \frac{T_f}{T} \right) \\ &= \delta \exp \left(\gamma - \frac{\gamma}{X_2} \right) = f(X_2). \end{aligned}$$

Detta leder till det icke-linjära differentialekvationssystemet

$$\begin{aligned} \frac{dX_1}{ds} &= U_1(1 - X_1) - X_1 f(X_2) && (= F_1(X, U)), \\ \frac{dX_2}{ds} &= U_1(1 - X_2) + \alpha X_1 f(X_2) - \beta(X_2 - U_2) && (= F_2(X, U)). \end{aligned}$$

(b) Stationära punkter ges av

$$\begin{aligned} 0 &= \bar{U}_1(1 - \bar{X}_1) - \bar{X}_1 f(\bar{X}_2), \\ 0 &= \bar{U}_1(1 - \bar{X}_2) + \alpha \bar{X}_1 f(\bar{X}_2) - \beta(\bar{X}_2 - \bar{U}_2). \end{aligned}$$

Vi löser ut styrvariablerna

$$\begin{aligned} \bar{U}_1 &= \frac{\bar{X}_1}{1 - \bar{X}_1} f(\bar{X}_2), \\ \bar{U}_2 &= \bar{X}_2 - \frac{1}{\beta} \left(\frac{\bar{X}_1}{1 - \bar{X}_1} (1 - \bar{X}_2) f(\bar{X}_2) + \alpha \bar{X}_1 f(\bar{X}_2) \right). \end{aligned}$$

Med $\alpha = 0.3$, $\gamma = 30$, $\delta = 0.1$ och $\bar{X}_1 = 0.5$, $\bar{X}_2 = 1$ får vi $f(1) = 0.1$ och $f'(1) = 3$, $\alpha f(1) = 0.03 \approx 0$. Dvs

$$\begin{aligned} \bar{U}_1 &= 0.1, \\ \bar{U}_2 &= 1 - \frac{0.015}{\beta}. \end{aligned}$$

(c) Nu linjäriserar vi kring \bar{X}, \bar{U} . Det linjäriserade systemet blir

$$x'(s) = Ax(s) + Bu(s), \quad s > 0; \quad x(0) = x_0,$$

med Jacobimatrisererna

$$A = \begin{bmatrix} \frac{\partial F_1}{\partial X_1}(\bar{X}, \bar{U}) & \frac{\partial F_1}{\partial X_2}(\bar{X}, \bar{U}) \\ \frac{\partial F_2}{\partial X_1}(\bar{X}, \bar{U}) & \frac{\partial F_2}{\partial X_2}(\bar{X}, \bar{U}) \end{bmatrix} = \begin{bmatrix} -\bar{U}_1 - f(\bar{X}_2) & -\bar{X}_1 f'(\bar{X}_2) \\ \alpha f(\bar{X}_2) & -\bar{U}_1 + \alpha \bar{X}_1 f'(\bar{X}_2) - \beta \end{bmatrix},$$

där $f'(\bar{X}_2) = \frac{\gamma}{\bar{X}_2^2} f(\bar{X}_2)$, och

$$B = \begin{bmatrix} \frac{\partial F_1}{\partial U_1}(\bar{X}, \bar{U}) & \frac{\partial F_1}{\partial U_2}(\bar{X}, \bar{U}) \\ \frac{\partial F_2}{\partial U_1}(\bar{X}, \bar{U}) & \frac{\partial F_2}{\partial U_2}(\bar{X}, \bar{U}) \end{bmatrix} = \begin{bmatrix} 1 - \bar{X}_1 & 0 \\ 1 - \bar{X}_2 & \beta \end{bmatrix}.$$

Vi sätter in siffervärden i A ,

$$A = \begin{bmatrix} -0.2 & -1.5 \\ 0.03 & 0.35 - \beta \end{bmatrix} \approx \begin{bmatrix} -0.2 & -1.5 \\ 0 & 0.35 - \beta \end{bmatrix}$$

Den sista matrisen har egenvärdena $\lambda_1 = -0.2$ och $\lambda_2 = 0.35 - \beta$. De är båda negativa, och det linjäriserade systemet asymptotiskt stabilt, om $\beta > 0.35$. Det betyder fysikaliskt att kylarens area måste vara tillräckligt stor, större än $0.35 \rho c_p q_f / \kappa$.

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