

Fakultät für Mathematik Institut für Angewandte und Numerische Mathematik

Geometric Numerical Integration, Serie 6

Exercise 18: Complete the proof of the following Theorem from the lecture.

Let f(y) be analytic in $B_{2R}(y_0)$, let the coefficients $d_j(y)$ of the numerical scheme be analytic in $B_R(y_0)$ and assume that

$$||f(y)|| \le M$$
 for $||y - y_0|| \le 2R$

and

$$\|d_j(y)\| \le \mu M \Big(\frac{2\kappa M}{R}\Big)^{j-1}$$
 for $\|y-y_0\| \le R$

hold. If $h \le h_0/4$ with $h_0 = R/(e\eta M)$, then there exists N = N(h) (namely N equals the largest integer satisfying $hN \le h_0 < h(N+1)$) such that the difference between the numerical solution $y_1 = \Phi_h(y_0)$ and the exact solution $\varphi_{N,t}(y_0)$ of the truncated modified equation satisfies

$$\|\Phi_h(y_0) - \varphi_{N,h}(y_0)\| \le h\gamma M e^{-h_0/h}$$

where $\gamma = e(2 + 1.65\eta + \mu)$ depends only on the method.

The proof proceeded in the following steps, where all parts except for part (d) were done in the lecture.

(a) Show that for $g(h) := \Phi_h(y_0) - \varphi_{N,h}(y_0)$ the bound

$$\|g(h)\| \leq \left(\frac{h}{\varepsilon}\right)^{N+1} \max_{|z| \leq \varepsilon} \|g(z)\|$$

holds for $0 \le h \le \varepsilon := eh_0/N$.

(b) Split the error in two parts,

$$||g(z)|| \le ||\Phi_z(y_0) - y_0|| + ||\varphi_{N,z}(y_0) - y_0||$$
,

which are estimated separately in the following two steps of the proof.

(c) Show that

$$\|\Phi_z(y_0) - y_0\| \le \varepsilon M(1+\mu)$$

holds.

(d) Show that

$$\|\varphi_{N,z}(y_0) - y_0\| \le \varepsilon M(1 + 1.65\eta)$$

for $\varphi_{N,z}(y_0) \in B_{R/2}(y_0)$.

(e) Finally combine both bounds to obtain the desired result.

Exercise 19:

Consider a differential equation

$$y'(t) = f(y(t)), \quad y(0) = y_0$$

which possesses the invariant I(y). We solve this equation numerically with a scheme $\Phi_h(y)$, that also conserves the invariant. Show that the modified equation conserves I(y) as well.

<u>Hint</u>: Show via induction that $\nabla I(y)f_j(y) = 0$, j = 1, ..., r holds. For this purpose let $\varphi_{r,t}^{(h)}$ be the flow of the truncated modified equation $\tilde{y}' = f(\tilde{y}) + hf_2(\tilde{y}) + ... + h^{r-1}f_r(\tilde{y})$, $\tilde{y}(0) = y_0$.

3.7.2012

Programming Exercise 7:

(a) Show that for a nonsymplectic method of order *r* applied to a Hamilton system the error in the energy grows linearly with the time, namely

$$H(y_n) - H(y_0) = \mathcal{O}(th^r)$$

where t = nh.

<u>Hint</u>: Use that the local error of the method is of the order h^{r+1} .

(b) Consider the pendulum equation given by the Hamiltonian

$$H(p,q) = \frac{p^2}{2} - \cos(q)$$

with initial condition $(p_0, q_0) = (2.5, 0)$. Solve this equation with the explicit and the symplectic Euler method on the time interval [0, T] = [0, 50] using the time step size h = 0.005. Plot the error in the energy H of both methods over time.

Programming Exercise 8:

Consider the pendulum equation given by the Hamilton function

$$H(p,q) = \frac{p^2}{2} - \cos(q)$$

(a) Let *K* be a compact subset of $\{(x, y) \in \mathbb{R}^2; |x| \le c\}$. Show that

$$||f(p,q)|| \le \sqrt{(c+2R)^2 + e^{4R}} = M$$

for $f(p,q) = -J\nabla H(p,q)$ and $||(p,q) - (p_0,q_0)|| \le 2R$ with $(p_0,q_0) \in K$.

- (b) Now choose c = 2 and R = 1/2 and compute the corresponding time step size h_0 for the midpoint rule to obtain good energy conservation.
- (c) Numerically compute the solution of the pendulum equations with the different initial conditions

$$(p_0, q_0) = (0, -1.5)$$
, $(0, -2.5)$, $(1.5, -\pi)$, $(2.5, -\pi)$.

Therefore use the midpoint rule with the time step size h_0 computed in the previous part on the time interval $[0, 200.000h_0]$ and plot the components p and q against each other. What happens to the results, if the time step size is increased?

Programming Exercise 9: Oscillatory example – Fermi-Pasta-Ulam problem

Consider the equation

$$x^{\prime\prime}(t) = -\omega^2 x(t)$$
 , $\omega \gg 1$

with initial conditions $x(0) = x_0$ and $x'(0) = x'_0$.

(a) Show, that the exact solution of the system is given by

$$\begin{pmatrix} \omega x(t) \\ x'(t) \end{pmatrix} = \begin{pmatrix} \cos \omega t & \sin \omega t \\ -\sin \omega t & \cos \omega t \end{pmatrix} \begin{pmatrix} \omega x_0 \\ x'_0 \end{pmatrix} \, .$$

(b) We consider the midpoint rule

$$y_{n+1} = y_n + hf\left(\frac{y_n + y_{n+1}}{2}\right)$$

for y' = f(y) as well as the Störmer-Verlet method

$$\begin{aligned} x'_{n+1/2} &= x'_n + \frac{h}{2}g(x_n) \\ x_{n+1} &= x_n + hx'_{n+1/2} \\ x'_{n+1} &= x'_{n+1/2} + \frac{h}{2}g(x_{n+1}) \end{aligned}$$

for x'' = g(x). Compute one step of each method applied to the above problem. Note that for the midpoint rule, the problem needs to be rewritten as a first order system. Analyze the stability of the methods with respect to the time step size *h*.

Hint: Write the numerical solution as

$$\begin{pmatrix} \omega x_{n+1} \\ x'_{n+1} \end{pmatrix} = M(h\omega) \begin{pmatrix} \omega x_n \\ x'_n \end{pmatrix}$$

and consider the eigenvalues of $M(h\omega)$.

(c) The Hamiltonian of the Fermi-Pasta-Ulam problem in the scaled expansions of the springs is given by

$$\begin{split} H(y,x) &= \frac{1}{2} \sum_{i=1}^{m} \left(y_{0i}^2 + y_{1i}^2 \right) + \frac{\omega^2}{2} \sum_{i=1}^{m} x_{1i}^2 \\ &+ \frac{1}{4} \left(x_{01} - x_{11} \right)^4 + \frac{1}{4} \sum_{i=1}^{m-1} \left(x_{0i+1} - x_{1i+1} - x_{0i} - x_{1i} \right)^4 + \frac{1}{4} \left(x_{0m} + x_{1m} \right)^4 \,, \end{split}$$

the oscillatory energy of the stiff springs is given by

$$I_j(x_{1j}, y_{1j}) = \frac{1}{2} \left(y_{1j}^2 + \omega^2 x_{1j}^2 \right)$$

and the total oscillatory energy is $I = I_1 + \ldots + I_m$. Consider for m = 3 and $\omega = 50$ the initial values

$$x_{01}(0) = 1$$
, $y_{01}(0) = 1$, $x_{11}(0) = \omega^{-1}$, $y_{11}(0) = 1$
 $x_{0i}(0) = y_{0i}(0) = x_{1i}(0) = y_{1i}(0) = 0$, $i = 2, 3$.

Solve this problem numerically on the time interval [0, 225] with

- (i) the implicit midpoint rule,
- (ii) the symplectic Euler method and
- (iii) the Störmer-Verlet method.

Employ the time step sizes h = 0.001 and h = 0.03 and plot the shifted Hamiltonian H(y, x) - 0.8 as well as the total oscillatory energy *I* and the energies of the three stiff springs I_j , j = 1, 2, 3.

Discussion in the exercise class on 12.7.2012.