BIFURCATION THEORY METHODS IN THE PROBLEMABOUT CAPILLARY-GRAVITY WAVES ON THE INTERFACE ON TWO-FLUIDS FLOW

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Bifurcation theory under group symmetry conditions and its interaction with group analysis of differential equations methods [1-4] allow to investigate free boundary value problems of capillary-gravity waves in fluid layers [5,6]. Here we consider the problem about waves on the interface of two nonmixing fluids flow in spatial layer [7] having geophysical and technical applications. Corresponding planar problem was considered in [8].

Potential flows of two nonmixing incompressible fluids with densities ρ_1 and ρ_2 , bifurcating from the flows with constant velocities V_1 and V_2 in the Oxaxis direction is describing by the following system in dimensionless variables

$$\Delta \Phi_1 = 0, \qquad -1 < z < f(x, y); \quad \Delta \Phi_2 = 0, \qquad f(x, y) < z < k$$

$$\frac{\partial \Phi_1}{\partial z}|_{z=-1} = 0; \frac{\partial \Phi_2}{\partial z}|_{z=k} = 0;$$
$$\frac{\partial \Phi_j}{\partial z} - \frac{\partial f}{\partial x} - \frac{\partial \Phi_j}{\partial x} \frac{\partial f}{\partial x} - \frac{\partial \Phi_j}{\partial y} \frac{\partial f}{\partial y} = 0, \quad z = f(x, y), \quad j = 1, 2$$

$$\begin{aligned} \frac{\partial \Phi_1}{\partial x} - \tilde{k}_0 \frac{\partial \Phi_2}{\partial x} + \frac{1}{2} |\nabla \Phi_1|^2 - \frac{k_0}{2} |\nabla \Phi_2|^2 - (1 - k_0) F^2 f &= \gamma F^2 \Big[\frac{\partial}{\partial x} \Big(\frac{f_x}{\sqrt{1 + f_x^2 + f_y^2}} \Big) + \\ &+ \frac{\partial}{\partial y} \Big(\frac{f_y}{\sqrt{1 + f_x^2 + f_y^2}} \Big) \Big], \quad z = f(x, y) \end{aligned}$$

Here $\Phi_j(x, y, z) = -V_j x + \varphi_j(x, y, z)$ - the potentials of fluids velocities, z = f(x, y) the interface fluids surface close to horizontal plane z = 0, $k = \frac{h_2}{h_1}$ the ratio of fluid layers thicknesses, $k_0 = \frac{\rho_2}{\rho_1}$, $\tilde{k}_0 = \frac{V_2^2}{V_1^2} k_0$, $F = \frac{\sqrt{h_1 g}}{V_1}$ is the inverse to Froud number magnitude, $\gamma = \frac{\sigma}{\rho_1 h_1^2 g}$ the Bond number and g the acceleration of gravity. It is posed the problem of periodical solutions construction with the periods $\frac{2\pi}{a} = a_1$ and $\frac{2\pi}{b} = b_1$ along the coordinate axes, and \prod_0 is the rectangle of periodicity.

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After straightening of the free boundary

$$\begin{aligned} \zeta &= \frac{(k+1)(z-f(x,y))}{(k-2f(x,y)-1)z+2k+(k-1)f(x,y)} \\ u_j(x,y,\zeta) &= \Phi_j(x,y,-\frac{(2k+kf-f)\zeta+(k+1)f(x,y)}{(k-2f-1)\zeta-k-1}) \end{aligned}$$

setting $F^2 = F_{m,n}^2 + \varepsilon$ we obtain nonlinear system linear part of which represents Fredholm operator $B : C^{2+\alpha}([-1,0] \times \prod_0) + C^{2+\alpha}([0,1] \times \prod_0) + C^{2+\alpha}(\prod_0) \rightarrow C^{\alpha}([-1,0] \times \prod_0) + C^{\alpha}([0,1] \times \prod_0) + C^{\alpha}(\prod_0)$. The case $k = \frac{h_2}{h_1} = 1$ is considered separately because of the impossibility of the direct substitution k = 1 in calculated formulae. The bifurcating points F_{mn}^2 are determined by the dispersion relation connecting the natural numbers m, n and periods a_1, b_1 with parameters F^2 and γ . It is shown that 2,4,6,8,10,12- multiple degenerations of the operator B are possible. However the regular hexagonal and octagonal lattices of periodicity are possible when more heavy fluid is setting over (physically impossible case).

By group analysis methods the bifurcation equations for these cases are constructed and asymptotics of bifurcating solutions are calculated.

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