

*Progress in the SSF project "CutFEM, Geometry, and Optimal design"*

# Computer simulation of multiscale problems

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University of Gothenburg and Uppsala University

Umeå 2015-05-20

# Research group and Outline

## **My research group consists of:**

Daniel Elfverson, Uppsala, Multiscale and UQ, 2011-2015

Fredrik Hellman, Uppsala, Multiscale and UQ, 2012-2017

Anna Persson, Göteborg, Time dep. ms problems, 2013-2018

Gustav Kettil, Göteborg, Simulation of paper, 2014-2016

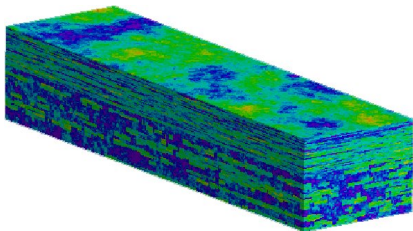
Tony Stillfjord, Göteborg, Time dep. PDE and splitting, 2015-2017

*Support from SSF project 2016-2018.*

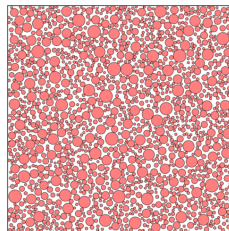
- Introduction to Localized Orthogonal Decomposition techniques for solving multiscale problems (Axel)
- Application to eigenvalue problems (Axel)
- Application to problems posed on complex geometry (Daniel)
- Recent development and openings for collaboration (Axel)

# Multiscale problems

Applications such as



▷ flow in a porous medium



▷ composite materials

require numerical solution of partial differential equations with rough data (e.g. permeability or module of elasticity).

Major challenge: Features on **multiple scales** in space and time.

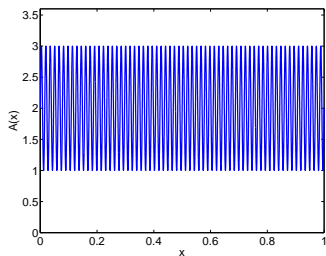
# Motivation for multiscale techniques

Poisson's equation

$$-\nabla \cdot \mathbf{A} \nabla u = f \quad \text{in } \Omega \quad u = 0 \quad \text{on } \partial\Omega$$

where  $\mathbf{A}$  has rapid oscillations.

**Example** (periodic coefficient):  $A(x) = 2 + \sin(2\pi x/\varepsilon)$ ,  $\varepsilon = 2^{-6}$ ,  $f = 1$



oscillatory coefficient

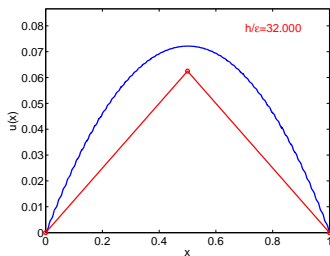
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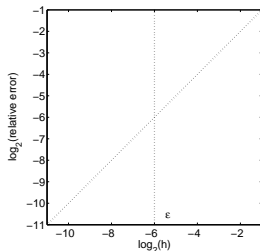
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solution and P1-FEM-approximation



$\log_2(H^1(\Omega) - \text{error})$  vs.  $\log_2(h)$

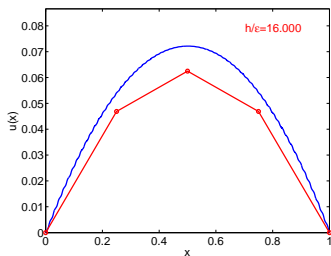
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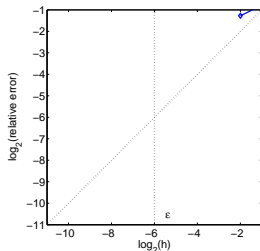
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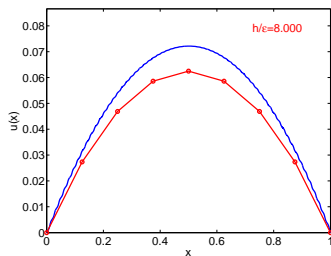
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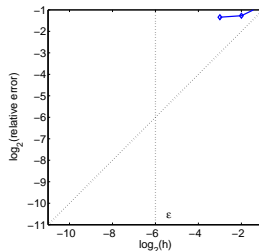
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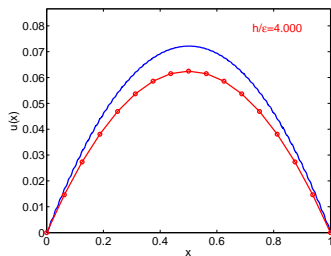
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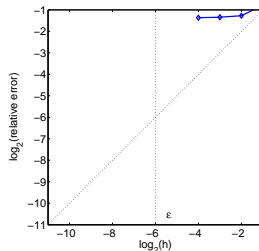
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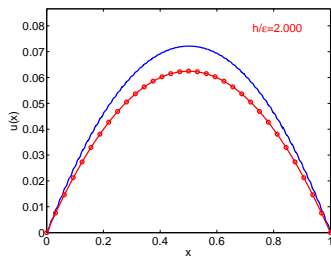
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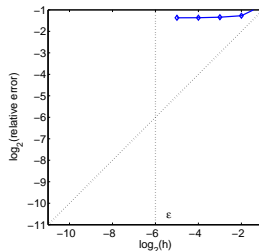
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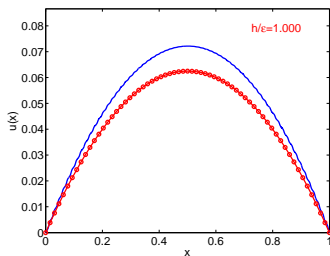
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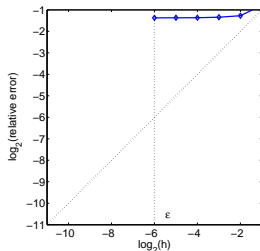
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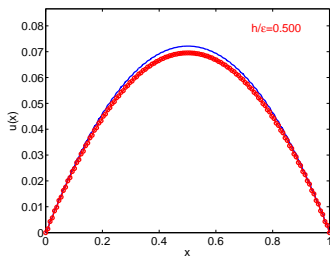
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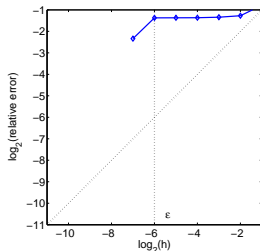
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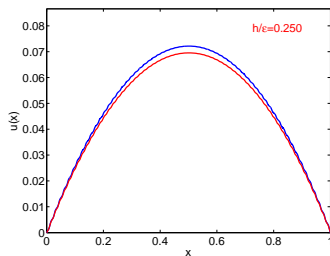
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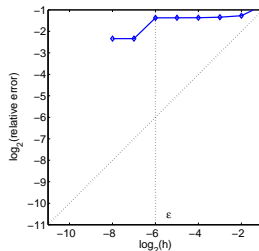
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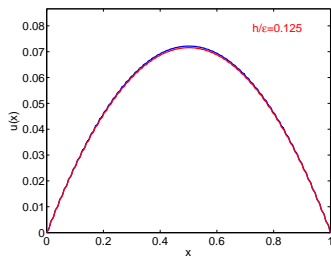
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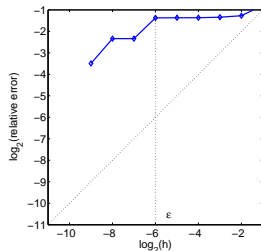
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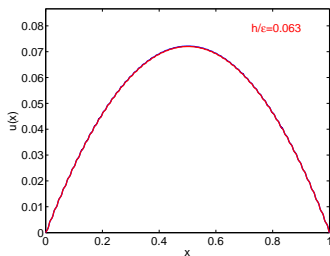
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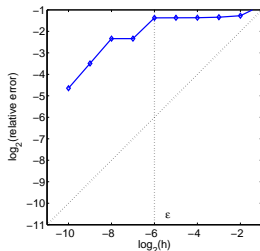
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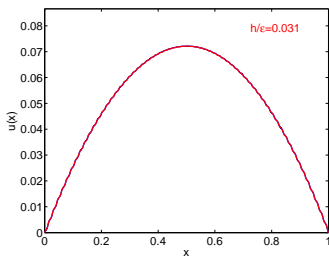
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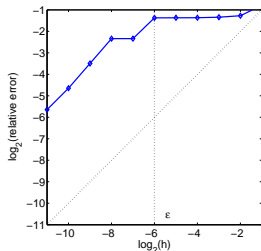
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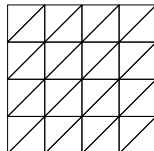
## Conclusion

- Fine scale features have to be resolved even to get coarse solution behavior right (both  $H^1$  and  $L^2$  errors are large).
- Resolution of the fine scales by a uniform mesh is very computationally expensive.
- Local mesh refinement is not an option.
- **The standard basis does not seem to be suitable for this problem.**



# Standard FE decomposition

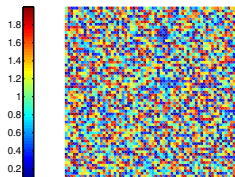
- Coarse FE mesh with parameter  $H$
- Piecewise linear continuous FE space  $V_H$
- $\mathfrak{I}_{\mathcal{T}} : V \rightarrow V_H$  interpolation operator



## Decomposition

$$V = V_H \oplus V^f \quad \text{with } V^f := \{v \in V \mid \mathfrak{I}_{\mathcal{T}} v = 0\}$$

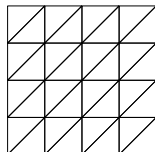
## Example:



rough coefficient

# Standard FE decomposition

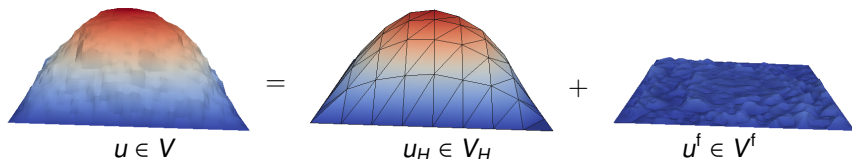
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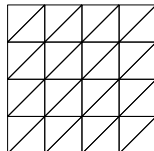
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## Example:



# Orthogonal multiscale decomposition

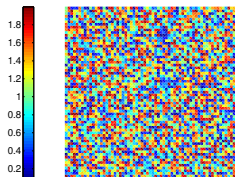
- Let  $\dim(V_H^{\text{ms}}) = \dim(V_H)$
- Start from  $V_H$  and add fine scale corrections in  $V^f$
- $\forall v_H^{\text{ms}} \in V_H^{\text{ms}}, v_f \in V^f$  holds  $\int_{\Omega} A \nabla v_H^{\text{ms}} \cdot \nabla v_f dx = 0$



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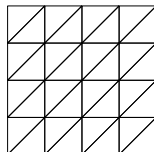
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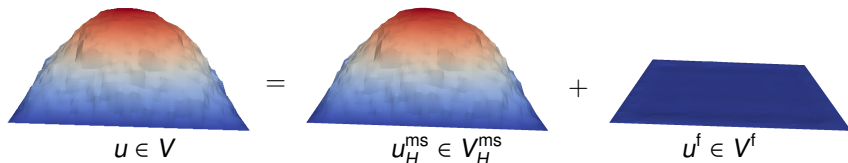
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## Decomposition

$$V = V_H^{\text{ms}} \oplus V^f \quad \text{with } V^f := \{v \in V \mid \mathfrak{I}_{\mathcal{T}} v = 0\}$$

## Example:



# Computing a basis

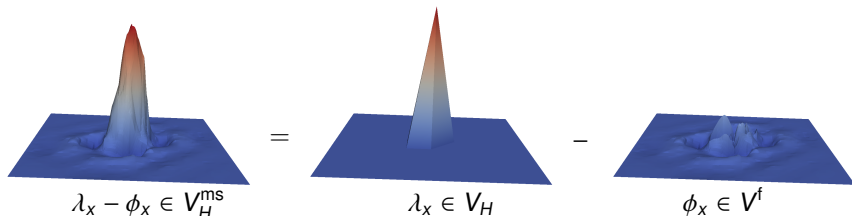
Let  $V_H = \text{span} \{\lambda_x \mid x \in \mathcal{N}\}$  and further let,  $\phi_x \in V^f$  solve

$$\int_{\Omega} A \nabla(\lambda_x - \phi_x) \cdot \nabla w \, dx = 0, \quad \text{for all } w \in V^f.$$

## Multiscale FE space

$$V_H^{\text{ms}} = \text{span} \{\lambda_x - \phi_x \mid x \in \mathcal{N}\}$$

Example:



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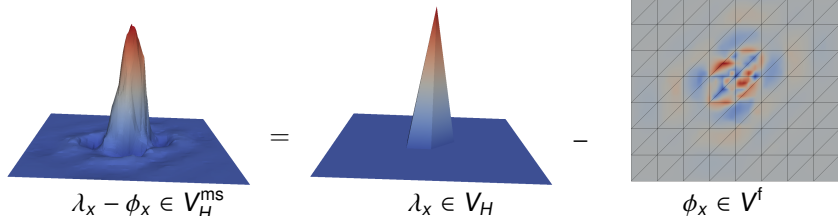
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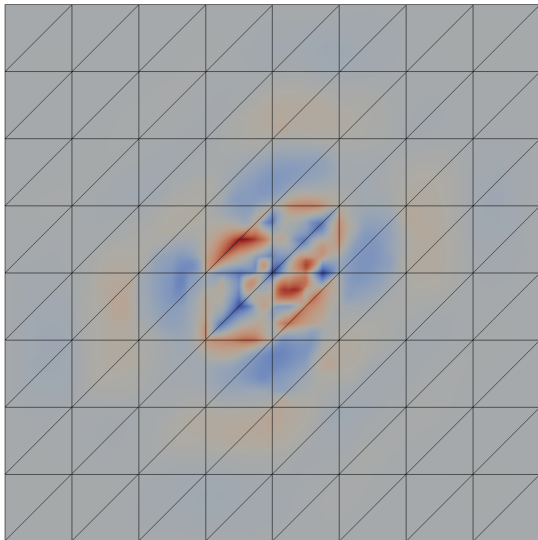
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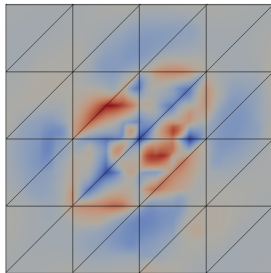


# Computing a basis



We have proven exponential decay of  $\lambda_x - \phi_x!$

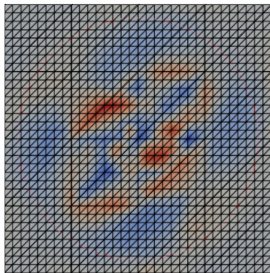
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This allows us to truncate to a patch.



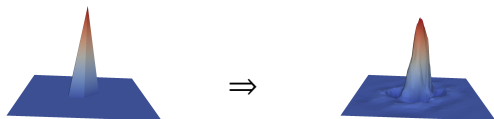
# Computing a basis



This allows us to truncate to a patch and fine scale discretization.

# Computing a basis

Localized patch with refined mesh gives computable basis functions



- The multiscale basis functions are computed by solving the PDE on local patches.
- The basis functions are totally independent, leading to trivial parallelization.
- Exponential decay has been proven and is crucial.
- Remember  $\dim(V_H^{\text{ms}}) = \dim(V_H)$ .

# Computing the multiscale approximation

**Multiscale approximation:**  $u_H^{\text{ms}} \in V_H^{\text{ms}}$  satisfies

$$\int_{\Omega} A \nabla u_H^{\text{ms}} \cdot \nabla v \, dx = \int_{\Omega} f \cdot v \, dx \quad \text{for all } v \in V_H^{\text{ms}}$$

- We have proven error bound (using  $k = \log(1/H)$  layers):

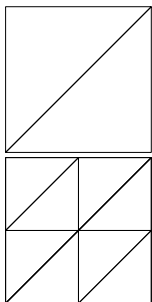
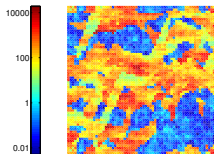
$$\|\nabla(u - u_H^{\text{ms}})\| \leq CH\|f\|,$$

where  $\|v\|^2 = \int_{\Omega} v^2 \, dx$  and  $C$  is independent on variations in  $A$ .

- Note that for the standard FEM with  $A = A(\frac{x}{\epsilon})$  we have,

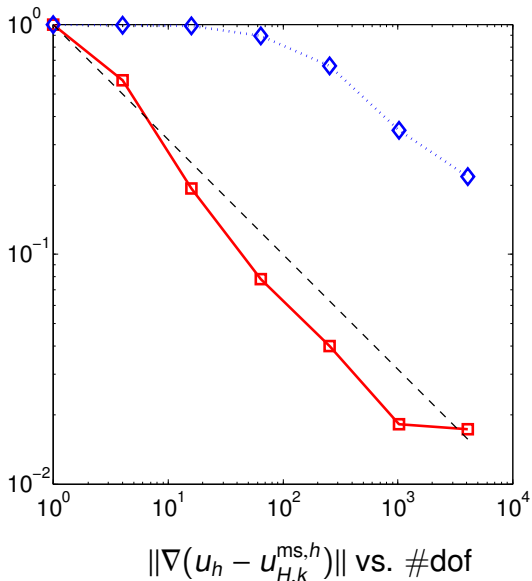
$$\|\nabla(u - u_H)\| \leq C \frac{H}{\epsilon} \|f\|.$$

# Numer. exp. (Poisson, $f \in L^2(H)$ , $f \in H_0^1(H^2)$ !)

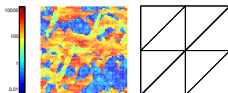


$$H = 2^{-1}, 2^{-2}, \dots, 2^{-7}$$

$$h = 2^{-9}, k = \log(1/H)$$



# Numer. exp. (linear eigenvalue problem, $H^4$ )



$\ell$	$\lambda_h^{(\ell)}$	$e^{(\ell)}(1/2\sqrt{2})$	$e^{(\ell)}(1/4\sqrt{2})$	$e^{(\ell)}(1/8\sqrt{2})$	$e^{(\ell)}(1/16\sqrt{2})$
1	21.4144522	5.472755371	0.237181706	0.010328293	0.000781683
2	40.9134676	-	0.649080539	0.032761482	0.002447049
3	44.1561133	-	1.687388874	0.097540102	0.004131422
4	60.8278691	-	1.648439518	0.028076168	0.002079812
5	65.6962136	-	2.071005692	0.247424446	0.006569640
6	70.1273082	-	4.265936007	0.232458016	0.016551520
7	82.2960238	-	3.632888104	0.355050163	0.013987920
8	92.8677605	-	6.850048057	0.377881216	0.049841235
9	99.6061234	-	10.305084010	0.469770376	0.026027378
10	109.1543283	-	-	0.476741452	0.005606426
11	129.3741945	-	-	0.505888044	0.062382302
12	138.2164330	-	-	0.554736550	0.039487317
13	141.5464639	-	-	0.540480876	0.043935515
14	145.7469718	-	-	0.765411709	0.034249528
15	152.6283573	-	-	0.712383825	0.024716759
16	155.2965039	-	-	0.761104705	0.026228034
17	158.2610708	-	-	0.749058367	0.091826207
18	164.1452194	-	-	0.840736127	0.118353184
19	171.1756923	-	-	0.946719951	0.111314058
20	179.3917590	-	-	0.928617606	0.119627862

Table: Errors  $e^{(\ell)}(H) =: \frac{\lambda_H^{(\ell)} - \lambda_h^{(\ell)}}{\lambda_h^{(\ell)}}$  and  $h = 2^{-7} \sqrt{2}$ .

# The quadratic eigenvalue problem

Consider discretized structure with damping,

$$Kx + \lambda Cx + \lambda^2 Mx = 0,$$

where  $K_{ij} = (A \nabla \phi_j, \nabla \phi_i)$  is stiffness,  $C_{ij} = c(\phi_j, \phi_i)$  is damping, and  $M_{ij} = (\phi_j, \phi_i)$  is mass matrix.

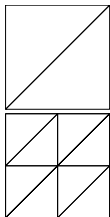
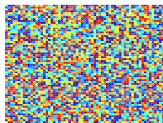
**Linearization:** ( $y = \lambda x$ )

$$Ax := \begin{bmatrix} K & 0 \\ 0 & M \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} = \lambda \begin{bmatrix} -C & -M \\ M & 0 \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} := \lambda Bx,$$

Note that  $T = A^{-1}B$  has eigenvalues  $\lambda^{-1}$ .  $T$  is not symmetric (even if  $C$  is), it is compact but not necessarily in the limit  $h \rightarrow 0$ .

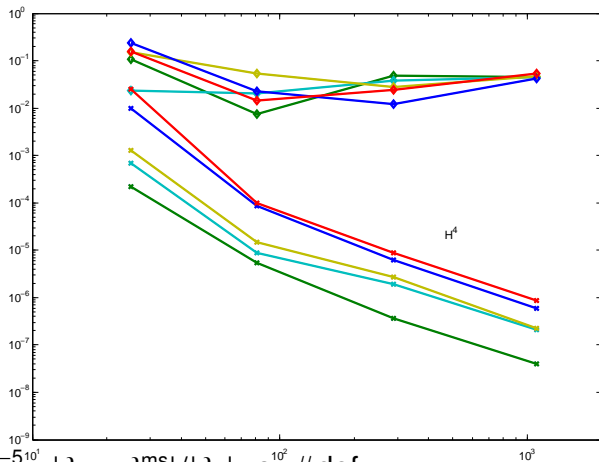
The analysis is very different from symmetric case. Semi-simple eigenvalues are ok but defective leads to unknown constants.

# Numerical experiment



$$H = 2^{-1}, 2^{-2}, \dots, 2^{-5 \cdot 10^1}$$

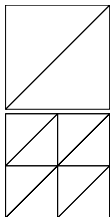
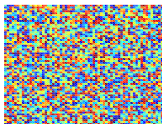
$$h = 2^{-6}, k = \infty$$



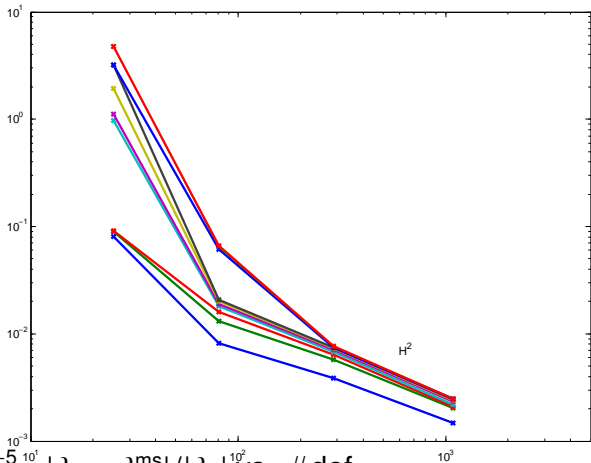
$|\lambda_h - \lambda_H^{ms}| / |\lambda_h|$  vs.  $\#dof$

$$A \text{ (pic)}, c(u, v) = \int_{\Omega} (1 + \sin(10x)) u v \, dx$$

# Numerical experiment



$H = 2^{-1}, 2^{-2}, \dots, 2^{-5}$   
 $h = 2^{-6}, k = \infty$



$$\pi c(u, v) = \int_{\Omega} \left( \text{atan}(10x - 5) + \frac{\pi}{2} \right) A \nabla u \nabla v \, dx$$



# Relevant papers

Some papers relevant for the SSF project (work package D):

- Poisson P1, M. & Peterseim, Math. Comp., 2014.
- Poisson DG, E., Georgoulis, M., & Petersiem, SINUM, 2013.
- Linear Eig., M. & Peterseim, Numer. Math. 2015.
- Gross-Pitaevskii, Henning, M., & Peterseim, SINUM, 2014.
- Helmholtz, Peterseim, two preprints

**LOD is applicable when:** The diffusion is present and lower order terms are not dominating.

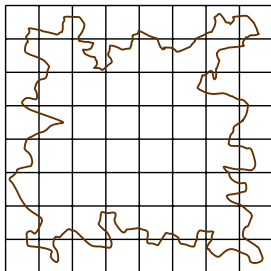
**LOD is efficient when:** basis can be reused many times in the calculation (nonlinear iteration, time stepping, ...).

# LOD – complex geometry

Consider the Poisons equation

$$\begin{aligned} -\Delta u &= f && \text{in } \Omega, \\ \nu \cdot \nabla u &= -\kappa u && \text{on } \Gamma_R, \\ u &= 0 && \text{on } \Gamma_D, \end{aligned}$$

where  $\Gamma_D$  and  $\Gamma_R$  are not resolved by the mesh



Construct a multiscale method for complicated geometries with the following properties:

- No pre-asymptotic effect
- Linear convergence without resolving the boundary
- Conditioning is independent of how the boundary cut the mesh
- Only enrich the coarse finite element space in the vicinity of the boundary

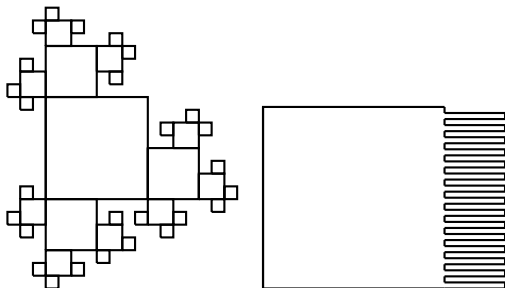
# Poincaré-Friedrich inequality

Difficulties: We need a inequality of the type

$$\inf_{c \in \mathbb{R}} \|u - c\|_{L^2(\omega)} \leq C(\omega) \text{diam}(\omega) \|\nabla u\|_{L^2(\omega)}$$

where

- $C(\omega)$  is bounded independent of  $\text{diam}(\omega)$



# Error analysis

- Let  $\omega_\Gamma$  be a vicinity of the complicated boundary

## Theorem (Locally enriched LOD method)

Given that  $u \in H_{\Gamma_D}^1 \cap H^2(\Omega \setminus \omega_\Gamma)$  and that  $u_H^{\Gamma,L} \in V_H^{\Gamma,L}$  is the LOD solution, then

$$\| \| u - u_H^{LOD} \| \|_h \lesssim \|\nabla(u - \mathfrak{I}_h u)\|_{L^2(\omega_\Gamma)} + H \|u\|_{H^2(\Omega \setminus \omega_\Gamma)} \\ + H \|f\|_{L^2(\omega_\Gamma)} + H^{-1} (L)^{d/2} \gamma^L \|f\|_{L^2(\Omega)}$$

holds. The condition number scale as  $\kappa \lesssim H^{-2}$ .

- $\mathfrak{I}_h u$ : interpolation operator on the fine mesh
- $L$ : Number of layers
- $0 < \gamma < 1$ : depends on the Poinaré-Friedrich constant

# Error analysis

- Choose  $L = \lceil C \log(H^{-1}) \rceil$ ,

## Theorem (Locally enriched LOD method)

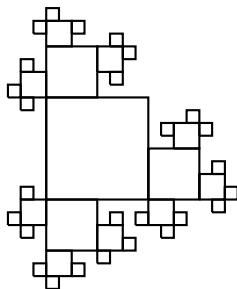
Given that  $u \in H_{\Gamma_D}^1 \cap H^2(\Omega \setminus \omega_{\Gamma}^{k-1})$  and  $u_H^{\Gamma, L} \in V_H^{\Gamma, L}$ , then

$$\| \| u - u_H^{LOD} \| \|_h \lesssim H$$

holds. The condition number scale as  $\kappa \lesssim H^{-2}$ .

# Numerical verification – Fractal domain

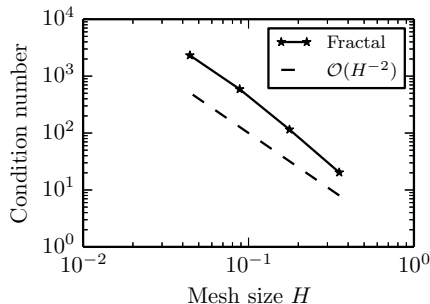
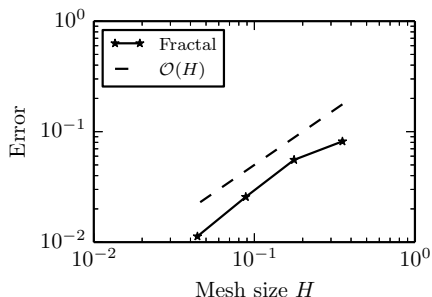
Consider the Poisson equation in the domain:



- Correctors are computed everywhere
- Mixed boundary condition

# Numerical verification – Fractal domain

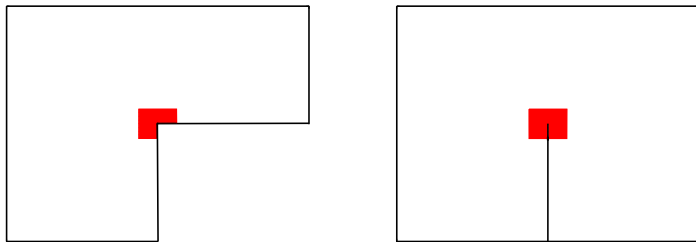
- The convergence in relative energy norm
- Conditioning





# Numerical verification – Local singularities

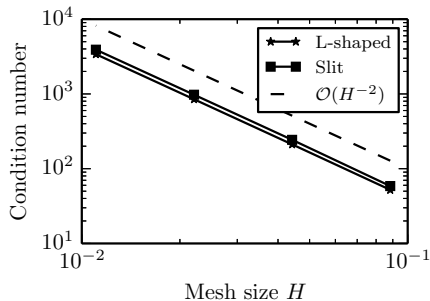
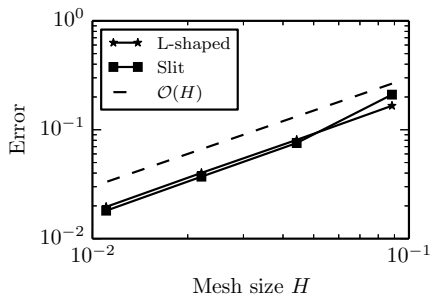
Consider the Poisson equation in the domains:



- Correctors are computed in the red areas
- Dirichlet boundary condition

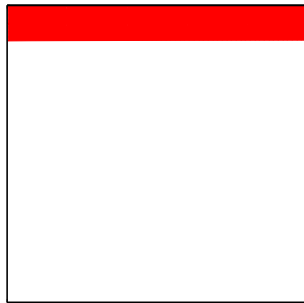
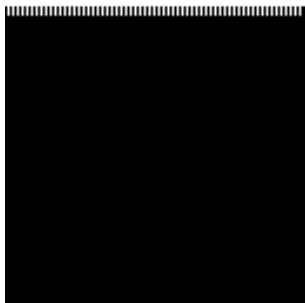
# Numerical verification – Local singularities

- The convergence in relative energy norm
- Conditioning



# Numerical verification – Saw tooth boundary

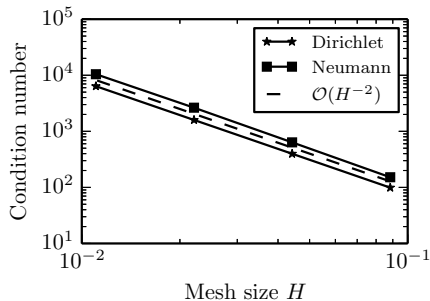
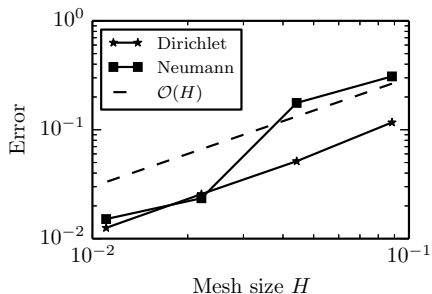
Consider the Poisson equation in the domain:



- Correctors are computed in the red area
- Test both Dirichlet and Neumann boundary condition on the saw tooth

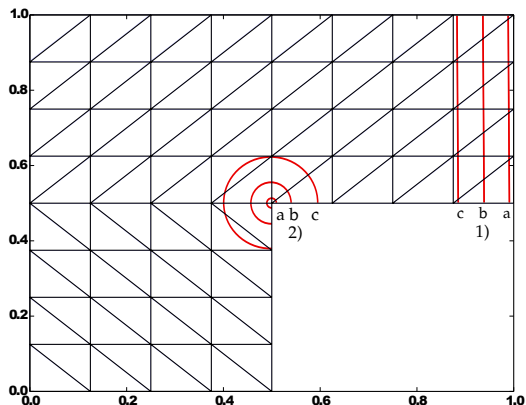
# Numerical verification – Saw tooth boundary

- The convergence in relative energy norm
- Conditioning



# Numerical verification – Cut domain

Consider the Poisson equation in the domain:



- Correctors are computed everywhere
- Consider a fix mesh and cut the domain in different ways

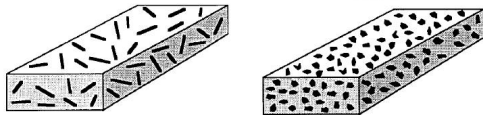
# Numerical verification – Cut domain

Cut 1	$e_{\text{rel}}(a)$	$e_{\text{rel}}(b)$	$e_{\text{rel}}(c)$
D	0.085	0.082	0.082
D-N	0.092	0.078	0.077
N-D	0.036	0.036	0.035
	cond( $a$ )	cond( $b$ )	cond( $c$ )
D	10.37	10.59	11.47
D-N	10.82	11.17	12.06
N-D	272.56	254.58	256.48

Cut 2	$e_{\text{rel}}(a)$	$e_{\text{rel}}(b)$	$e_{\text{rel}}(c)$
D	0.086	0.091	0.100
D-N	0.086	0.081	0.080
N-D	0.037	0.052	0.068
	cond( $a$ )	cond( $b$ )	cond( $c$ )
D	10.32	10.08	9.49
D-N	10.41	11.98	13.56
N-D	222.78	99.95	71.06

# Recent development and collaboration

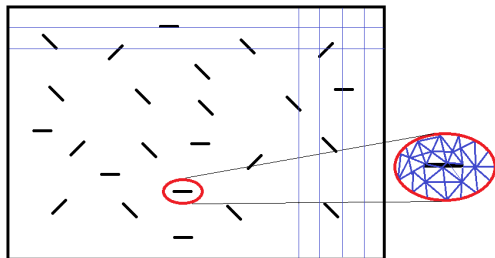


We consider a Poisson type model problem with a diffusion coefficient  $A$  that varies between two different values:

$$-\nabla \cdot A \nabla u = f \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega.$$

- The objects may be short fibres/particles or cracks/holes
- the solution has low regularity at the boundary of the objects
- local mesh refinement work for single configuration but not when multiple configurations or "moving" objects are allowed
- Monte Carlo simulation or optimization procedure

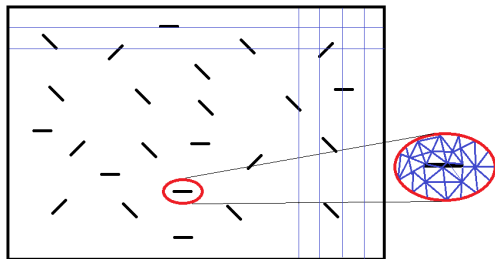
# Recent development



- In order to treat multiple configurations we want a regular static coarse background mesh.
- We only compute LOD basis functions in a region around each object and use standard P1 FEM basis functions in between.
- The size of the surrounding region can be controlled by  $H$  to meet tolerance,  $p$  adaptivity, DG. .
- If the regions overlap we compute basis functions for the union of the overlapping domains.

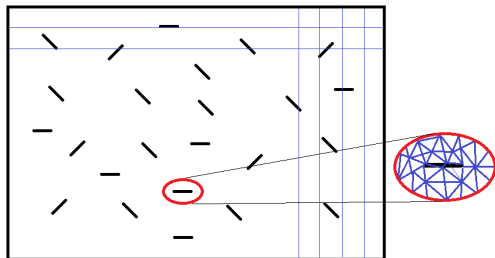


# Recent development



- CutFEM allows us (we hope) to "glue" the LOD computations to the background mesh avoiding remeshing.
- The assembly of the LOD parts is independent of the orientation of the object (rotation, translation).
- If the objects have the same shape huge savings are possible.
- Compared to CutFEM using spatially adapted subgrids around the objects the LOD approach leads to much smaller systems of equations for each configuration of the objects.

# Recent development



- Random distribution of objects: Monte-Carlo simulation can be performed on the coarse scale with "fine scale" accuracy.
- Design of materials: Optimization procedure where the objects are moved in order to minimize some output. Again repeated solves on the coarse scale still gives "fine" scale accuracy.

# Comments/questions/possible collaborations

- The problem is very similar to complex boundary, holes is a direct application.
- Any comments from or difficulties seen by the CutFEM experts?
- We got the idea from porous media flow applications, where there is a great uncertainty in the location and direction of defects. We seek applications in composite materials which fits our framework.
- Uncertainty quantification in material science is a possible application for this approach.
- What about applicability to optimal design problems?

We will start this work after the summer, in collaboration with Fredrik Hellman