

# Soliton scattering in the hyperbolic relativistic Calogero-Moser system

Martin Hallnäs  
(w/ Simon Ruijsenaars)

Chalmers University of Technology  
& University of Gothenburg

BMC-BAMC Glasgow  
April 2021

# The hyperbolic relativistic Calogero-Moser system

## Commuting Hamiltonians:

$$H_k(x) = \sum_{\substack{I \subset \{1, \dots, N\} \\ |I|=k}} \prod_{\substack{m \in I \\ n \notin I}} f_-(x_m - x_n) \prod_{j \in I} \exp\left(-i\hbar\beta \frac{\partial}{\partial x_j}\right) \prod_{\substack{m \in I \\ n \notin I}} f_+(x_m - x_n),$$

where

$$k = 1, \dots, N, \quad \beta > 0$$

and

$$f_{\pm}(z) = \left( \frac{\sinh(\mu(z \pm i\beta g)/2)}{\sinh(\mu z/2)} \right)^{1/2}.$$

**Physical picture:**  $w/\beta = 1/mc$  ( $c =$  speed of light)

$$H_{\text{rel}} = mc^2(H_1(x) + H_1(-x)) \quad (\text{time transl.})$$

$$P_{\text{rel}} = mc(H_1(x) - H_1(-x)) \quad (\text{space transl.})$$

$$B = -m \sum_{j=1}^N x_j \quad (\text{Lorentz boost})$$

yield a representation of the Lie algebra of the Poincaré group.

# The hyperbolic relativistic Calogero-Moser system

## Length scales:

$$a_+ = 2\pi/\mu \quad (\text{imaginary period / interaction length})$$

$$a_- = \hbar/mc \quad (\text{shift step size / Compton wave length})$$

Since  $f_{\pm}(z)$  is  $ia_+$ -periodic,

$$[H_k, H_l]_{a_+ \leftrightarrow a_-} = 0, \quad \forall k, l = 1, \dots, N,$$

i.e. there are  $2N$  commuting Hamiltonians  $H_k, H_l|_{a_+ \leftrightarrow a_-}$ ,  $k = 1, \dots, N$ .

## Differential / non-relativistic limit:

$$\lim_{c \rightarrow \infty} (H_{\text{rel}} - 2Nmc^2) =$$
$$-\frac{\hbar^2}{m} \sum_{j=1}^N \frac{\partial^2}{\partial x_j^2} + \frac{2g(g - \hbar)}{m} \sum_{1 \leq j < k \leq N} \frac{\mu^2}{4 \sinh^2(\mu(x_j - x_k)/2)}$$

(The Schrödinger op. for the standard (non-rel.) hyperbolic Calogero-Moser system.)

## Motivation and Background

**Soliton scattering:** conservation of momenta and factorization of each scattering event into a series of pair-scattering events.

The rel. CM sys. introduced by Ruijsenaars & Schneider (1986, classical) and Ruijsenaars (1987, quantum) to reproduce the scattering of  $N$  sine-Gordon solitons, w/ eq. of motion

$$\left( \frac{\partial^2}{\partial x^2} - \frac{\partial^2}{\partial t^2} \right) \phi = \sin \phi,$$

in a model of  $N$  point particles.

Classical case settled by Ruijsenaars (1988). We settle the quantum case.

## Joint eigenfunctions

**Hyperbolic gamma function:**  $G(a_+, a_-; z)$  is a  $a_+ \leftrightarrow a_-$ -invariant 'minimal' meromorphic solution of

$$\frac{G(z + ia_\delta/2)}{G(z - ia_\delta/2)} = 2 \cosh(\pi z/a_{-\delta}), \quad \delta = +, -.$$

**Kernel function:**

$$\Psi_N^\sharp(b; x, z) = S_N^\sharp(b; x, z) (W_N(b; x) W_{N-1}(b; z))^{1/2}$$

where

$$S_N^\sharp(b; x, z) = \prod_{j=1}^N \prod_{k=1}^{N-1} \frac{G(x_j - z_k - ib/2)}{G(x_j - z_k + ib/2)}$$

and

$$W_N(b; x) = \prod_{1 \leq j \neq k \leq N} \frac{1}{c(b; x_j - x_k)}, \quad c(b; z) = \frac{G(z + i(a_+ + a_-)/2 - ib)}{G(z + i(a_+ + a_-)/2)}.$$

Lemma

$$H_k^{(N)}(x) \Psi_N^\sharp(x, z) = \left( H_k^{(N-1)}(-z) + H_{k-1}^{(N-1)}(-z) \right) \Psi_N^\sharp(x, z)$$

## Joint eigenfunctions

**Recursive scheme:** Let

$$\mathcal{F}_1(x, y) = \exp(i\alpha xy), \quad \alpha = \frac{2\pi i}{a_+ a_-},$$

and define

$$\begin{aligned} & \mathcal{F}_N(b; x, y) \\ &= \frac{1}{N!} \exp\left(i\alpha y_N \sum_{j=1}^N x_j\right) \left(\prod_{j=1}^{N-1} \prod_{\delta=\pm, -} \frac{1}{c(a_+ + a_- - b; \delta(y_j - y_N))}\right)^{1/2} \\ & \quad \cdot \int_{\mathbb{R}^{N-1}} \Psi_N^\sharp(b; x, z) \mathcal{F}_{N-1}(b; z, (y_1 - y_N, \dots, y_{N-1} - y_N)) dz \end{aligned}$$

(recursively in  $N = 2, 3, \dots$ )

**Theorem (HR 2014)**

For  $k = 1, \dots, N$ ,

$$H_k(x) \mathcal{F}_N(x, y) = S_k(e^{2\pi y_1/a_+}, \dots, e^{2\pi y_N/a_+}) \mathcal{F}_N(x, y)$$

(w/  $S_k$  the elementary symmetric func. of degree  $k$ ).

## Soliton scattering / factorized asymptotics

Theorem (HR 2018 & 2021)

For  $b \in (0, \max(a_+, a_-)]$ , we have

$$\mathcal{F}_N(b; x, y) \sim C_N \sum_{\sigma \in S_N} (-1)^{|\sigma|} \prod_{\substack{j < k \\ \sigma^{-1}(j) < \sigma^{-1}(k)}} u(b; x_j - x_k)^{1/2} \\ \cdot \prod_{\substack{j < k \\ \sigma^{-1}(j) > \sigma^{-1}(k)}} u(b; x_j - x_k)^{-1/2} \cdot \exp(i\alpha\sigma(x) \cdot y)$$

as  $y_j - y_{j+1} \rightarrow \infty$ ,  $j = 1, \dots, N-1$ , w/ the 'scattering' function

$$u(b; z) = -\frac{c(b; z)}{c(b; -z)}, \quad c(b; z) = \frac{G(z + i(a_+ + a_-)/2 - ib)}{G(z + i(a_+ + a_-)/2)}$$

(satisfying  $|u(b; z)| = 1$  for  $b, z \in \mathbb{R}$ ).

**Rmk:** Soliton scattering in the sine-Gordon model recovered for  $b = a_{\pm}/2$ .

## Idea of proof

### Basic ingredients:

- ▶  $G(a_+, a_-; z)$  is meromorphic w/

$$\text{poles: } z = -i(a_+ + a_-)/2 - ika_+ - ila_-$$

$$\text{zeros: } z = i(a_+ + a_-)/2 + ika_+ + ila_-$$

where  $k, l = 0, 1, 2, \dots$

- ▶  $G(a_+, a_-; z) = \exp(ig(a_+, a_-; z))$  w/

$$g(a_+, a_-; z) \sim \mp \left( \frac{\pi}{24} \left( \frac{a_+}{a_-} + \frac{a_-}{a_+} \right) + \frac{\alpha}{4} z^2 \right), \quad \text{Re } z \rightarrow \pm\infty.$$

For simplicity, consider

$$\mathcal{F}_2(x, y) = \text{Prefactor} \cdot \exp(i\alpha y_2(x_1 + x_2))$$

$$\cdot \int_{\mathbb{R}} \exp(i\alpha z(y_1 - y_2)) \prod_{j=1}^2 \frac{G(x_j - z - ib/2)}{G(x_j - z + ib/2)} dz.$$



## Idea of proof

G-asymptotics entails

$$\text{Prefactor} \sim C_1(x) \exp(\alpha(a - b/2)(y_1 - y_2))$$

$$\text{Integrand} \sim C_2(x, z) \exp(-\alpha \text{Im } z (y_1 - y_2))$$

as  $y_1 - y_2 \rightarrow +\infty$ .

Shift contour  $\mathbb{R} \rightarrow \mathbb{R} + i(a - b/2) + ir$ , w/  $0 < r < \min(a_+, a_-)$ , picking up residues only at (simple) poles

$$z = x_j + i(a - b/2), \quad j = 1, 2,$$

to get

$$\mathcal{F}_2(b; x, y) = \text{Prefactor} \times \exp(i\alpha y_2(x_1 + x_2))$$

$$\cdot \int_{\mathbb{R} + i(a - b/2) + ir} \exp(i\alpha z(y_1 - y_2)) \prod_{j=1}^2 \frac{G(x_j - z - ib/2)}{G(x_j - z + ib/2)} dz$$

$$+ \text{Prefactor} \cdot \exp(-\alpha(a - b/2)(y_1 - y_2)) \left[ u(b; x_1 - x_2)^{1/2} \exp(i\alpha(x_1 y_1 + x_2 y_2)) \right. \\ \left. - u(b; x_2 - x_1)^{-1/2} \exp(i\alpha(x_2 y_1 + x_1 y_2)) \right].$$

## Global meromorphy

Consider the function

$$J_N(b; x, y) = (W_N(b; x)W_N(2a - b; y))^{-1/2} \cdot \mathcal{F}_N(b; x, y).$$

**Conjecture:** We expect that  $J_N(b; x, y)$  is meromorphic in  $(b, x, y)$  on

$$S(2a) \times \mathbb{C}^N \times \mathbb{C}^N,$$

where

$$S(\epsilon) = \{b \in \mathbb{C} \mid 0 < \operatorname{Re} b < \epsilon\}.$$

At the moment, we can prove:

**Theorem (HR, in preparation)**

$J_N(b; x, y)$  is meromorphic in  $(b, x, y)$  on

$$\left\{ (b, x, y) \in S(\epsilon_N) \times \mathbb{C}^N \times \mathbb{C}^N \mid \max_{1 \leq j < k \leq N} |\operatorname{Im}(y_j - y_k)| < \operatorname{Re} b \right\},$$

where  $\epsilon_N = 2 \max(a_+, a_-)/(3N - 2)$ .

(We have also obtained the locations of the poles and bounds on their orders.)

## References

M. H. & S. Ruijsenaars

*Joint eigenfunctions for the relativistic Calogero-Moser Hamiltonians of hyperbolic type.*

- ▶ *Part I. First steps*  
Int. Math. Res. Not. IMRN 2014.
- ▶ *Part II. The two- and three-variable cases*  
Int. Math. Res. Not. IMRN 2018.
- ▶ *Part III. Factorized asymptotics*  
Int. Math. Res. Not. IMRN 2021.
- ▶ *Part IV. Global meromorphy*  
in preparation.