

# *Error estimation and adaptivity for multiscale methods*

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# Outline and Papers

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## Outline

- Model problem
- Motivation and previous work
- Derivation of proposed method with examples
- A posteriori error estimates
- Adaptive algorithm
- Application to oil recovery problem
- Conclusions and future work

## Papers

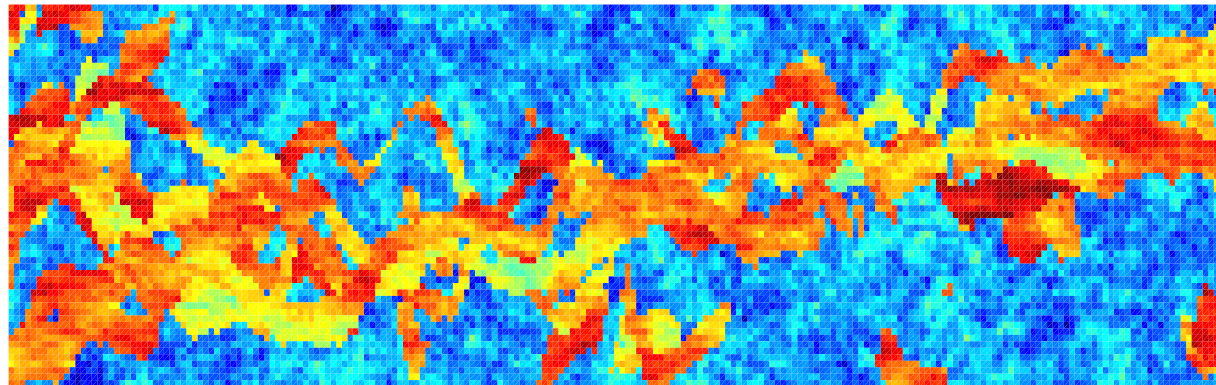
- M.G. Larson and A. Målqvist, *Adaptive Variational Multiscale Methods Based on A Posteriori Error Estimation: Energy Norm Estimates for Elliptic Problems*, CMAME 2007
- A. Målqvist, *Multiscale methods for elliptic problems* (in review MMS)

## Model Problem

**Elliptic equation:** Find  $u$  such that

$$\mathcal{L}u = -\nabla \cdot a \nabla u + \nabla \cdot (bu) + cu = f \quad \text{in } \Omega$$

where  $\mathcal{L}$  is elliptic with multiscale coefficients,  $f$  is a given function, and  $\Omega \subset \mathbb{R}^d$  is a domain.



**Weak form:** Let  $a(v, w) = (\mathcal{L}v, w)$  and  $l(w) = (f, w)$  for all  $v, w \in \mathcal{V}$ , for an appropriate function space  $\mathcal{V}$ . Find  $u \in \mathcal{V}$  such that,

$$a(u, v) = l(v), \quad \text{for all } v \in \mathcal{V}.$$

## Motivation: Why Resolve the Coefficients?

**Simple periodic example:** Consider the Poisson equation

$$\mathcal{L}u = -\nabla \cdot a \nabla u = f,$$

with periodic coefficient  $a = a(x/\epsilon)$  solved using the finite element method on a mesh of size  $H$ , we have (Hou-Wu-Cai),

$$\|\sqrt{a} \nabla (u - u_h)\|_{L^2(\Omega)} \leq C \frac{H}{\epsilon} \|f\|_{L^2(\Omega)},$$

- $\epsilon < H$  will give unreliable results even with exact quadrature.
- $\epsilon > H$  will be too computationally expensive to solve on a single mesh.

**From now on we assume nothing on the coefficients, more than what is needed to guarantee existence and uniqueness.**

## Some Previous Works and Related Methods

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- Upscaling techniques: Durlinsky et al. 98, Nielsen et al. 98
- Variational multiscale method: Hughes et al. 95, Arbogast 04, Larson-Målqvist 05, Nolen et al. 08, Nordbotten 09
- Multiscale finite element method: Hou-Wu 96, Efendiev-Ginting 04, Aarnes-Lie 06
- Multiscale finite volume method: Jenny et al. 03
- Heterogeneous multiscale method: Engquist-E 03, E-Ming-Zhang 04
- Equation free: Kevrekidis et al. 05
- ...

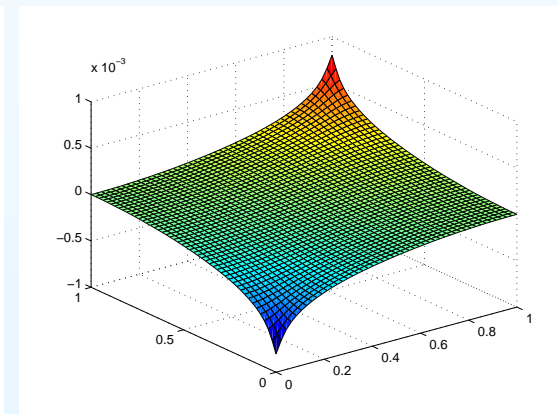
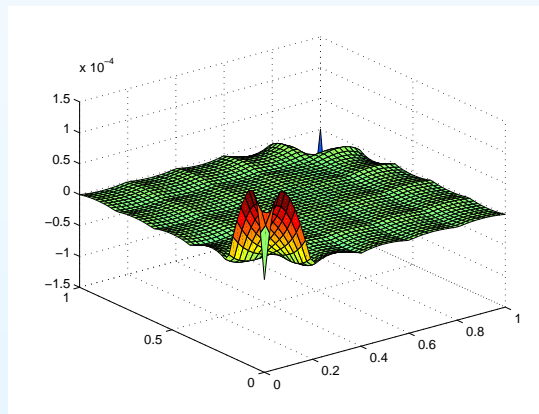
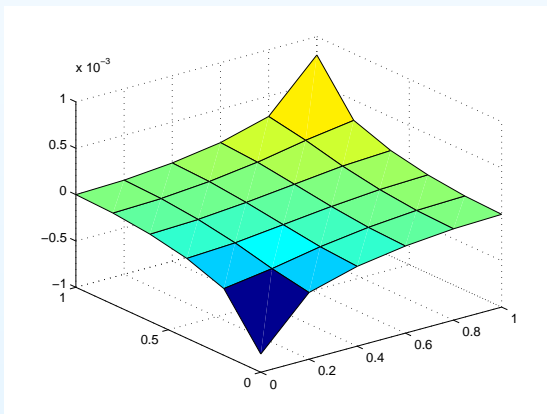
Local approximations (in parallel) on a fine scale are used to modify the coarse scale equation.

# The Variational Multiscale Method (VMS)

The weak form reads: find  $u \in \mathcal{V}$  such that,

$$a(u, v) = l(v), \quad \text{for all } v \in \mathcal{V}.$$

Now let  $\mathcal{V}_c \oplus \mathcal{V}_f = \mathcal{V}$ .



- $\mathcal{V}_c$  is a finite dimensional approximation of  $\mathcal{V}$ . (FE space)
- $\mathcal{V}_f$  can be chosen as e.g. hierarchical basis,  $L^2(\Omega)$ -orthogonal to  $\mathcal{V}_c$ , or wavelet modified hierarchical basis.

## $a$ -Orthogonal Split

We want to decouple the coarse and the fine scales. We introduce two multiscale projection operators, let  $\mathcal{T}, \mathcal{T}^* : \mathcal{V}_c \rightarrow \mathcal{V}_f$  such that

$$\begin{aligned} a(v_c + \mathcal{T}v_c, v_f) &= 0, & \text{for all } v_c \in \mathcal{V}_c, v_f \in \mathcal{V}_f \\ a(v_f, v_c + \mathcal{T}^*v_c) &= 0, & \text{for all } v_c \in \mathcal{V}_c, v_f \in \mathcal{V}_f \end{aligned}$$

Let  $u = u_c + \mathcal{T}u_c + u_f \in \mathcal{V}_c \oplus \mathcal{V}_f$  and  $v = v_c + \mathcal{T}^*v_c + v_f$  in the weak form,

$$\begin{aligned} a(u_c + \mathcal{T}u_c, v_c + \mathcal{T}^*v_c) &= l(v_c + \mathcal{T}^*v_c) & \text{for all } v_c \in \mathcal{V}_c, \\ a(u_f, v_f) &= l(v_f) & \text{for all } v_f \in \mathcal{V}_f, \end{aligned}$$

since  $a(u_c + \mathcal{T}u_c, v_f) = a(u_f, v_c + \mathcal{T}^*v_c) = 0$ .

## Modified Coarse Scale Equations

Find  $u_c \in \mathcal{V}_c$  such that

$$a(u_c + \mathcal{T}u_c, v_c + \mathcal{T}^*v_c) = l(v_c + \mathcal{T}^*v_c) - a(u_f, v_c + \mathcal{T}^*v_c)$$

for all  $v_c \in \mathcal{V}_c$  (note that  $a(u_f, v_f) = l(v_f)$  for all  $v_f \in \mathcal{V}_f$ ).

- If  $a$  is symmetric we get  $\mathcal{T}^*v_c = \mathcal{T}v_c$  i.e. a symmetric formulation
- In standard VMS  $\mathcal{T}^*v_c := 0$  and  $u_f$  is included in the coarse scale. The computation of  $\mathcal{T}u_c + u_f \in \mathcal{V}_f$  is decoupled and done analytically on each coarse element using homogeneous Dirichlet boundary conditions.
- In MsFEM  $u_f = 0$  and therefore not present in the coarse scale equations. Here the computation of  $v_c + \mathcal{T}v_c$  is decoupled and solved numerically on each element (or larger domains) using approximate boundary conditions.



## Approximation of $\mathcal{T}$ , $\mathcal{T}^*$ , and $u_f$

Let  $u_c = \sum_i u_c^i \phi_i$  with  $\{\phi_i\}$  a basis in  $\mathcal{V}_c$ .

We have  $\mathcal{T}u_c = \sum_i u_c^i \mathcal{T}\phi_i$ ,  $\mathcal{T}^*v_c = \sum_i v_c^i \mathcal{T}^*\phi_i$ , and  $u_f = \sum_i u_{f,i}$  where,

$$a(\mathcal{T}\phi_i, v_f) = -a(\phi_i, v_f) \quad \text{for all } v_f \in \mathcal{V}_f,$$

$$a(v_f, \mathcal{T}^*\phi_i) = -a(v_f, \phi_i) \quad \text{for all } v_f \in \mathcal{V}_f,$$

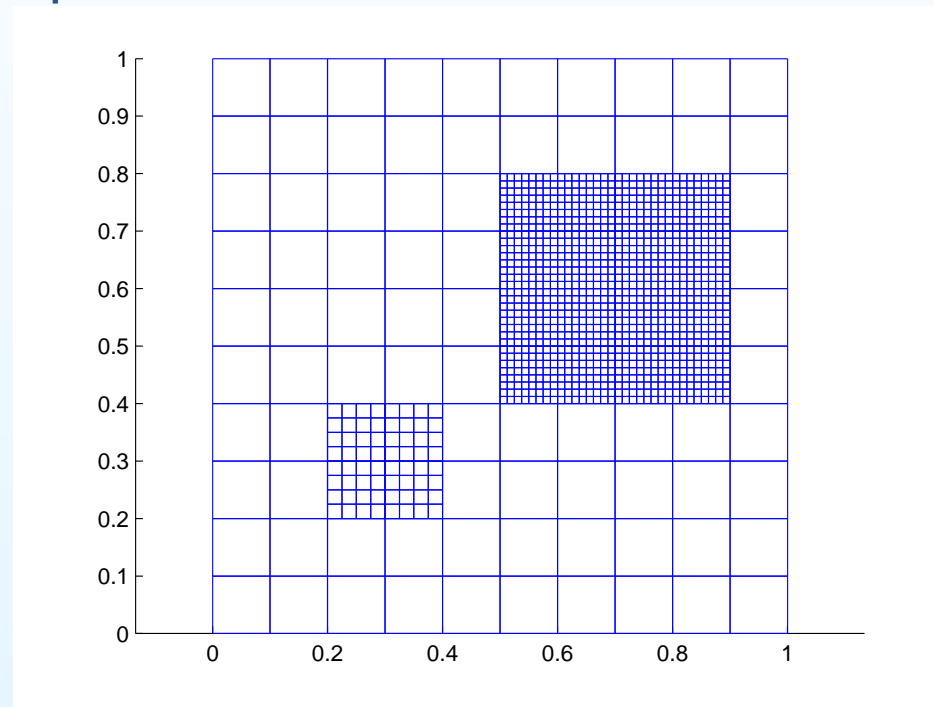
$$a(u_{f,i}, v_f) = (f\phi_i, v_f) \quad \text{for all } v_f \in \mathcal{V}_f.$$

We compute approximations  $\tilde{\mathcal{T}}\phi_i$ ,  $\tilde{\mathcal{T}}^*\phi_i$ , and  $U_{f,i}$  by

- restricting to a localized patches  $\text{supp}(\phi_i) \subset \omega_i$ ,
- discretizing with a fine subgrid on  $\omega_i$ ,
- and using homogeneous boundary conditions on the patch.

## Refinements and Layers

We let  $H$  be coarse scale mesh size and  $h$  be fine scale mesh size. Further we let  $L$  denote the number of layers of coarse elements in the patch.



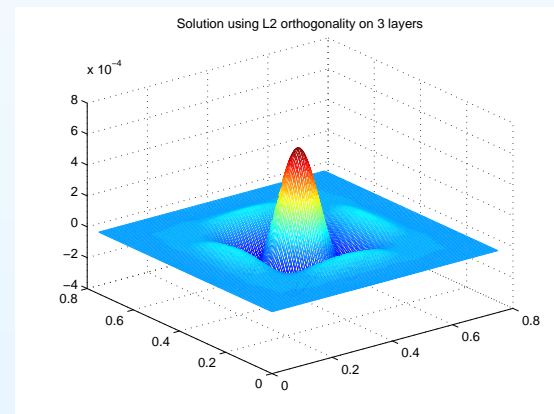
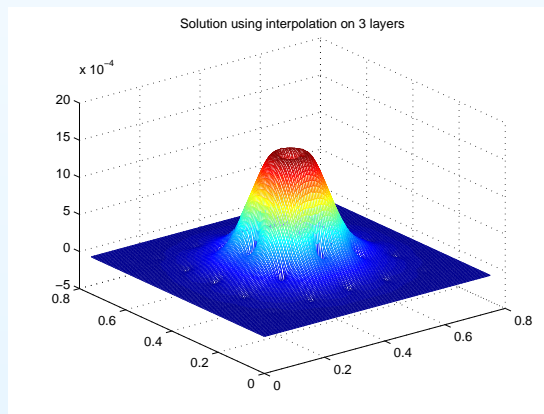
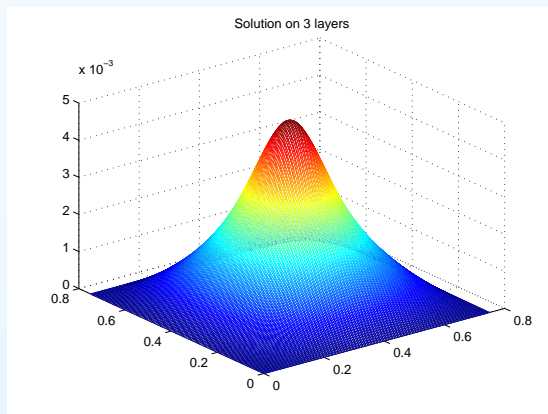
Since we exclude  $\mathcal{V}_c$  from the fine scale the condition number is only  $\sim \left(\frac{H}{h}\right)^2 \log(H/h)$  using an hierarchical split (Marion-Xu 95).

## Simple Observation About Decay in $\mathcal{V}_f$ (Fourier)

Consider the Laplace equation,

$$-\Delta u = \varphi_i \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega,$$

where  $\varphi_i$  has local support in  $\Omega$ . The weak form reads: find  $u \in \mathcal{W}$  s.t.,  $(\nabla u, \nabla v) = (\varphi_i, v)$  for all  $v \in \mathcal{W}$ .

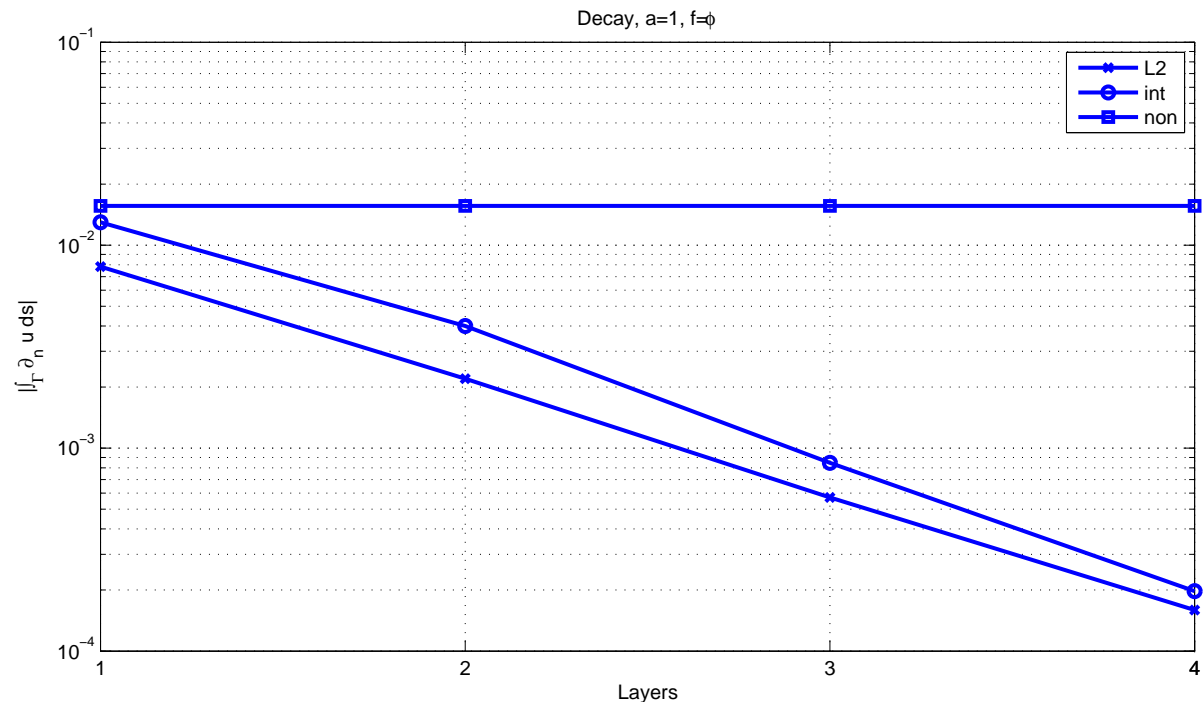


To the left  $\mathcal{W} = \mathcal{V}_c \oplus \mathcal{V}_f$ , middle  $\mathcal{W} = \mathcal{V}_f$  using hierarchical split, and right  $\mathcal{W} = \mathcal{V}_f$  using  $L^2$ -orthogonal split.

Constraints are realized using Lagrangian multipliers.

## Simple Observation About Decay in $\mathcal{V}_f$

Decay of flux integrated over the boundary.



We see exponential decay with respect distance measured in number of coarse elements. This effect gives rapid convergence as the patch size increases.

## The Proposed Multiscale Method

- Let  $\tilde{\mathcal{T}}, \tilde{\mathcal{T}}^*$  be the computable approximations of  $\mathcal{T}, \mathcal{T}^*$
- Let  $U_f$  be the computable approximation of  $u_f$

We get: find  $U_c \in \mathcal{V}_c$  such that

$$a(U_c + \tilde{\mathcal{T}}U_c, v_c + \tilde{\mathcal{T}}^*v_c) = l(v_c + \tilde{\mathcal{T}}^*v_c) - a(U_f, v_c + \tilde{\mathcal{T}}^*v_c)$$

for all  $v_c \in \mathcal{V}_c$ . On matrix form this leads to a system,

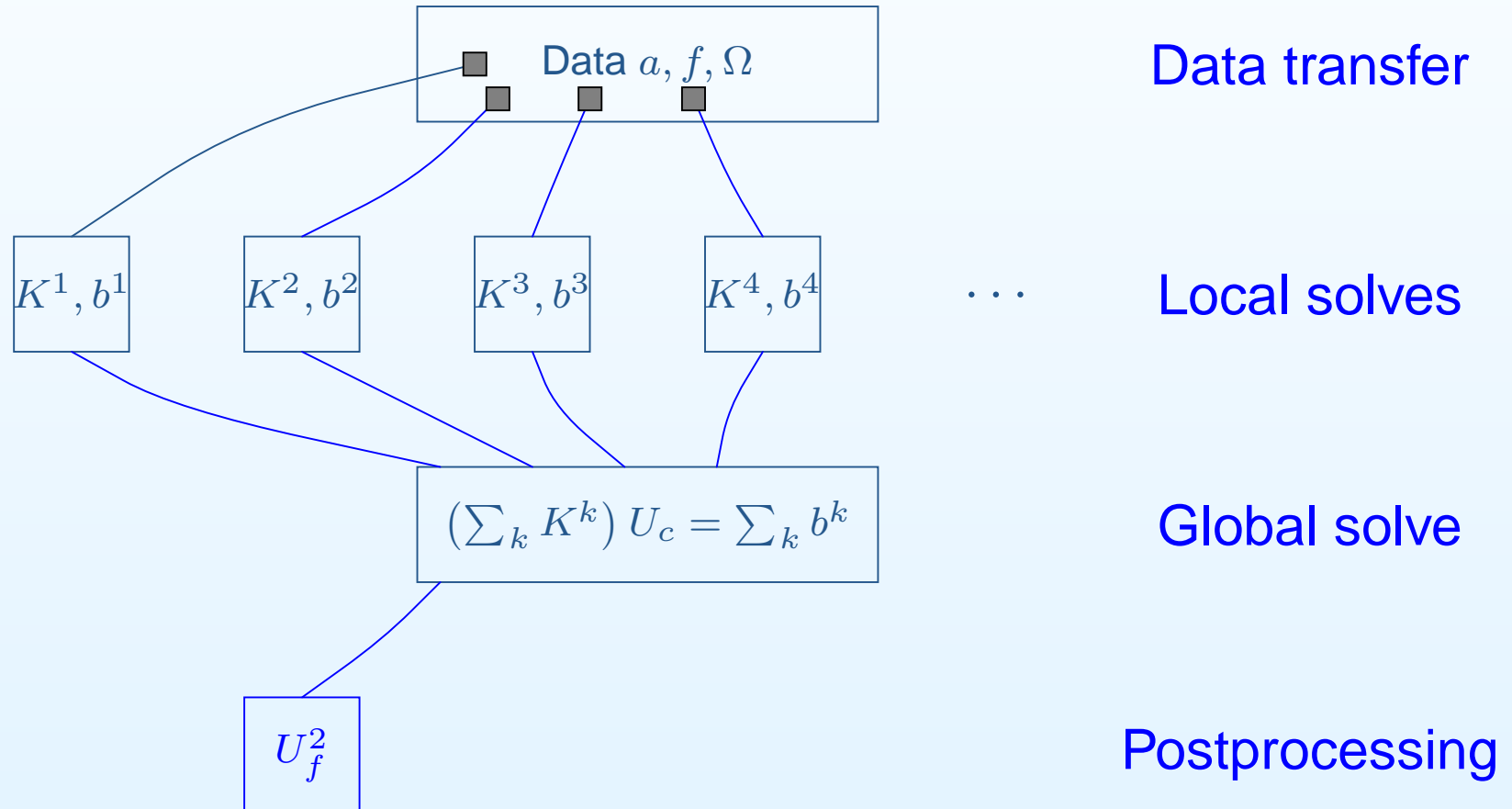
$$KU_c = b$$

Given  $U_c, U_f$ , and  $\tilde{\mathcal{T}}, U$  can be computed.

When  $a(U_f, v_c + \tilde{\mathcal{T}}^*v_c)$  is included, the error is orthogonal to  $v_c + \tilde{\mathcal{T}}^*v_c$  i.e.  $a(u - U_c - \tilde{\mathcal{T}}U_c - U_f, v_c + \tilde{\mathcal{T}}^*v_c) = 0$ .

# Parallel Structure

One local problem for each coarse dof, minimal communication.



## Examples of Applications

**Oil recovery:** We seek water saturation  $s$  (oil is  $1 - s$ )

$$\sigma - a\lambda(s)\nabla u = 0$$

$$-\nabla \cdot \sigma = q$$

$$\dot{s} + \sigma \cdot \nabla f(s) = 0$$

$f(s)$  is fractional flow function,  $\lambda(s)$  is total mobility,  $a$  is permeability, and  $q$  is a source term.

**Organic Semi-Conductor:** We seek electric potential  $\phi$ , hole concentration  $p$ , and electron concentration  $n$ ,

$$\dot{n} - \nabla \cdot (D_n \nabla n - \mu_n n \nabla \phi) = 0$$

$$\dot{p} - \nabla \cdot (D_p \nabla p + \mu_p p \nabla \phi) = 0$$

$$-\nabla \cdot \epsilon \nabla \phi = p - n,$$

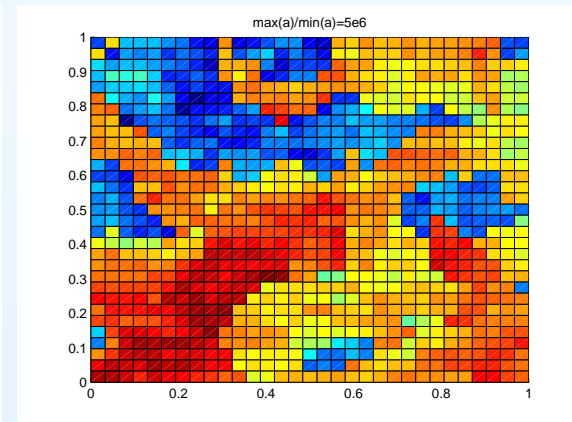
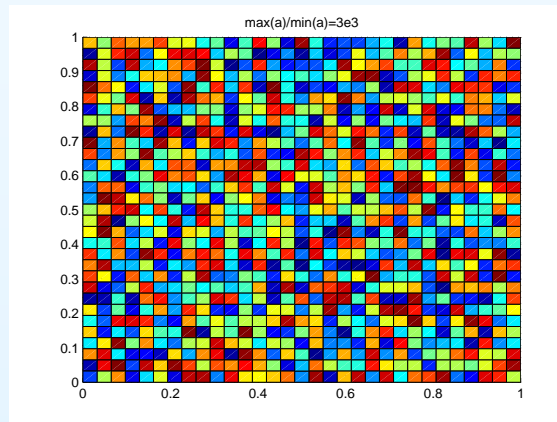
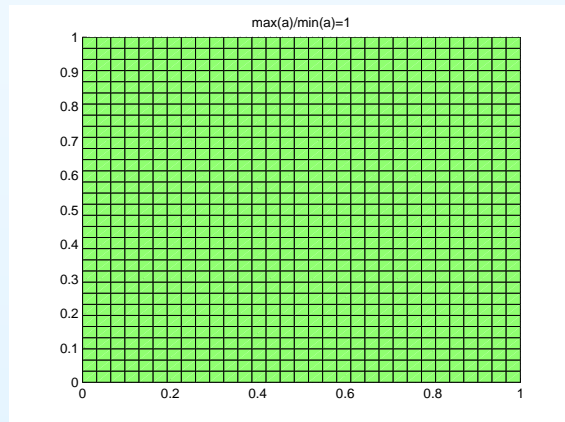
where  $D_i$  are diffusions,  $\mu_i$  mobilities,  $i = n, p$ ,  $\epsilon$  dielectric func.

# The Mixed Problem in Oil Recovery

Poisson equation on mixed form:

$$\begin{cases} \frac{1}{a}\sigma - \nabla u = 0 & \text{in } \Omega \\ -\nabla \cdot \sigma = f & \text{in } \Omega \\ n \cdot \sigma = 0 & \text{on } \Gamma \end{cases}$$

where the permeability  $a$  is constant, random, or taken from the SPE data set (upperness in log-scale),



We let  $a(v, w) = (\frac{1}{a}v_1, w_1) + (v_2, \nabla \cdot w_1) + (\nabla \cdot v_1, w_2)$  and  $l(w) = -(f, w_2)$ .



## Splitting Based on RT-elements

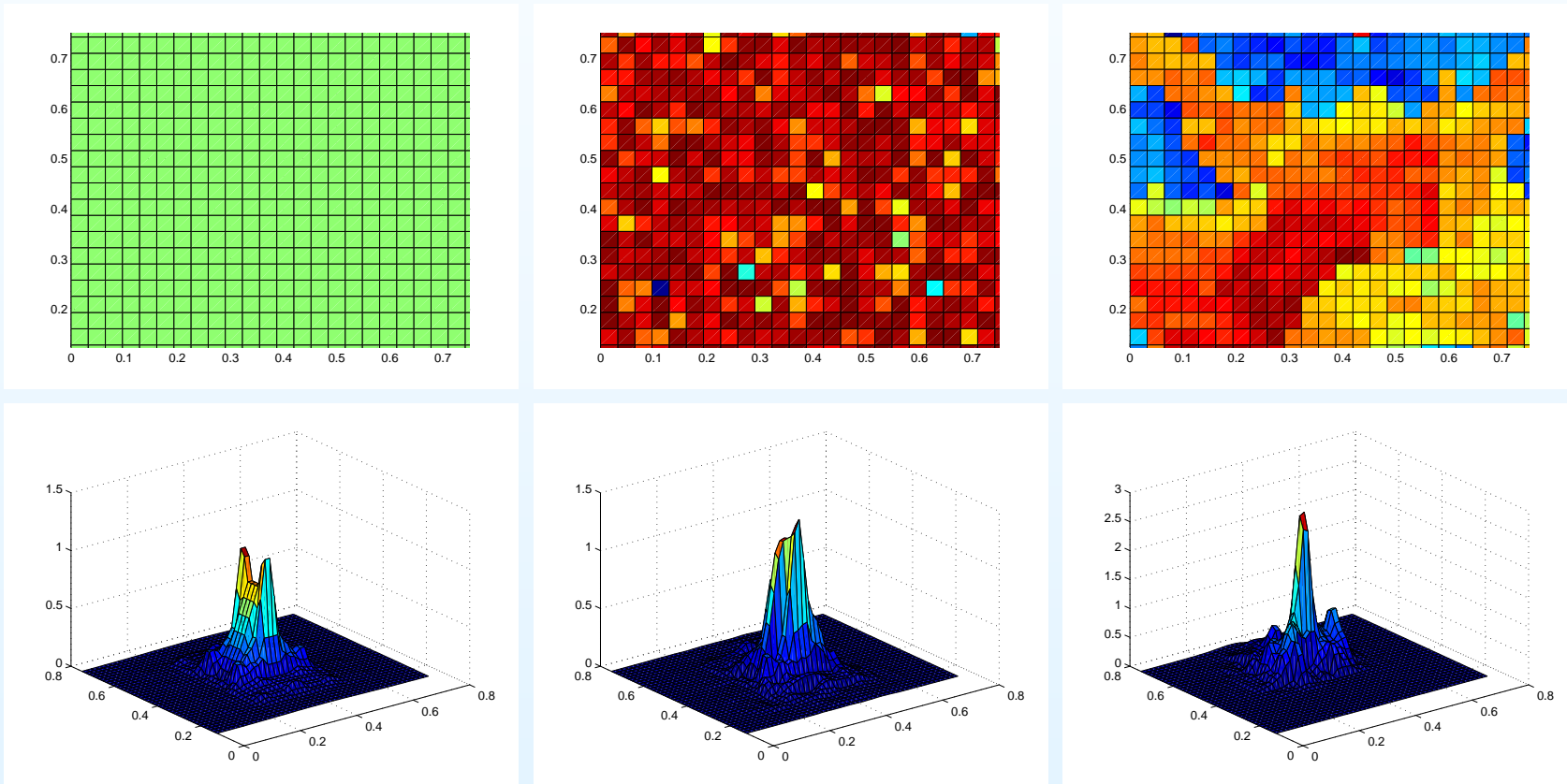
We use lowest order RT basis functions together with piecewise constants.

- Let  $\Pi_c$  be the RT-interpolant onto the space of lowest order RT functions  $\mathcal{V}_c$  and  $P_c$  be the  $L^2$ -projection onto the space of piecewise constants  $W_c$
- We define  $\mathcal{W}_f = (I - P_c)\mathcal{W}$ ,  $\mathcal{W} = L^2(\Omega)$
- We define  $\mathcal{V}_f = (I - \Pi_c)\mathcal{V}$ ,  $\mathcal{V} = H(\text{div}; \Omega)$
- This means  $\sigma_c = \pi_c \sigma \in \mathcal{V}_c$     $u_c = P_c u \in W_c$ .
- Thus we are using an  $L^2$ -orthogonal splitting in the scalar variable.

Hierarchical split for lagrangian elements leads to nodal exactness in the coarse solution  $U_c \approx \pi_c u$  while here we get exactness of average values on coarse elements  $U_c \approx P_c u$ .

## Example of Local Solutions $\tilde{\mathcal{T}}_\sigma \phi_i$

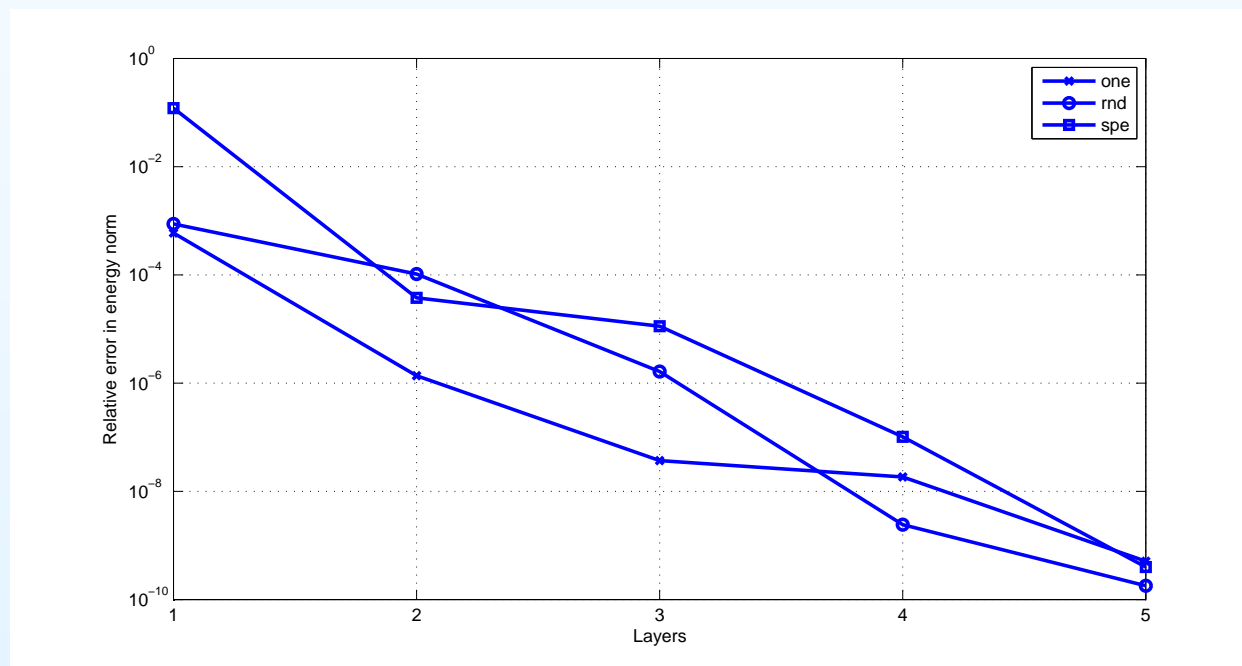
$$\begin{cases} (\frac{1}{a} \tilde{\mathcal{T}}_\sigma \phi_i, v_f) + (\tilde{\mathcal{T}}_u \phi_i, \nabla \cdot v_f) = -(\frac{1}{a} \phi_i, v_f) \\ -(\nabla \cdot \tilde{\mathcal{T}}_\sigma \phi_i, w_f) = 0. \end{cases}$$



We use 3 layer patches and plot absolute value of the flux  $|\tilde{\mathcal{T}}_\sigma \phi_i|$ .

## Example of Convergence

- The reference mesh has  $32 \times 32$  elements
- The coarse mesh has  $8 \times 8$  elements.
- We let  $f = 1$  lower left corner and  $f = -1$  in upper right, otherwise  $f = 0$ .

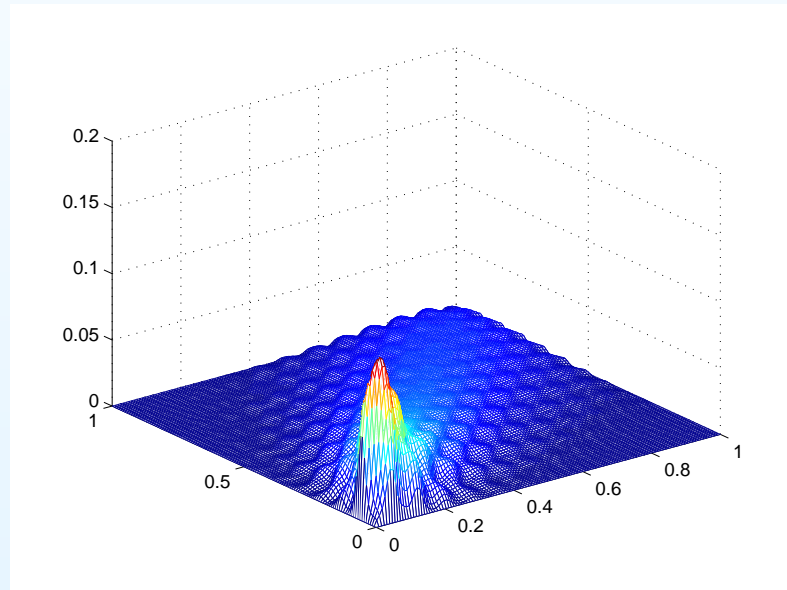
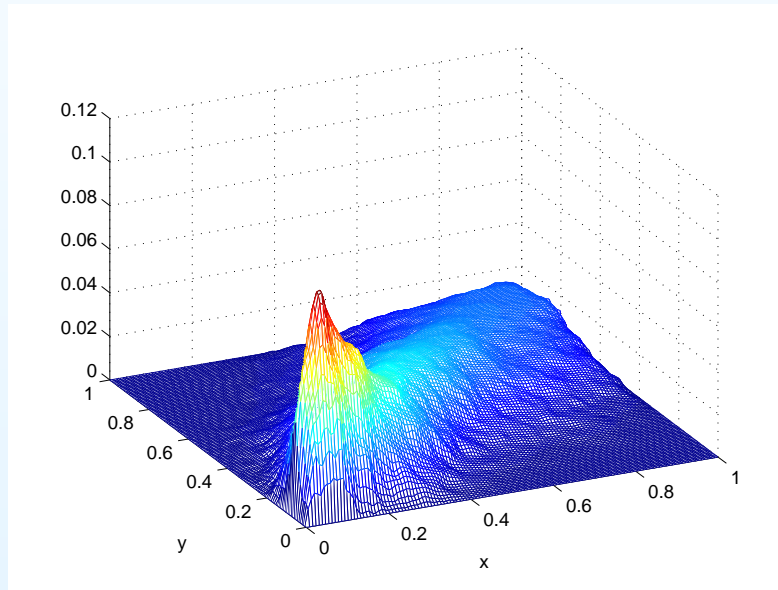


Error compared to reference solution.

## Two Convection Dominated Problems

$$\mathcal{L}u = -\epsilon\Delta u + \nabla \cdot (bu) = f \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega,$$

- i. Let  $b = [N(0, 1) + 0.25, N(0, 1) + 0.25]$  and  $\epsilon = 0.008$
- ii. Let  $b = [\sin(25\pi y) + 0.3, \cos(25\pi x) + 0.3]$ , and  $\epsilon = 0.005$ .



In both examples  $f = 1$  for  $0 \leq x, y \leq 0.1$  and zero otherwise.  
We let  $a(v, w) = (\epsilon \nabla v, \nabla w) + (\nabla \cdot (bv), w)$  and  $l(w) = (f, w)$ .

## Split Based on Lagrangian Basis Functions

- We let the coarse mesh consist of  $16 \times 16$  and the reference mesh of  $128 \times 128$  rectangular elements.
- We let  $\pi_c$  be the interpolant onto bilinear functions of the coarse mesh and let  $\mathcal{V}_f = (1 - \pi_c)\mathcal{V}$ , i.e. an hierarchical split
- We remember the definition of the fine scale equations,

$$a(\mathcal{T}\phi_i, v_f) = -a(\phi_i, v_f),$$

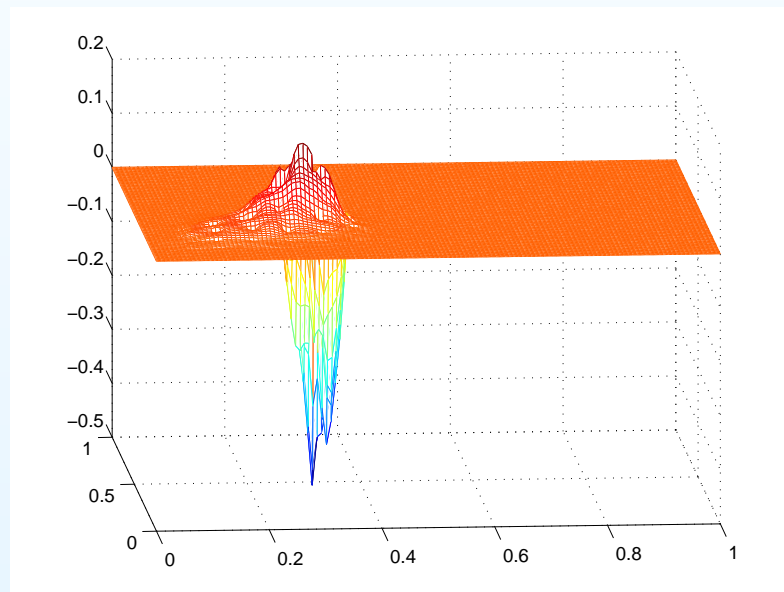
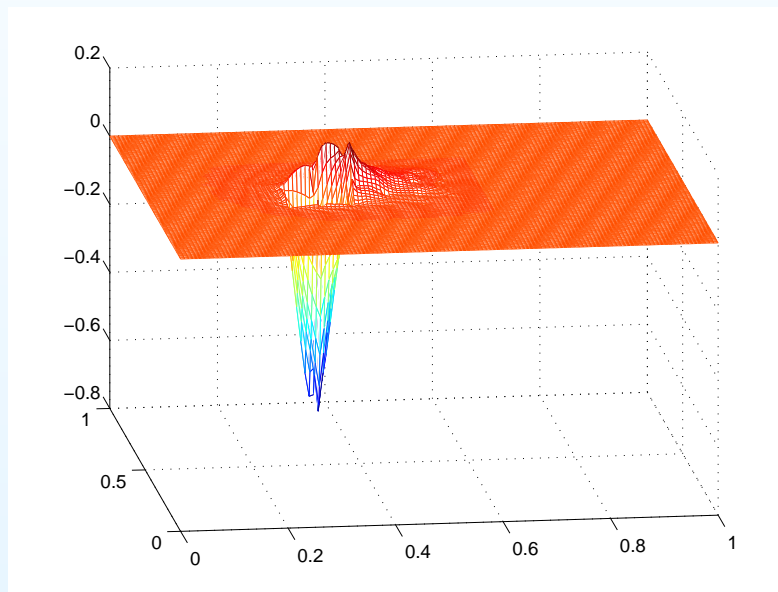
$$a(v_f, \mathcal{T}^*\phi_i) = -a(v_f, \phi_i),$$

$$a(u_{f,i}, v_f) = l(\phi_i v_f),$$

for all  $v_f \in \mathcal{V}_f$  and  $i \in \mathcal{N}$  (coarse nodes).

## Decay of Local Solutions

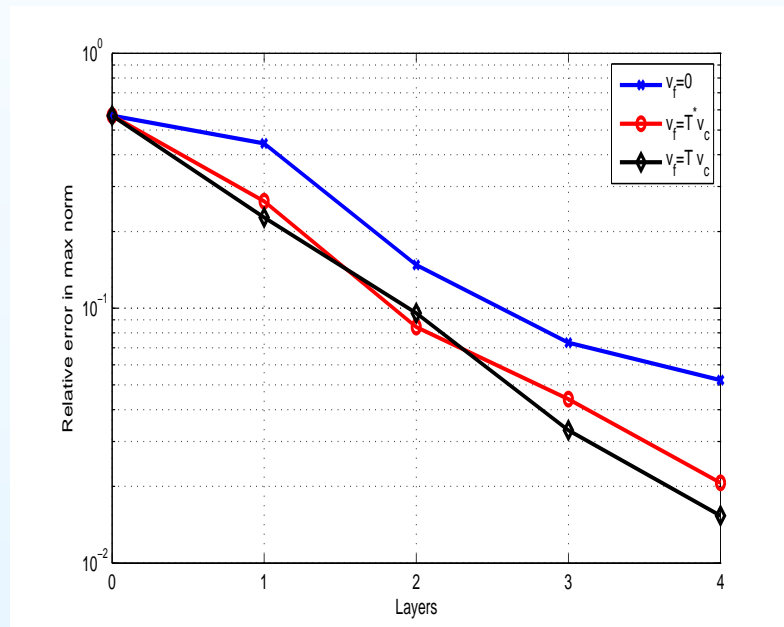
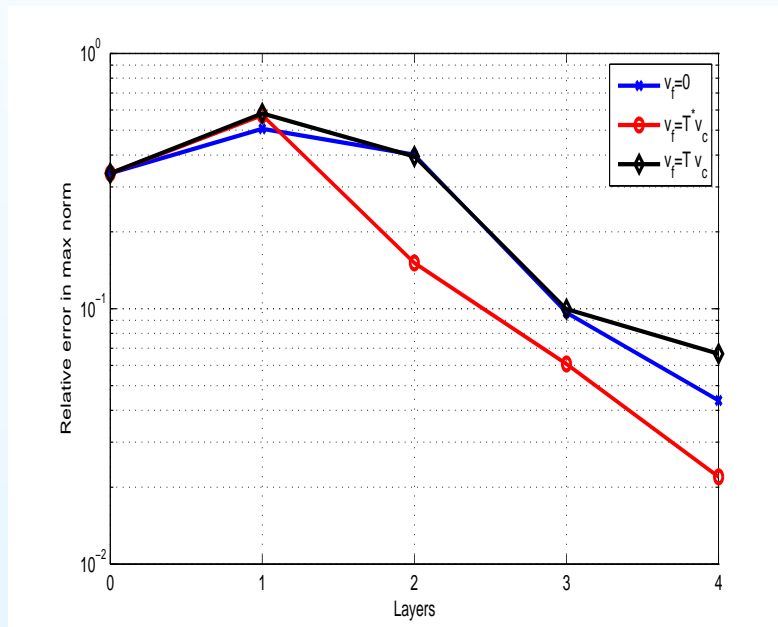
We use three refinements in all local problems and vary the number of layers. We plot local solutions  $\tilde{\mathcal{T}}\phi_i$  and  $\tilde{\mathcal{T}}^*\phi_i$  using four layers (random).



We still get decay but not as quick as in the non-convective problem. Directed patches would reduce the work.

# Error in Multiscale Solution

We now plot the max norm of the error  $\|u - U\|_{L^\infty(\Omega)}$  with and without  $\tilde{\mathcal{T}}^*$  in the right slot.



We observe quick decay (slower) starting from two layers.

The random coefficient appears to be more difficult to resolve.

# Adaptive Multiscale Method

The adaptive version of the method (sometimes referred to as AVMS) builds on the following ingredients:

- Error estimation framework
- Adaptive strategy for tuning of critical discretization parameters

The method is designed so that:

$$\text{error} \rightarrow 0 \text{ when } h \rightarrow 0 \text{ and } L \rightarrow \infty$$

- A priori error estimates in progress.
- To circumvent difficulties with choosing discretization parameters  $h$  and  $L$  we use an adaptive algorithm based on a posteriori error estimates



## A Posteriori Error Estimate (Poisson, mixed)

The following energy norm bound holds

$$\left\| \frac{1}{\sqrt{a}} (\sigma - \Sigma) \right\|^2 \leq C_a \sum_i (R_{\omega_i}^2 + R_{\partial\omega_i}^2)$$

where

$$R_{\omega_i}^2 = \left\| \frac{1}{a} (\Sigma_c^i(\phi_i + \tilde{\mathcal{T}}_\sigma \phi_i) + \Sigma_f^i) - \nabla U_f^{i,*} \right\|_{\omega_i}^2 + \left\| \frac{h}{a} (f\psi_i + \nabla \cdot (\Sigma_c^i(\phi_i + \tilde{\mathcal{T}}_\sigma \phi_i) + \Sigma_f^i)) \right\|_{\omega_i}^2 + \sum_{K \in \omega_i} \left\| h^{-1/2} [U_f^{i,*}] \right\|_{\partial K}^2$$

$$R_{\partial\omega_i}^2 = \left\| h^{-1/2} U_f^{i,*} \right\|_{\partial\omega_i \setminus \Gamma}^2$$

$U^*$  is a post processed version (Lovadina and Stenberg 06)

of  $U$ ,  $C_a \sim \|\sqrt{a}\|_{L^\infty(\omega_i)}$ .

# Adaptive Algorithm

We have the error bound

$$\|\sigma - \Sigma\|_a^2 \leq C_a \sum_i (R_{\omega_i}^2 + R_{\partial\omega_i}^2)$$

Let  $h = H/2$  and  $L = 1$  for all  $i$ .

**for**  $i = 1, \dots, n$  **do**

    Compute  $\tilde{T}\phi_i$ ,  $\tilde{T}^*\phi_i$  (if non-symmetric), and  $U_{f,i}$ .

**end for**

Compute the solution  $\{U, \Sigma\}$ .

**for**  $i = 1, \dots, n$  **do**

    Compute residuals  $R_{\omega_i}^2$  and  $R_{\partial\omega_i}^2$ .

**end for**

Mark large entries.

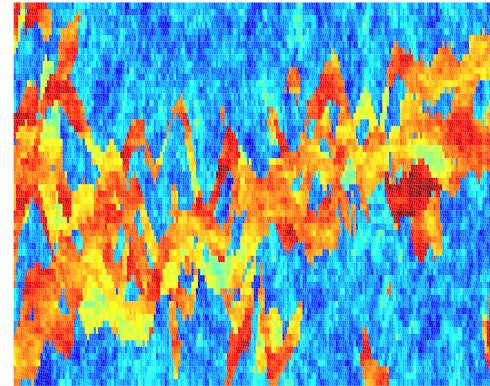
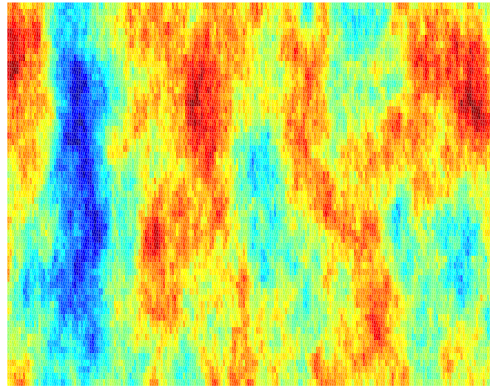
For marked entries in  $R_{\omega_i}^2$  let  $h := h/2$ .

For marked entries in  $R_{\partial\omega_i}^2$  let  $L := L + 1$ .

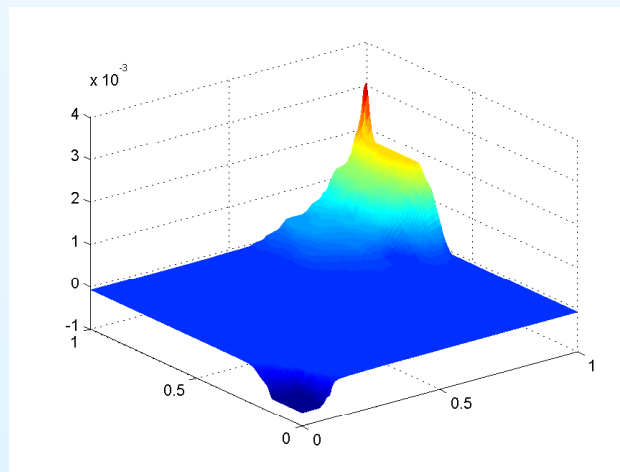
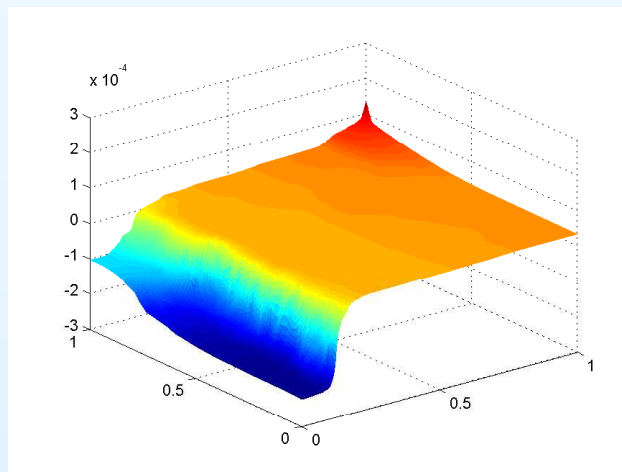
Return to step two or stop if estimators are small enough.

## Back to the Oil Recovery Example

Layer 1 and 50 in the SPE comparative sol. proj. (log scale).

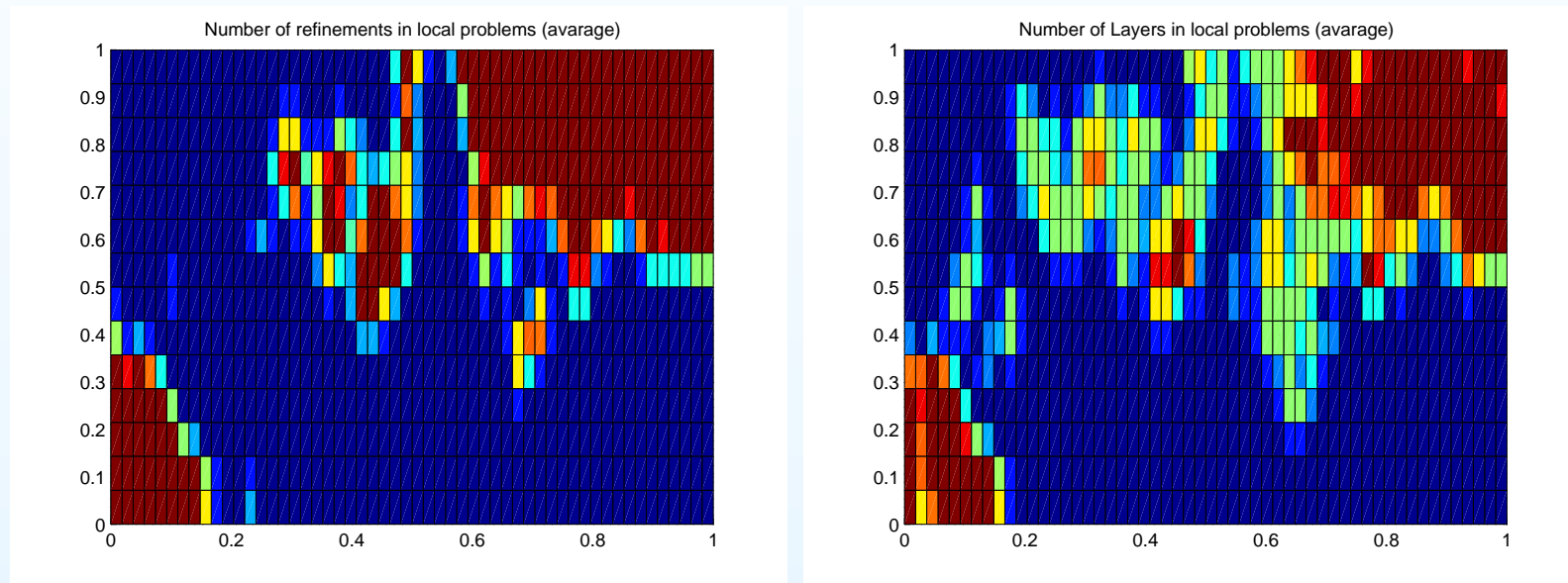


Plot of the sol. (pressure),  $q = 1$  upper right  $q = -1$  lower left.



## Refinements and Layers SPE50

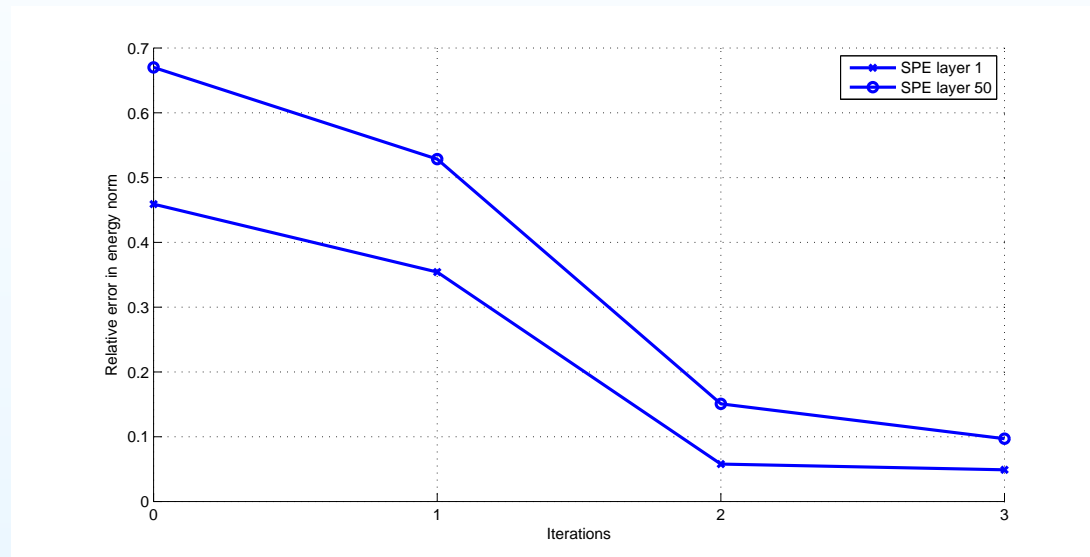
We use  $55 \times 15$  coarse elements and a reference mesh with  $440 \times 120$  elements.



We start the adaptive algorithm with **one refinement** and **one layer** in all local problems. After three iterations in the algorithm marking 30%.

# Convergence of Adaptive Algorithm

We compare error in energy norm with reference solution.



- Critical areas are found
- A majority of the patches uses one layer and one refinement.
- As the water front travels only local problems at the front need to be recomputed.

## Summary and Future Work

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The adaptive variational multiscale method (AVMS) provides:

- Systematic technique for construction of a computable approximation of the fine scale part of the solution using decoupled localized subgrid problems.
- A posteriori error estimation framework
- Adaptive algorithms for automatic tuning of critical discretization parameters

The decay in  $\mathcal{V}_f$  together with the adaptive strategy makes the method efficient.

Future work includes: 3D implementation, a priori error analysis, time-dependent convection dominated problems, other applications