

**TMA372/MMG800: Partial Differential Equations, 2011–03–14; kl 8.30-13.30.**

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Calculators, formula notes and other subject related material are not allowed.

Each problem gives max 5p. Valid bonus points will be added to the scores.

Breakings: **3:** 15-20p, **4:** 21-27p och **5:** 28p- For GU students **G:**15-27p, **VG:** 28p-

For solutions and gradings information see the course diary in:

<http://www.math.chalmers.se/Math/Grundutb/CTH/tma372/1011/index.html>

1. Formulate and prove the Lax-Milgram theorem in the case of symmetric bilinear form.

2. Let  $\alpha(x)$  be a bounded positive function on  $[0, 1]$ , i.e.  $0 \leq \alpha(x) \leq M$ . Prove an a priori and an a posteriori error estimate for the cG(1) finite element method for the problem

$$-u'' + \alpha u = f, \quad 0 < x < 1, \quad u(0) = u(1) = 0, \quad (1)$$

in the energy norm defined by  $\|w\|_E^2 := \int_I ((w')^2 + \alpha(w)^2) dx, \quad I = (0, 1)$ .

3. Let  $\varepsilon$  be a positive constant, and  $f \in L_2(I)$ . Consider the problem

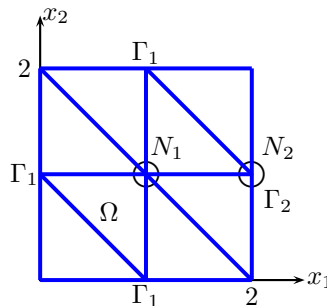
$$-\varepsilon u'' + xu' + u = f \quad \text{in } I = (0, 1), \quad u(0) = u'(1) = 0,$$

Prove that  $\|\varepsilon u''\| \leq \|f\|$ , ( $\|\cdot\|$  is the  $L_2(I)$ - norm).

4. In the square domain  $\Omega := (0, 1)^2$ , with the boundary  $\Gamma := \partial\Omega$ , consider the problem of solving

$$\begin{cases} -\frac{\partial^2 u}{\partial x_1^2} - 2\frac{\partial^2 u}{\partial x_2^2} = 1, & \text{in } \Omega = \{x = (x_1, x_2) : 0 < x_1 < 2, 0 < x_2 < 2\}, \\ u = 0 \text{ on } \Gamma_1 := \Gamma \setminus \Gamma_2, & \frac{\partial u}{\partial x_1} = 0 \text{ on } \Gamma_2 = \{x = (x_1, x_2) : x_1 = 2, 0 < x_2 < 2\}. \end{cases} \quad (2)$$

Determine the stiffness matrix and load vector if the cG(1) finite element method with piecewise linear approximation is applied to the equation (2) above and on the following triangulation:



5. Consider the Poisson equation

$$-\Delta u = f, \quad \text{in } \Omega \in \mathbf{R}^2, \quad \text{with} \quad -\mathbf{n} \cdot \nabla u = k u, \quad \text{on } \partial\Omega, \quad (3)$$

where  $k > 0$  and  $\mathbf{n}$  is the outward unit normal to  $\partial\Omega$  ( $\partial\Omega$  is the boundary of  $\Omega$ ).

a) Prove the Poincare inequality:  $\|u\|_{L_2(\Omega)} \leq C_\Omega (\|u\|_{L_2(\partial\Omega)} + \|\nabla u\|_{L_2(\Omega)})$ .

b) Use the inequality in a) and the boundary data to show that  $\|u\|_{L_2(\partial\Omega)} \rightarrow 0$  as  $k \rightarrow \infty$ .

6. a) Formulate a relevant minimization problem for the solution of the Poisson equation

$$-\Delta u = f, \quad \text{in } \Omega \in \mathbf{R}^2, \quad \text{with} \quad \mathbf{n} \cdot \nabla u = b(g - u), \quad \text{on } \partial\Omega,$$

where  $f > 0$ ,  $b > 0$  and  $g$  are given functions.

b) Derive an a priori error estimate for cG(1) approximation in the corresponding energy-norm.

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void!

1. See Lecture Notes or text book chapter 21.

2. Multiply (1) by  $v \in H_0^1(I)$  and integrate over  $I$ . Then, partial integration yields

$$\int_I (u'v' + \alpha uv) dx = \int_I f v dx. \quad (4)$$

Hence we may formulate the *Variational formulation*: as Find  $u \in H_0^1(I)$  such that

$$(VF) \quad \int_I (u'v' + \alpha uv) dx = \int_I f v dx, \quad \forall v \in H_0^1(I). \quad (5)$$

The corresponding finite element formulation, for the cG(1) method reads: Find  $U \in V_h^0$  such that

$$(FEM) \quad \int_I (U'v' + \alpha Uv) dx = \int_I f v dx, \quad \forall v \in V_h^0. \quad (6)$$

Then, with  $e := u - U$  (5)-(6) gives the *Galerkin orthogonality*:

$$(G^\perp) \quad \int_I (e'v' + \alpha ev) dx = 0, \quad \forall v \in V_h^0 \subset H_0^1(I). \quad (7)$$

We define the scalar product

$$(v, w)_E := \int_I (v'w' + \alpha vw) dx,$$

which, by letting  $w = v$  ends up to our choice of the energy norm:

$$\|v\|_E^2 := (v, v)_E = \int_I \left( (v')^2 + \alpha(v)^2 \right) dx.$$

Note also that (7) is written in concise form as

$$(G^\perp) \quad (e, v)_E = 0, \quad \forall v \in V_h^0 \subset H_0^1(I). \quad (8)$$

*A priori error estimate*: Using (8) with  $v = \pi_h u - U$  and Cauchy-Schwarz inequality, we can write

$$\|e\|_E^2 = (e, e)_E = (e, u - U)_E = (e, u - \pi_h u + \pi_h u - U)_E = (e, u - \pi_h)_E \leq \|e\|_E \|u - \pi_h\|_E. \quad (9)$$

On the other hand by the definition of the energy norm, and interpolation estimates

$$\|u - \pi_h\|_E^2 \leq \|(u - \pi_h)'\|_{L_2(I)}^2 + \|\sqrt{\alpha}(u - \pi_h)\|_{L_2(I)}^2 \leq C_i^2 \|hu''\|_{L_2(I)}^2 + C_i^2 M \|h^2 u''\|_{L_2(I)}^2. \quad (10)$$

Thus, combining (9) and (10) we get the a priori error estimate, viz

$$\|e\|_E \leq C_i \left( \|hu''\|_{L_2(I)} + \|h^2 u''\|_{L_2(I)} \right). \quad (11)$$

*A posteriori error estimate*: (This time we aim to eliminate  $u$ , and keep  $U$  and  $f$ ).

$$\begin{aligned} \|e\|_E^2 &= \int_I (e'e' + \alpha ee) dx = \int_I (u - U)'e' dx + \int_I \alpha(u - U)e dx = \{v = e \text{ in (5)}\} \\ &= \int_I f e - \int_I (U'e' + \alpha Ue) dx = \{v = \pi_h e \text{ in (6)}\} \\ &= \int_I f(e - \pi_h e) - \int_I \left( U'(e - \pi_h e)' + \alpha U(e - \pi_h e) \right) dx = \{P.I. \text{ over each subinterval}\} \\ &= \int_I (f + U'' - \alpha U)(e - \pi_h e) \equiv \int_I R(U)(e - \pi_h e) dx, \end{aligned} \quad (12)$$

where  $R(U) = f + U'' - \alpha U = f - \alpha U$  ( $U'' \equiv 0$ , since  $U$  is linear on each subinterval).

Now using Cauchy-Schwarz and interpolation estimate we get

$$\begin{aligned} \|e\|_E^2 &\leq \|hR(U)\|_{L_2(I)} \|h^{-1}(e - \pi_h e)\|_{L_2(I)} \\ &\leq C_i \|hR(U)\|_{L_2(I)} \|e'\|_{L_2(I)} \leq C_i \|hR(U)\|_{L_2(I)} \|e\|_E. \end{aligned} \quad (13)$$

Hence, finally we have the a posteriori error estimate:

$$\|e\|_E \leq C_i \|hR(U)\|_{L_2(I)}.$$

**3.** Multiply the equation  $-\varepsilon u'' + xu' + u = f$ , by  $-\varepsilon u''$  and integrate over  $I = (0, 1)$ :

$$\|\varepsilon u''\|_{L_2(I)}^2 - \varepsilon \int_0^1 xu'u'' dx - \varepsilon \int_0^1 uu'' dx = \int_0^1 (-\varepsilon u'')f dx. \quad (14)$$

Integration by parts, and using the boundary data  $u'(1) = 0$  yields

$$\int_0^1 xu'u'' dx = [xu'^2]_0^1 - \int_0^1 (u' + xu'')u' dx = - \int_0^1 u'^2 dx - \int_0^1 xu'u'' dx. \quad (15)$$

Thus

$$\int_0^1 xu'u'' dx = -\frac{1}{2} \int_0^1 u'^2 dx. \quad (16)$$

Similarly,

$$\int_0^1 uu'' dx = [uu']_0^1 - \int_0^1 u'^2 dx. \quad (17)$$

Inserting (16)-(17) in (14) we have

$$\|\varepsilon u''\|_{L_2(I)}^2 + \frac{\varepsilon}{2} \int_0^1 u'^2 dx + \varepsilon \int_0^1 u'^2 dx = \int_0^1 (-\varepsilon u'')f dx. \quad (18)$$

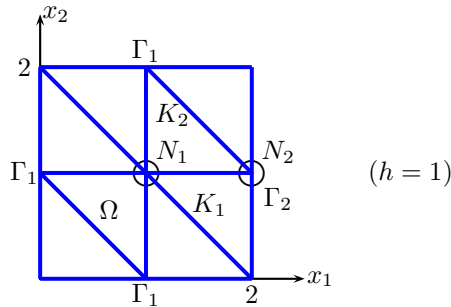
Hence, using Cauchy-Schwarz inequality

$$\|\varepsilon u''\|_{L_2(I)}^2 \leq \int_0^1 (-\varepsilon u'')f dx \leq \{C - S\} \leq \|\varepsilon u''\|_{L_2(I)} \|f\|_{L_2(I)}, \quad (19)$$

and we have the desired result:

$$\|\varepsilon u''\|_{L_2(I)} \leq \|f\|_{L_2(I)}. \quad (20)$$

**4.** Recall that the mesh size is  $h = 1$ . Further, the first triangle (the triangle with nodes at  $(0, 0)$ ,  $(1, 0)$  and  $(0, 1)$ ) is not in the support of the test function of  $N_1$ , whereas the last triangle (the triangle with nodes at  $(4, 4)$ ,  $(2, 4)$  and  $(4, 2)$ ) is in the support of the test function for  $N_2$ !. Thus, the nodal bases functions  $\varphi_1$  and  $\varphi_2$  share the two triangles  $K_1$  and  $K_2$ , see figure below. We



define the test function space

$$V = \{v : v \in H^1(\Omega), \quad v = 0 \quad \text{on } \Gamma_1\}. \quad (21)$$

We multiply the differential equation in (2) by  $v \in V$  and integrate over  $\Omega$ . Using Green's formula, the boundary data ( $v = 0$  on  $\Gamma_1$  and  $\frac{\partial u}{\partial x_1} = 0$  on  $\Gamma_2$ ), and the standard notation  $\vec{n} = (n_1, n_2)$  for the outward unit normal on  $\Gamma_1 \cup \Gamma_2$ , we end up with

$$\begin{aligned} \int_{\Omega} \left( \frac{\partial u}{\partial x_1} \frac{\partial v}{\partial x_1} + 2 \frac{\partial u}{\partial x_2} \frac{\partial v}{\partial x_2} \right) dx_1 dx_2 - \int_{\Gamma} \left( \frac{\partial u}{\partial x_1} v n_1 + 2 \frac{\partial u}{\partial x_2} v n_2 \right) ds \\ = \int_{\Omega} \left( \frac{\partial u}{\partial x_1} \frac{\partial v}{\partial x_1} + 2 \frac{\partial u}{\partial x_2} \frac{\partial v}{\partial x_2} \right) dx_1 dx_2 = \int_{\Omega} v dx_1 dx_2. \end{aligned} \quad (22)$$

Hence, we have the *variational formulation*: Find  $u \in V$  such that

$$\int_{\Omega} \left( \frac{\partial u}{\partial x_1} \frac{\partial v}{\partial x_1} + 2 \frac{\partial u}{\partial x_2} \frac{\partial v}{\partial x_2} \right) dx_1 dx_2 = \int_{\Omega} v dx_1 dx_2, \quad \forall v \in V, \quad (23)$$

and the corresponding *finite element method*: Find  $U \in V_h$  such that

$$\int_{\Omega} \left( \frac{\partial U}{\partial x_1} \frac{\partial v}{\partial x_1} + 2 \frac{\partial U}{\partial x_2} \frac{\partial v}{\partial x_2} \right) dx_1 dx_2 = \int_{\Omega} v dx_1 dx_2, \quad \forall v \in V_h (\subset V), \quad (24)$$

where

$$V_h := \{v : v \text{ is piecewise linear and continuous on the partition of } \Omega, v = 0 \text{ on } \Gamma_1\}. \quad (25)$$

A basis for  $V_h$  consists of  $\{\varphi_i\}_{i=1}^2$ , where

$$\begin{cases} \varphi_i \in V_h, & i = 1, 2 \\ \varphi_i(N_j) = \delta_{ij}, & i, j = 1, 2. \end{cases}$$

Then, (25) is equivalent to: find  $U \in V_h$  such that

$$\int_{\Omega} \left( \frac{\partial U}{\partial x_1} \frac{\partial \varphi_i}{\partial x_1} + 2 \frac{\partial U}{\partial x_2} \frac{\partial \varphi_i}{\partial x_2} \right) dx_1 dx_2 = \int_{\Omega} \varphi_i dx_1 dx_2, \quad i = 1, 2. \quad (26)$$

Now, we make the ansatz:  $U = \sum_{j=1}^2 \xi_j \varphi_j$ . Inserting in (26) gives

$$\sum_{j=1}^2 \xi_j \left\{ \int_{\Omega} \left( \frac{\partial \varphi_j}{\partial x_1} \frac{\partial \varphi_i}{\partial x_1} + 2 \frac{\partial \varphi_j}{\partial x_2} \frac{\partial \varphi_i}{\partial x_2} \right) dx_1 dx_2 \right\} = \int_{\Omega} \varphi_i dx_1 dx_2, \quad i = 1, 2, \quad (27)$$

which can be written in the equivalent form as

$$A\xi = b, \quad a_{ij} = \int_{\Omega} \left( \frac{\partial \varphi_j}{\partial x_1} \frac{\partial \varphi_i}{\partial x_1} + 2 \frac{\partial \varphi_j}{\partial x_2} \frac{\partial \varphi_i}{\partial x_2} \right) dx_1 dx_2, \quad b_i = \int_{\Omega} \varphi_i dx_1 dx_2. \quad (28)$$

We can easily compute that

$$\begin{aligned} a_{11} &= \int_{\Omega} \left( \frac{\partial \varphi_1}{\partial x_1} \frac{\partial \varphi_1}{\partial x_1} + 2 \frac{\partial \varphi_1}{\partial x_2} \frac{\partial \varphi_1}{\partial x_2} \right) dx_1 dx_2 = 6, & b_1 &= \int_{\Omega} \varphi_1 dx_1 dx_2 = 1 \\ a_{22} &= \int_{\Omega} \left( \frac{\partial \varphi_2}{\partial x_2} \frac{\partial \varphi_2}{\partial x_2} + 2 \frac{\partial \varphi_2}{\partial x_1} \frac{\partial \varphi_2}{\partial x_1} \right) dx_1 dx_2 = \frac{a_{11}}{2} = 3, & b_2 &= \int_{\Omega} \varphi_2 dx_1 dx_2 = \frac{b_1}{2} = \frac{1}{2}, \end{aligned} \quad (29)$$

and

$$\begin{aligned} a_{12} = a_{21} &= \int_{\Omega} \left( \frac{\partial \varphi_2}{\partial x_1} \frac{\partial \varphi_1}{\partial x_2} + 2 \frac{\partial \varphi_2}{\partial x_2} \frac{\partial \varphi_1}{\partial x_1} \right) dx_1 dx_2 = \{ \text{see Fig.} \} = \int_{K_1} \dots + \int_{K_2} \dots \\ &= \left( \frac{1}{h} \left(-\frac{1}{h}\right) + 2 \cdot \frac{1}{h} \cdot 0 \right) \cdot \frac{h^2}{2} + \left( \frac{1}{h} \left(-\frac{1}{h}\right) + 2 \cdot 0 \cdot \left(-\frac{1}{h}\right) \right) \cdot \frac{h^2}{2} = -\frac{1}{2} - \frac{1}{2} = -1. \end{aligned} \quad (30)$$

So, in summary we have that the stiffness matrix  $A$ , and the load vector  $b$  are given by

$$A = \begin{bmatrix} 6 & -1 \\ -1 & 2 \end{bmatrix} \quad b = \begin{bmatrix} 1 \\ 1/2 \end{bmatrix}.$$

5. a) There is smooth function  $\phi$  such that  $\Delta\phi = 1$  so that, using Greens formula

$$\begin{aligned}\|u\|_{\Omega}^2 &= \int_{\Omega} u^2 \Delta\phi = \int_{\partial\Omega} u^2 \partial_n\phi - \int_{\Omega} 2u\nabla u \cdot \nabla\phi \\ &\leq C_1\|u\|_{\partial\Omega}^2 + C_2\|u\|\|\nabla u\| \leq C_1\|u\|_{\partial\Omega}^2 + \frac{1}{2}\|u\|_{\Omega}^2 + \frac{1}{2}C_2^2\|\nabla u\|_{\Omega}^2.\end{aligned}$$

This yields

$$\|u\|_{\Omega}^2 \leq 2C_1\|u\|_{\partial\Omega}^2 + C_2^2\|\nabla u\|_{\Omega}^2 \leq C^2(\|u\|_{\partial\Omega}^2 + \|\nabla u\|_{\Omega}^2),$$

where  $C^2 = \max(2C_1, C_2^2)$ ,  $C_1 = \max_{\partial\Omega} |\partial_n\phi|$ , and  $C_2 = \max_{\Omega} (2|\nabla\phi|)$ .

b) Multiply the equation  $-\Delta u = f$  by  $u$  and integrate over  $\Omega$ . Partial integration together with the boundary data  $-\partial_n u = ku$  and Cauchy's inequality, yields

$$\begin{aligned}\|\nabla u\|_{\Omega}^2 + k\|u\|_{\partial\Omega}^2 &= \int_{\Omega} \nabla u \cdot \nabla u + \int_{\partial\Omega} u(-\partial_n u) = \int_{\Omega} u(-\Delta u) = \int_{\Omega} fu \\ &\leq \|u\|_{\Omega}\|f\|_{\Omega} \leq C_{\Omega}(\|u\|_{\partial\Omega} + \|\nabla u\|_{\Omega})\|f\|_{\Omega} = \|u\|_{\partial\Omega}C_{\Omega}\|f\|_{\Omega} + \|\nabla u\|_{\Omega}C_{\Omega}\|f\|_{\Omega} \\ &\leq \frac{1}{2}\|u\|_{\partial\Omega}^2 + \frac{1}{2}\|\nabla u\|_{\Omega}^2 + C_{\Omega}^2\|f\|_{\Omega}^2.\end{aligned}$$

Subtracting  $\frac{1}{2}\|u\|_{\partial\Omega}^2 + \frac{1}{2}\|\nabla u\|_{\Omega}^2$  from the both sides, we end up with

$$(k - \frac{1}{2})\|u\|_{\partial\Omega}^2 \leq \frac{1}{2}\|\nabla u\|_{\Omega}^2 + (k - \frac{1}{2})\|u\|_{\partial\Omega}^2 \leq C_{\Omega}^2\|f\|_{\Omega}^2,$$

which gives that  $\|u\|_{\partial\Omega} \rightarrow 0$  as  $k \rightarrow \infty$ .

6. a) Multiply the equation by  $v$ , integrate over  $\Omega$ , partial integrate, and use the boundary data to obtain

$$\int_{\Omega} fv = - \int_{\Omega} (\Delta u)v = - \int_{\Gamma} (\mathbf{n} \cdot \nabla u)v + \int_{\Omega} \nabla u \cdot \nabla v = \int_{\Gamma} buv - \int_{\Gamma} bgv + \int_{\Omega} \nabla u \nabla v,$$

where  $\Gamma := \partial\Omega$ . This can be rewritten as

$$\underbrace{\int_{\Omega} \nabla u \cdot \nabla v + \int_{\Gamma} buv}_{:=a(u,v)} = \underbrace{\int_{\Omega} fv + \int_{\Gamma} bgv}_{:=l(v)}.$$

Let now

$$F(w) = \frac{1}{2} = a(w, w) - l(w) = \frac{1}{2} \int_{\Omega} \nabla w \cdot \nabla w + \int_{\Gamma} bww - \int_{\Omega} fw + \int_{\Gamma} bgv,$$

and choose  $w = u + v$ , then

$$\begin{aligned}F(w) &= F(u + v) = F(u) + \\ &+ \underbrace{\int_{\Omega} \nabla u \cdot \nabla v + \int_{\Gamma} buv - \int_{\Omega} fv + \int_{\Gamma} bgv}_{=0} + \underbrace{\frac{1}{2} \int_{\Omega} \nabla v \cdot \nabla v + \frac{1}{2} \int_{\Gamma} bvv}_{\geq 0} \geq F(u).\end{aligned}$$

This gives  $F(u) \leq F(w)$  for arbitrary  $w$ .

b) Make the discrete ansatz  $U = \sum_{j=1}^M U_j \varphi_j$ , and set  $v = \varphi_i$ ,  $i = 1, 2, \dots, M$  in the variational formulation. Then we get the equation system  $AU = B$ , where  $U$  is the column vector with entries  $U_j$ ,  $B$  is the load vector with elements

$$B_i = \int_{\Omega} f\varphi_i + \int_{\Gamma} bg\varphi_i,$$

and  $A$  is the matrix with elements

$$A_{ij} = \int_{\Omega} \nabla \varphi_i \cdot \nabla \varphi_j + \int_{\Gamma} b\varphi_i \varphi_j.$$

Here  $\varphi_j = \varphi_j(x)$  is the basis function (*hat-functions*) for the set of all piecewise linear polynomials functions on a triangulation of the domain  $\Omega$ .

Finally for the energy-norm  $\|v\| = a(v, v)^{1/2}$ , using the definition for  $U = U(x)$ , and the Galerkin orthogonality, we estimate the error  $e = u - U$  as

$$\begin{aligned}\|e\|^2 &= a(e, e) = a(e, u - U) = a(e, u) - a(e, U) = a(e, u) \\ &= a(e, u) - a(e, v) = a(e, u - v) \leq \|e\| \|u - v\|.\end{aligned}$$

This gives  $\|u - U\| = \|e\| \leq \|u - v\|$ , for arbitrary piecewise linear function  $v$ , due to the fact that for such  $U$  and  $v$  Galerkin orthogonality gives  $a(e, U) = 0$  and  $a(e, v) = 0$ : Just notice that both  $U$  and  $v$  are the linear combination of the basis functions  $\varphi_j$  for which according to the definition of  $U$  we have that

$$a(e, \varphi_j) = a(u, \varphi_j) - a(U, \varphi_j) = l(\varphi_j) - l(\varphi_j) = 0.$$

In particular, we may chose the piecewise linear function  $v$  to be the interpolant  $u$  and hence get

$$\|u - U\| \leq \|u - v\| \leq C \|h D^2 u\|,$$

where  $h$  is the mesh size and  $C$  is an interpolation constant independent of  $h$  and  $u$ .

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