## ROBIN BOUNDARY CONDITIONS

## 1. Modeling

As an example, we consider the following mathematical model of a stationary reaction-diffusion process involving a single substance,

$$-(au')' + cu = f, x_{\min} < x < x_{\max},$$

(1) 
$$a(x_{\min})u'(x_{\min}) = \gamma(x_{\min})(u(x_{\min}) - g_D(x_{\min})) + g_N(x_{\min}), -a(x_{\max})u'(x_{\max}) = \gamma(x_{\max})(u(x_{\max}) - g_D(x_{\max})) + g_N(x_{\max}),$$

where u(x), denoting the *concentration* of the substance, is the unknown function that we wish to compute. The following functions are data to the problem:

$$\begin{array}{ll} a(x): & \textit{diffusion coefficient} \\ c(x): & \textit{rate coefficient} \end{array} \qquad \begin{array}{ll} (a(x)>0) \\ (c(x)\geq 0) \end{array}$$

f(x): source

 $g_D(x_{\min}), g_D(x_{\max}): ambient concentration$ 

 $g_N(x_{\min}), g_N(x_{\max})$ : externally induced flux through the boundary

First, we consider the case  $g_N(x_{\min}) = g_N(x_{\max}) = 0$ , for which Robin boundary conditions are a mathematical model of the physical fact that the outward flux is proportional to the concentration difference between the domain boundary and its surroundings. We have the following special cases:

**Homogeneous Neumann boundary condition:** This boundary condition physically corresponds to the case of an *impermeable* boundary, i.e., one for which  $\gamma = 0$ , implying zero flux through the boundary: u' = 0.

**Dirichlet boundary condition:** This boundary condition physically corresponds to the case of a very high permeability, i.e.,  $\gamma \to +\infty$ , implying that the concentration at the boundary adapts to the ambient concentration:  $u = g_D$ .

We can also imagine a case where we externally control the flux through the boundary. This case can be modelled by choosing  $\gamma = 0$  and  $g_N \neq 0$ :

**Inhomogeneous Neumann boundary condition:** This boundary condition prescribes the flux through the boundary:

$$a(x_{\min})u'(x_{\min}) = g_N(x_{\min}), -a(x_{\max})u'(x_{\max}) = g_N(x_{\max}).$$

## 2. Variational Formulation

To derive the variational formulation of (1), we multiply the differential equation by v(x) and integrate over  $[x_{\min}, x_{\max}]$ ,

$$- \int_{x_{\min}}^{x_{\max}} (au')' v \, dx + \int_{x_{\min}}^{x_{\max}} cuv \, dx = \int_{x_{\min}}^{x_{\max}} fv \, dx.$$

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We now integrate by parts,

$$[-(au')v]_{x=x_{\min}}^{x=x_{\max}} + \int_{x_{\min}}^{x_{\max}} au'v' \, dx + \int_{x_{\min}}^{x_{\max}} cuv \, dx = \int_{x_{\min}}^{x_{\max}} fv \, dx.$$

Use the boundary conditions in (1),

$$\begin{array}{lcl} a(x_{\min})u'(x_{\min}) & = & \gamma(x_{\min})(u(x_{\min}) - g_D(x_{\min})) + g_N(x_{\min}), \\ -a(x_{\max})u'(x_{\max}) & = & \gamma(x_{\max})(u(x_{\max}) - g_D(x_{\max})) + g_N(x_{\max}), \end{array}$$

to obtain,

$$\gamma(x_{\max})u(x_{\max})v(x_{\max}) \; + \; \gamma(x_{\min})u(x_{\min})v(x_{\min}) \; + \; \int_{x_{\min}}^{x_{\max}} au'v' \, dx \; + \; \int_{x_{\min}}^{x_{\max}} cuv \, dx \; = \\ (\gamma(x_{\max})g_D(x_{\max}) - g_N(x_{\max}))v(x_{\max}) \; + \; (\gamma(x_{\min})g_D(x_{\min}) - g_N(x_{\min}))v(x_{\min}) \; + \; \int_{x_{\min}}^{x_{\max}} fv \, dx.$$

We now state the following variational formulation of (1):

Find  $u(x) \in H^1([x_{\min}, x_{\max}]) := \left\{ v(x) : \int_{x_{\min}}^{x_{\max}} v(x)^2 dx < \infty, \int_{x_{\min}}^{x_{\max}} v'(x)^2 dx < \infty \right\}$ , such that

$$\gamma(x_{\max})u(x_{\max})v(x_{\max}) + \gamma(x_{\min})u(x_{\min})v(x_{\min}) + \int_{x_{\min}}^{x_{\max}} au'v' dx + \int_{x_{\min}}^{x_{\max}} cuv dx =$$

(2) 
$$(\gamma(x_{\max})g_D(x_{\max}) - g_N(x_{\max}))v(x_{\max}) + (\gamma(x_{\min})g_D(x_{\min}) - g_N(x_{\min}))v(x_{\min}) + \int_{x_{\min}}^{x_{\max}} fv \, dx, \quad \forall v \in H^1([x_{\min}, x_{\max}]).$$

## 3. The Finite Element Method (FEM)

3.1. **Discretization.** Introducing the vector space,  $V_h$ , of continuous, piecewise linear functions on a partition,  $x_{\min} = x_1 < x_2 < \ldots < x_{N-1} < x_N = x_{\max}$ , of  $[x_{\min}, x_{\max}]$ , we now state the cG(1) method<sup>1</sup> as the following discrete counterpart of (2):

Find  $U(x) \in V_h$ , such that

$$\gamma(x_N)U(x_N)v(x_N) + \gamma(x_1)U(x_1)v(x_1) + \int_{x_1}^{x_N} aU'v' dx + \int_{x_1}^{x_N} cUv dx =$$

$$(3) \qquad (\gamma(x_N)g_D(x_N) - g_N(x_N))v(x_N) + (\gamma(x_1)g_D(x_1) - g_N(x_1))v(x_1) +$$

$$\int_{x_1}^{x_N} fv dx, \quad \forall v \in V_h.$$

 $<sup>^{1}</sup>$ In cG(1), the letter c stands for continuous and the number 1 stands for linear, expressing the fact that this finite element method is based on continuous, piecewise linear approximation. The letter G stands for Galerkin. Boris Grigorievich Galerkin (1871 - 1945) was a Russian mathematician who made pioneering contributions to the field of numerical solution of differential equations. The Galerkin method is the method of rewriting the differential equation in variational form, and discretize this. A Finite Element Method (FEM), is a Galerkin method that utilises piecewise polynomials as approximating functions.

3.2. **Ansatz.** We now seek a solution, U(x), to (3), expressed in the basis of hat functions  $\{\varphi_i\}_{i=1}^N \subset V_h$ , defined by  $\varphi_i \in V_h$ ,  $\varphi_i(x_j) = \delta_{ij}$ ,  $i, j = 1, \ldots, N$ . (Here,  $\delta_{ij}$  denotes the Kronecker delta function, which is defined by the property  $\delta_{ij} = \begin{cases} 1, & \text{if } i = j, \\ 0, & \text{if } i \neq j. \end{cases}$ ) In other words, we make the Ansatz

(4) 
$$U(x) = \sum_{j=1}^{N} \xi_j \varphi_j(x),$$

and seek to determine the coefficient vector,

$$\xi = \left[ egin{array}{c} \xi_1 \ \xi_2 \ dots \ \xi_N \end{array} 
ight] = \left[ egin{array}{c} U(x_1) \ U(x_2) \ dots \ U(x_N) \end{array} 
ight],$$

of nodal values of U(x), in such a way that (3) is satisfied.

3.3. Construction of discrete system of linear equations. We substitute (4) into (3),

$$\gamma(x_{N})\xi_{N}v(x_{N}) + \gamma(x_{1})\xi_{1}v(x_{1}) + \sum_{j=1}^{N} \xi_{j} \left\{ \int_{x_{1}}^{x_{N}} a \varphi_{j}'v' dx + \int_{x_{1}}^{x_{N}} c \varphi_{j}v dx \right\} =$$

$$(5) \qquad (\gamma(x_{N})g_{D}(x_{N}) - g_{N}(x_{N}))v(x_{N}) + (\gamma(x_{1})g_{D}(x_{1}) - g_{N}(x_{1}))v(x_{1}) +$$

$$\int_{x_{1}}^{x_{N}} fv dx, \quad \forall v \in V_{h}.$$

Since  $\{\varphi_i\}_{i=1}^N \subset V_h$  is a *basis* of  $V_h$ , (5) is equivalent to,

$$\gamma(x_{N})\xi_{N}\varphi_{i}(x_{N}) + \gamma(x_{1})\xi_{1}\varphi_{i}(x_{1}) + \sum_{j=1}^{N} \xi_{j} \left\{ \int_{x_{1}}^{x_{N}} a \, \varphi_{j}' \varphi_{i}' \, dx + \int_{x_{1}}^{x_{N}} c \, \varphi_{j}\varphi_{i} \, dx \right\} =$$
(6)
$$(\gamma(x_{N})g_{D}(x_{N}) - g_{N}(x_{N}))\varphi_{i}(x_{N}) + (\gamma(x_{1})g_{D}(x_{1}) - g_{N}(x_{1}))\varphi_{i}(x_{1}) +$$

$$\int_{x_{1}}^{x_{N}} f\varphi_{i} \, dx, \quad i = 1, \dots, N,$$

which is a quadratic system of N linear equations and N unknowns. Introducing the notation

$$a_{ij} = \int_{x_1}^{x_N} a \, \varphi_j' \varphi_i' \, dx,$$
  $m_{c\,ij} = \int_{x_1}^{x_N} c \, \varphi_j \varphi_i \, dx,$   $b_i = \int_{x_1}^{x_N} f \varphi_i \, dx,$ 

and taking into account that  $\varphi_i(x_1) = \begin{cases} 1, & \text{if } i = 1, \\ 0, & \text{if } i \neq 1, \end{cases}$  and  $\varphi_i(x_N) = \begin{cases} 1, & \text{if } i = N, \\ 0, & \text{if } i \neq N, \end{cases}$  we can

$$\begin{cases} (\gamma(x_{1}) + a_{11} + m_{c \, 11})\xi_{1} & + & \dots & + & (a_{1\mathrm{N}} + m_{c \, 1\mathrm{N}})\xi_{\mathrm{N}} & = & b_{1} + \gamma(x_{1})g_{D}(x_{1}) - g_{N}(x_{1}) \\ (a_{21} + m_{c \, 21})\xi_{1} & + & \dots & + & (a_{2\mathrm{N}} + m_{c \, 2\mathrm{N}})\xi_{\mathrm{N}} & = & b_{2} \\ & \vdots & & \vdots & & \vdots & & \vdots \\ (a_{N-1\, 1} + m_{c \, N-1\, 1})\xi_{1} & + & \dots & + & (a_{N-1\, N} + m_{c \, N-1\, N})\xi_{\mathrm{N}} & = & b_{N-1} \\ (a_{N1} + m_{c \, N1})\xi_{1} & + & \dots & + & (a_{N\mathrm{N}} + m_{c \, N\mathrm{N}} + \gamma(x_{\mathrm{N}}))\xi_{\mathrm{N}} & = & b_{\mathrm{N}} + \gamma(x_{\mathrm{N}})g_{D}(x_{\mathrm{N}}) - g_{\mathrm{N}}(x_{\mathrm{N}}) \end{cases}$$

In matrix form, this reads,

$$(A + M_c + R) \xi = b + rv,$$

where 
$$A = \left[ \begin{array}{ccc} a_{11} & \dots & a_{1\mathrm{N}} \\ \vdots & \ddots & \vdots \\ a_{\mathrm{N}1} & \dots & a_{\mathrm{N}\mathrm{N}} \end{array} \right]$$
 is the  $\mathit{stiffness\ matrix},$ 

$$M_c = \left[ egin{array}{cccc} m_{c \; 11} & \ldots & m_{c \; 1N} \ & dots & \ddots & dots \ & m_{c \; \mathrm{N}1} & \ldots & m_{c \; \mathrm{N}N} \end{array} 
ight] ext{ is the } extit{mass matrix},$$

$$R = \left[ \begin{array}{ccccc} \gamma(x_1) & 0 & \dots & 0 & 0 \\ 0 & 0 & \dots & 0 & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & \dots & 0 & 0 \\ 0 & 0 & \dots & 0 & \gamma(x_{\mathrm{N}}) \end{array} \right] \text{ contains the boundary contributions to the system matrix,}$$

$$b = \left[egin{array}{c} b_1 \ dots \ b_{
m N} \end{array}
ight] ext{ is the } load\ vector, ext{ and }$$

$$rv = \left[ \begin{array}{c} \gamma(x_1)g_D(x_1) - g_N(x_1) \\ \\ 0 \\ \\ \vdots \\ \\ 0 \\ \\ \gamma(x_{\rm N})g_D(x_{\rm N}) - g_N(x_{\rm N}) \end{array} \right] \mbox{contains the boundary contributions to the right hand side.}$$