

Periodic linear systems Floquet theory.

$$\dot{x} = A(t)x ; A(t+\tau) = A(t), \tau > 0$$

A - periodic function.

$x(t+\tau)$ is also solution if $x(t)$ is.

The main result for periodic linear system is.

Theorem 3.15. (Floquet)

The principal matrix solution to the periodic linear system $\dot{x} = A(t)$, $A(t+\tau) = A(t)$ has the form:

$$\Omega(t, t_0) = P(t, t_0) \exp\{(t - t_0) Q(t_0)\}$$

where $Q_0(t_0)$ is a constant matrix.

and $P(t+\tau, t_0) = P(t, t_0)$; $P(t_0, t_0) = I$.

$Q_0(t_0)$ and its eigenvalues are defined not uniquely and can be complex. Matrix $P(t, t_0)$ can also be complex.

It makes it possible to describe stability of solutions similarly to the case of constant matrix A , because term $P(t, t_0)$ is periodic and bounded.

Lemma 3.14. Principal matrix solution of a periodic system satisfies the relation ^{(periodic with respect to pair (t, t_0))}

$$\Pi(t+\tau, t_0+\tau) = \Pi(t, t_0)$$

This relation follows from the observation that $\Pi(t, t_0)$ maps initial data $x(t_0) = x_0$ to the corresponding solution ^{$x(t)$} of the equation $\dot{x} = A(t)x$ with periodic $A(t+\tau) = A(t)$ at time t .

If we solve this equation with initial data $x(t_0+\tau) = x_0$ instead, we get exactly the same solution "shifted" in time because the "matrix $A(t)$ is periodic and solutions are unique (repeats itself after τ)". It implies that $\Pi(t+\tau, t_0+\tau) = \Pi(t, t_0)$

We can express this argument by

$$\begin{aligned} (\Pi(t+\tau, t_0+\tau))' &= A(t+\tau)\Pi(t+\tau, t_0+\tau) = \\ &= A(t)\Pi(t+\tau, t_0+\tau), \end{aligned}$$

$\Pi(t_0+\tau, t_0+\tau) = I \Rightarrow \Pi(t+\tau, t_0+\tau)$ solves the equation $\Pi' = A(t)\Pi$ and satisfies $\Pi(t_0+\tau, t_0+\tau) = I \Rightarrow$ it must be equal $\Pi(t, t_0)$ by uniqueness. and $\Rightarrow$ by definition of $\Pi(t, t_0)$

The idea of the proof to the Theorem 3.15 by Floquet is based on a Poincaré map in time variable. Consider a monodromy matrix

$M(t_0) = \Omega(t_0 + T, t_0)$ - describing a shift of the solution on T in time.

By the lemma 3.14.

$$M(t_0 + T) = M(t_0)$$

Application of this operator l -times is equivalent to applying the principal matrix solution on the interval lT .

$$\begin{aligned} \Omega(t_0 + lT, t_0) &= \Omega(t_0 + lT, t_0 + (l-1)T) \times \\ &\times \Omega(t_0 + (l-1)T, t_0 + (l-2)T) \dots \Omega(t_0 + T, t_0) = \\ &= (M(t_0))^l. \end{aligned}$$

We see that shift in time in lT corresponds to the power l of the monodromy matrix $M(t_0)$.

We like to use this idea to express a general shift in time as an exponent for an arbitrary time t .

We like to find a matrix $Q(t_0)$ such that

$$M(t_0) = \exp(T \cdot Q(t_0))$$

It would give possibility to calculate

$$\left| (M(t_0))^{t/T} \text{ as } = \exp(t \cdot Q(t_0)) \right|$$

(4)

Formally we can write that we express $Q(t_0)$ as

$$Q(t_0) = \ln(M(t_0)) \frac{1}{T}$$

but we need to give a sense to this expression. In the same way as we did for $\exp(A)$, we use power series for logarithm

$$\ln(1+x) = \sum_{j=1}^{\infty} \frac{(-1)^{j+1}}{j} x^j, \quad |x| < 1$$

to define $\ln(A) = B$ for $A = \exp(B)$. We can consider only the case when A is a Jordan canonical matrix and even

$A = \lambda I + N$ consider only one block:
- Jordan block.

Because if $A = U \gamma U^{-1}$ with γ - Jordan canonical matrix,

$\ln(A) = U \ln(\gamma) U^{-1}$ in the same way as it is shown for $\exp(A)$ or for any other function with convergent power series.

$$\ln \left(\underbrace{\begin{bmatrix} \lambda & 1 & 0 \\ 0 & \ddots & \vdots \\ 0 & \vdots & \lambda \end{bmatrix}}_U \right) = \ln(\lambda) I + \sum_{j=1}^{n-1} \frac{(-1)^{j+1}}{j} N^j$$

(higher powers of N are all zero)

Lemma 3.34 →

It shows that $\ln(A)$ can be defined if and only if all eigenvalues of A are non zero ($\det A \neq 0$).

(5)

If some of eigenvalues are negative or complex, $\ln(A)$ will be complex because $\ln(z) = \ln|z| + i \arg(z)$ for a complex number z . Argument of a complex number is defined only up to $2\pi i$ (we can add a complete angle without changing the position of the number on the complex plane). Therefore $\ln(A)$ is not unique either.

For real $\alpha < 0$ $\ln(\alpha) = \ln(|\alpha|) + i\pi$

$\det(M(t_0)) \neq 0$ because it is a value of principal matrix solution at a particular time $M(t_0) = \Omega(t_0 + T, t_0)$ and the last one is not degenerate. (it follows also from Liouville's formula)

$$\det(M(t_0)) = \exp\left(\int_{t_0}^{t_0+T} \text{tr}(A(s)) ds\right).$$

It finishes argumentation for the expression $Q(t_0) = \ln(M(t_0)) \frac{1}{T}$.

Eigenvalues β_i of the monodromy matrix $M(t_0)$ are called Floquet multipliers.

Eigenvalues β_i of $Q(t_0)$ are called Floquet exponents. $\beta_i = e^{T \gamma_i}$. Taking $M(t_1) = \Omega(t_1 + T, t_1) = \Omega(t_1 + T, t_0 + T) M(t_0) \Omega(t_0, t_1)$. Eigenvalues β_i of $M(t_0)$ are independent of t_0 because $M(t_0)$ and $M(t_1)$ are similar:

$$\underline{M(t_1)} = \underline{\Omega(t_1 + T, t_0 + T) M(t_0) \Omega(t_0, t_1)} = \underline{\Omega(t_1, t_0) M(t_0) \Omega(t_0, t_1)}$$

Finally we show that defining (5a)
$$P(t, t_0) = \Pi(t, t_0) e^{-Q(t_0)(t-t_0)}$$

we observe that

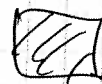
$$\Pi(t, t_0) = P(t, t_0) e^{Q(t_0)(t-t_0)}$$

and we check periodicity of $P(t, t_0)$

$$\begin{aligned} P(t+\tau, t_0) &= \Pi(t+\tau, t_0) e^{-Q(t_0)(\tau+t-t_0)} \\ &= \left(\Pi(t+\tau, t_0) e^{-Q(t_0)\tau} \right) e^{-Q(t_0)(t-t_0)} \\ &= \Pi(t, t_0) e^{-Q(t_0)(t-t_0)} = P(t, t_0) \end{aligned}$$

It finishes proof of the
corollary.

$$\left(P(t+\tau, t_0) = P(t, t_0) \right)$$



Stability of solutions to linear periodic systems.

(6)

Corollary Periodic linear system is stable if it has all Floquet multipliers $|g_i| \leq 1$ (Floquet exponents $\operatorname{Re} \gamma_i \leq 0$) and for those g_i with $|g_i| = 1$ (correspondingly $\operatorname{Re} \gamma_i = 0$) the algebraic and the geometric multiplicities of these eigenvalues must be equal (corresponding Jordan blocks diagonal)

A Periodic linear system is asymptotically stable if all Floquet multipliers $|g_i| < 1$ (Floquet exponents have $\operatorname{Re} \gamma_i < 0$)

Proof. It follows from the observation that $\Pi(t, t_0) = P(t, t_0) e^{Q(t_0)(t-t_0)}$

where $P(t, t_0)$ is periodic continuous and therefore bounded matrix.

Therefore boundedness of $\Pi(t, t_0)$ for all t and property $\|\Pi(t, t_0)\| \rightarrow 0$ is governed only by properties $t \rightarrow \infty$ of eigenvalues of $Q(t_0)$. Requirements in the corollary just repeat requirements

for stability and asymptotic stability of the linear system (7.

$$\dot{x} = Q(t_0)x \quad \text{with constant matrix } Q(t_0)$$

If this system is stable (asymptotically stable) then the original periodic system is stable (asymptotically stable).