Adaptive Variational Multiscale Methods

Axel Målqvist

axel.malqvist@it.uu.se

Division of Scientific Computing Uppsala University

Motivation

- Important applications including materials with microstructure, flow in porous media, etc
- These problems are computationally challenging and calls for multiscale methods
- Understanding of propagation and amplification of error in such multiscale methods are important
- Many discretization parameters creates a need for adaptive algorithms for automatic tuning

Outline and Papers

Outline

- Variational Multiscale Method (VMS)
- Adaptive VMS
- Symmetric VMS
- Extension to mixed problems
- Adaptive Symmetric VMS with application
- Conclusions and Outlook

Papers

- A. Målqvist, Adaptive Variational Multiscale Methods, PhD Thesis September 2005, Chalmers University of Technology, Sweden
- M.G. Larson and A. Målqvist, Adaptive Variational Multiscale Methods Based on A Posteriori Error Estimation: Energy Norm Estimates for Elliptic Problems, CMAME 2007
- M.G. Larson and A. Målqvist, A Mixed Adaptive Variational Multiscale Method with Applications in Oil Reservoir Simulation M3AS 2008

Poisson Equation. Find $u : \Omega \rightarrow \mathbf{R}$ such that

 $-\nabla \cdot a \nabla u = f$ in Ω , u = 0 on $\partial \Omega$

where $a(x) \ge a_0 > 0$ bounded with multiscale features, $f \in L^2$, and Ω is a domain in \mathbb{R}^d , d = 1, 2, 3

Weak Form. Find $u \in H_0^1(\Omega)$ such that

 $a(u,v) = (a\nabla u, \nabla v) = (f,v)$ for all $v \in H_0^1(\Omega)$,

where $(v, w) = \int_{\Omega} v \cdot w \, dx$.

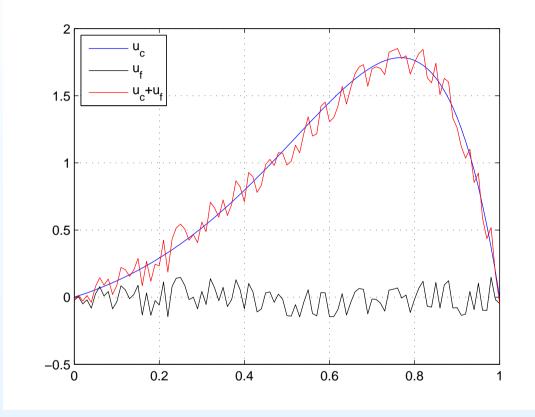
Variational Multiscale Method

We introduce two spaces \mathcal{V}_c and \mathcal{V}_f such that

$$\mathcal{V}_c \oplus \mathcal{V}_f = H_0^1(\Omega)$$

- \mathcal{V}_c is a finite dimensional approximation of $H_0^1(\Omega)$. (finite element space)
- \mathcal{V}_f can be chosen in different ways
 - Hierarchical basis
 - $\circ L^2(\Omega)$ -orthogonal to \mathcal{V}_c
 - Wavelet modified hierarchical basis

Variational Multiscale Method



Lines are u_c , u_f , and $u_c + u_f$

Variational Multiscale Method

Find $u_c \in \mathcal{V}_c$ and $u_f \in \mathcal{V}_f$ such that

$$a(u_c, v_c) + a(u_f, v_c) = (f, v_c) \text{ for all } v_c \in \mathcal{V}_c$$
$$a(u_f, v_f) = (f, v_f) - a(u_c, v_f)$$
$$:= (R(u_c), v_f) \text{ for all } v_f \in \mathcal{V}_f$$

Fine scale information is used to modify the coarse scale equation: Find $u_c \in \mathcal{V}_c$ such that

$$a(u_c, v_c) + a(\mathcal{T}R(u_c), v_c) = (f, v_c) \quad \forall v_c \in \mathcal{V}_c$$

where \mathcal{T} is the solution operator defined by

$$a(\mathcal{T}R(u_c), v_f) = (R(u_c), v_f)$$
 for all $v_f \in \mathcal{V}_f$

Adaptive VMS

The Adaptive Variational Multiscale Method (AVMS) builds on the following ingredients:

• Systematic technique for computing u_f . We approximate u_f by a sum

$$u_f = \sum_i u_{f,i}$$

of solutions to localized (subdomain-subgrid) problems

- Error estimation framework
- Adaptive strategy for tuning of critical discretization parameters

Decoupled Fine Scale Equations

Recall the fine scale equations:

$$a(u_f, v_f) = (R(u_c), v_f), \text{ for all } v_f \in \mathcal{V}_f$$

Include a partition of unity $\sum_{i=1}^{n} \varphi_i = 1$,

$$a(u_f, v_f) = (R(u_c), v_f) = \sum_{i=1}^n (\varphi_i R(u_c), v_f)$$

let $u_f = \sum_{i=1}^n u_{f,i}$ where

$$a(u_{f,i}, v_f) = (\varphi_i R(u_c), v_f)$$

 φ_i is typically a coarse basis function

Solution on Subdomains

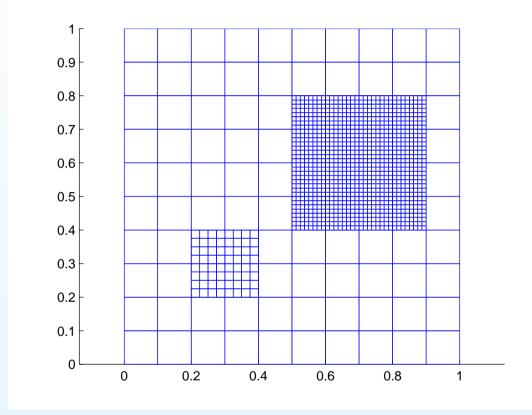
We approximate $u_{f,i}$ by

- Restricting to a localized patch ω_i with $supp(\varphi_i) \subset \omega_i$
- Discretizing using a subgrid on ω_i .
- We expect rapid decay of $u_{f,i}$ and therefore we use Dirichlet conditions $u_{f,i} = 0$ on $\partial \omega_i$.

We get the discrete problem: find $U_{f,i} \in \mathcal{V}_f^h(\omega_i)$ such that

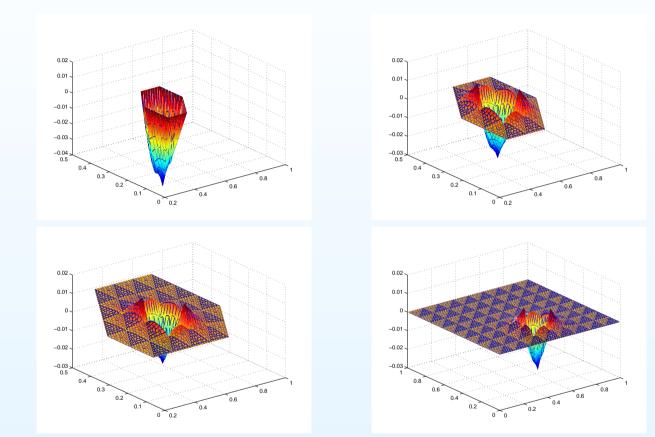
$$a(U_{f,i}, v_f) = (\varphi_i R(U_c), v_f) \text{ for all } v_f \in \mathcal{V}_f^h(\omega_i)$$

Refinement and Layers



One and two layer stars.

Localized Fine Scale Solution



The AVMS Method

Find
$$U_c \in \mathcal{V}_c$$
 and $U_f = \sum_{i=1}^n U_{f,i}$ with $U_{f,i} \in \mathcal{V}_f^h(\omega_i)$ such that

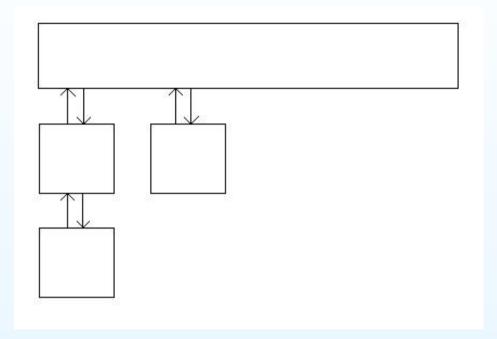
$$a(U_c, v_c) + a(U_f, v_c) = (f, v_c)$$
 for all $v_c \in \mathcal{V}_c$

and

$$a(U_{f,i}, v_f) = (\varphi_i R(U_c), v_f) \text{ for all } v_f \in \mathcal{V}_f^h(\omega_i)$$

Can be solved iteratively or directly by replacing U_c in RHS of fine scale with φ_j .

Parallel Structure of AVMS



- No communication between subgrid problems
- Multiple levels possible
- Different physics in subgrid model possible. Need transfer operator.

Error Estimates

The method is designed so that:

error $\rightarrow 0$ when $h \rightarrow 0$ and $L \rightarrow \infty$

- A priori error estimates in progress.
- To circumvent difficulties with choosing discretization parameters h and L we use an adaptive algorithm based on a posteriori error estimates

Energy Norm Error Estimate

$$\|\sqrt{a}\nabla e\|^{2} \leq \sum_{i\in\mathcal{C}} C_{i} \|H\mathcal{R}(U_{c})\|_{\omega_{i}}^{2}$$
$$+ \sum_{i\in\mathcal{F}} C_{i} \left(\|\sqrt{H}\Sigma(U_{f,i})\|_{\partial\omega_{i}}^{2} + \|h\mathcal{R}_{i}(U_{f,i})\|_{\omega_{i}}^{2}\right)$$

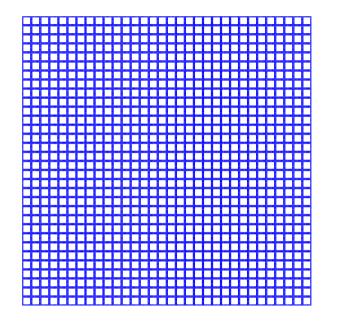
- The first term measures the coarse mesh error
- The second term is the normal derivative of the fine scale solutions on ∂ω_i and measures the error due to restriction to subdomains
- The third term measures the fine scale error caused by discretization

Adaptive Strategy

$$\|\sqrt{a}\nabla e\|^{2} \leq \sum_{i \in \mathcal{C}} C_{i} \|H\mathcal{R}(U_{c})\|_{\omega_{i}}^{2}$$
$$+ \sum_{i \in \mathcal{F}} C_{i} \left(\|\sqrt{H}\Sigma(U_{f,i})\|_{\partial\omega_{i}}^{2} + \|h\mathcal{R}_{i}(U_{f,i})\|_{\omega_{i}}^{2}\right)$$

- Calculate residuals for each $i \in \{\text{coarse fine}\}$
- Large values $i \in \text{coarse} \rightarrow \text{solve local problems}$
- Large values $i \in fine \rightarrow more$ layers or smaller h

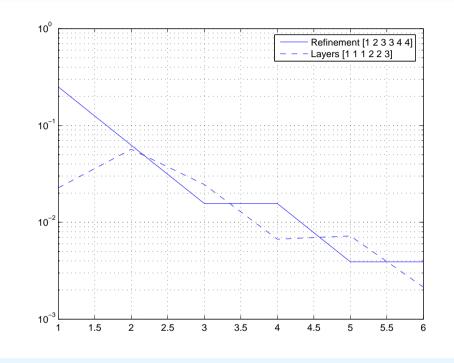
Microstructure Problem



Piece of microstructure.

- $-\nabla \cdot a \nabla u = 1$. a = 1 (white) and a = 0.05 (blue).
- Coarse grid has $129 \times 129 \approx 16$ kDoF. 3 refinement levels resolves lattice.

Convergence of Residuals



Convergence of subdomain boundary (dashed) and interior residuals (solid) during the adaptive refinement. 4 ref levels correspond to 4.25 MDoF.

Model problem. Find $u \in \mathcal{V}$ such that

a(u,v) = l(v) for all $v \in \mathcal{V}$

Assume the problem is well posed. $a(\cdot, \cdot)$ a bilinear form and $l(\cdot)$ linear form.

Splitting. We introduce two spaces V_c and V_f such that

$$\mathcal{V} = \mathcal{V}_c \oplus \mathcal{V}_f$$

Starting from the model problem: find $u \in \mathcal{V}$ such that

a(u,v) = l(v) for all $v \in \mathcal{V}$

and setting

$$u = u_c + u_f \quad v = v_c + v_f$$

we get: find $u_c + u_f \in \mathcal{V}_c \oplus \mathcal{V}_f$ such that

 $a(u_c + u_f, v_c + v_f) = l(v_c + v_f)$ for all $v_c + v_f \in \mathcal{V}_c \oplus \mathcal{V}_f$

Note that $u_f \in \mathcal{V}_f$ satisfies the equation

$$a(u_f, v_f) = l(v_f) - a(u_c, v_f)$$
 for all $v_f \in \mathcal{V}_f$

We next write

$$u_f = u_{f,RH} + u_{f,c}$$

with

• $u_{f,RH} \in \mathcal{V}_f$ the contribution from the right hand side

$$a(u_{f,RH}, v_f) = l(v_f)$$
 for all $v_f \in \mathcal{V}_f$

• $u_{f,c} \in \mathcal{V}_f$ the contribution from the coarse scale part

$$a(u_{f,c}, v_f) = -a(u_c, v_f)$$
 for all $v_f \in \mathcal{V}_f$

Let $\mathcal{T}: \mathcal{V}_c \to \mathcal{V}_f$ denote the solution operator to $a(u_{f,c}, v_f) = -a(u_c, v_f)$ for all $v_f \in \mathcal{V}_f$

i.e.

$$u_{f,c} = \mathcal{T}u_c$$

Using the resulting decomposition

$$u = u_c + \mathcal{T}u_c + u_{f,RH}$$

we get

$$a(u_c + Tu_c + u_{f,RH}, v_c + v_f) = (f, v_c + v_f)$$

for all $v_c \in \mathcal{V}_c$ and $v_f \in \mathcal{V}_f$.

Since $u_{f,RH}$ is directly determined we get the following problem for u_c : find $u_c \in \mathcal{V}_c$ such that

 $a(u_c + \mathcal{T}u_c, v_c + \mathcal{T}v_c) = (f, v_c + \mathcal{T}v_c) - a(u_{f,RH}, v_c + \mathcal{T}v_c)$

for all $v_c \in \mathcal{V}_c$.

- Here we chose $v_f = T v_c$ to get a symmetric formulation
- Note that in the standard VMS procedure one basically sets $v_f = 0$ in this step and thus the resulting problem is not necessarily symmetric
- However, when the exact operator \mathcal{T} is used, the methods are equivalent.

Let

• $\tilde{\mathcal{T}}$ be a computable approximation of \mathcal{T}

• $U_{f,RH}$ be a computable approximation of $u_{f,RH}$ We get the method: find $U_c \in \mathcal{V}_c$ such that

$$a(U_c + \tilde{\mathcal{T}}U_c, v_c + \tilde{\mathcal{T}}v_c) = (f, v_c + \tilde{\mathcal{T}}v_c) - a(U_{f,RH}, v_c + \tilde{\mathcal{T}}v_c)$$

for all $v_c \in \mathcal{V}_c$.

On matrix form this leads to,

 $K_{\rm mod}U_c=b_{\rm mod}$

Given U_c and $U_{f,RH}$, U_f can be computed.

Construction of $\tilde{\mathcal{T}}$

Recall that $u_c = \sum_i \xi_i N_{c,i}$ with $\{N_{c,i}\}$ a basis in \mathcal{V}_c and let

$$a(\mathcal{T}N_{c,i}, v_f) = -a(N_{c,i}, v_f) \text{ for all } v_f \in \mathcal{V}_f$$

By linearity

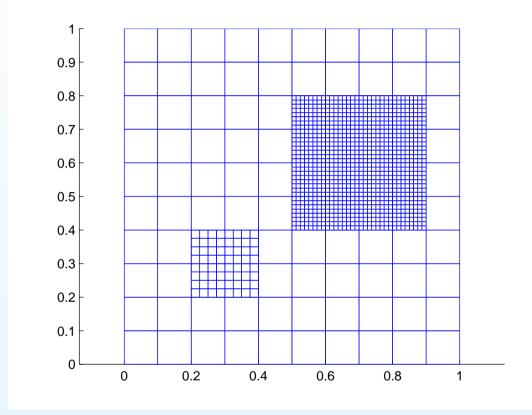
$$\mathcal{T}u_c = \sum_i \xi_i \mathcal{T}N_{c,i}$$

and thus we are led to computing $\mathcal{T}N_{c,i}$ for each coarse basis function $N_{c,i}$.

We define $\tilde{\mathcal{T}}$ by solving these problems approximately by

- Restricting to a localized patch problem $supp(N_{c,i}) \subset \omega_i$
- Discretizing using a fine subgrid on ω_i

Refinement and Layers



One and two layer stars.

Construction of $U_{f,RH}$

Recall that $u_{f,RH} \in \mathcal{V}_f$ solves

$$a(u_{f,RH}, v_f) = l(v_f)$$
 for all $v_f \in \mathcal{V}_f$

Using a partition of unity φ_i we can split the right hand side as follows

$$l(v_f) = \sum_i l(\varphi_i v_f)$$

and write

$$u_{f,RH} = \sum_{i} u_{f,RH,i}$$

where

$$a(u_{f,RH,i},v_f) = l(\varphi_i v_f)$$

Mixed Model Problem

Poisson equation on mixed form:

$$\int \frac{1}{a} \boldsymbol{\sigma} - \nabla u = 0 \quad \text{in } \Omega$$
$$-\nabla \cdot \boldsymbol{\sigma} = f \quad \text{in } \Omega$$
$$n \cdot \boldsymbol{\sigma} = 0 \quad \text{on } \Gamma$$

where the permeability *a* has multiscale features.

Typical application is the pressure equation in oil reservoir simulation

Weak Form

Find

$$oldsymbol{\sigma} \in oldsymbol{V} = \{oldsymbol{v} \in H({
m div};\Omega):oldsymbol{n}\cdotoldsymbol{v} = 0 ext{ on } \Gamma\}$$
 $u \in W = L^2(\Omega)/{f R}$

such that

$$(\frac{1}{a}\boldsymbol{\sigma},\boldsymbol{v}) + (u,\nabla\cdot\boldsymbol{v}) = 0$$
$$-(\nabla\cdot\boldsymbol{\sigma},w) = (f,w)$$

for all $\boldsymbol{v} \in \boldsymbol{V}$ and $w \in W$.

Here (\cdot, \cdot) denotes the $L^2(\Omega)$ scalar product for vector and scalar functions.

Coarse and Fine Scales

We introduce splittings

 $V = V_c \oplus V_f$ $W = W_c \oplus W_f$

where

- V_c is a finite dimensional approximation of $H(\text{div}; \Omega)$. We use lowest order Raviart-Thomas basis functions.
- W_c is an approximation of $L^2(\Omega)$. We use piecewise constants.

Splitting Based on RT-elements

- The RT-elements are normal continuous but allow tangential discontinuities.
- The RT-interpolator π_c is defined by

$$\int_F n \cdot (\boldsymbol{\sigma} - \pi_c \boldsymbol{\sigma}) ds = 0$$

for each face $F \subset \partial K$.

• We define

$$\boldsymbol{\sigma} = \pi_c \boldsymbol{\sigma} \oplus (I - \pi_c) \boldsymbol{\sigma}$$

and thus

$$\boldsymbol{\sigma}_f = (I - \pi_c)\boldsymbol{\sigma} \quad \boldsymbol{\sigma}_c = \pi_c \boldsymbol{\sigma}$$

Conservation

Due to the splitting in coarse and fine scales we have

$$\int_F n \cdot \boldsymbol{\sigma}_f ds = 0$$

Thus the integrated total flux through a coarse scale face is rep-

resented by the coarse scale flux $\sigma_c!$

Splitting Based on Constants

- Let P_c be the L^2 -projection onto W_c , piecewise constants on the coarse mesh
- Given v we have the splitting

$$v = P_c v \oplus (I - P_c)v$$

- Thus we are using an L^2 -orthogonal splitting
- $\int_{K} v_f \, dx = 0$, for coarse elements K

Coarse and Fine Scales

Find $\sigma_c \in V_c$, $\sigma_f \in V_f$, $u_c \in W_c$, and $u_f \in W_f$ such that,

$$\begin{cases} (\frac{1}{a}(\boldsymbol{\sigma}_{c} + \boldsymbol{\sigma}_{f}), \boldsymbol{v}_{c} + \boldsymbol{v}_{f}) + (u_{c} + u_{f}, \nabla \cdot (\boldsymbol{v}_{c} + \boldsymbol{v}_{f})) = 0 \\ -(\nabla \cdot (\boldsymbol{\sigma}_{c} + \boldsymbol{\sigma}_{f}), w_{c} + w_{f}) = (f, w_{c} + w_{f}) \end{cases}$$

for all
$$\boldsymbol{v}_c \in \boldsymbol{V}_c$$
, $\boldsymbol{v}_f \in \boldsymbol{V}_f$, $w_c \in W_c$, and $w_f \in W_f$.

We want to approximate the red terms by solving decoupled local problems.

Fine Scale Equations

If we let $v_c = 0$ and $w_c = 0$ we get the following relation between the fine and coarse solutions: find $\sigma_f \in V_f$ and $u_f \in W_f$ such that,

$$(\frac{1}{a}\boldsymbol{\sigma}_f, \boldsymbol{v}_f) + (u_f, \nabla \cdot \boldsymbol{v}_f) = -(\frac{1}{a}\boldsymbol{\sigma}_c, \boldsymbol{v}_f) - (u_c, \nabla \cdot \boldsymbol{v}_f)$$
$$-(\nabla \cdot \boldsymbol{\sigma}_f, w_f) = (f, w_f) + (\nabla \cdot \boldsymbol{\sigma}_c, w_f)$$

for all $v_f \in V_f$ and $w_f \in W_f$.

Using the properties of our splittings this problem can be simplified!

Orthogonality

We use a hierarchical basis of Raviart-Thomas functions for the flux and piecewise constants for the pressure.

$$(w_c, \nabla \cdot \boldsymbol{v}_f) = \sum_K (w_c, \nabla \cdot \boldsymbol{v}_f)_K = \sum_K w_c^K \int_{\partial K} \boldsymbol{n} \cdot \boldsymbol{v}_f \, dx = 0$$

where w_c^K is the constant at coarse element K.

$$(w_f, \nabla \cdot \boldsymbol{v}_c) = \sum_K (w_f, \nabla \cdot \boldsymbol{v}_c)_K = \sum_K \nabla \cdot \boldsymbol{v}_c^K \int_K w_f \, dx = 0$$

since $\nabla \cdot \boldsymbol{v}_c^K$ is constant on K.

Simplified Fine Scale Equations

• Find $\sigma_f \in V_f$ and $u_f \in W_f$ such that,

$$(\frac{1}{a}\boldsymbol{\sigma}_f, \boldsymbol{v}_f) + (u_f, \nabla \cdot \boldsymbol{v}_f) = -(\frac{1}{a}\boldsymbol{\sigma}_c, \boldsymbol{v}_f) - (u_c, \nabla \cdot \boldsymbol{v}_f) - (\nabla \cdot \boldsymbol{\sigma}_f, w_f) = (f, w_f) + (\nabla \cdot \boldsymbol{\sigma}_c, w_f)$$

for all $v_f \in V_f$ and $w_f \in W_f$.

The red terms vanish and we end up with

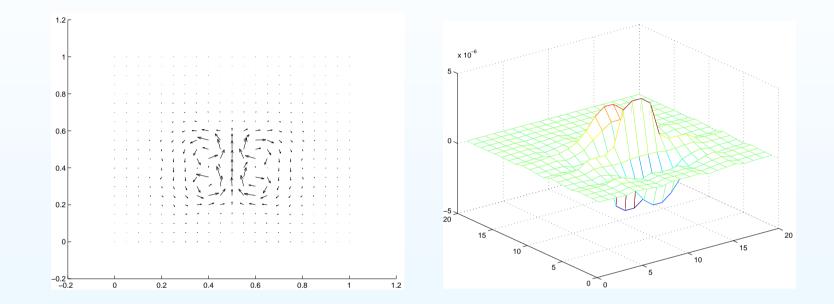
$$\begin{cases} \left(\frac{1}{a}\boldsymbol{\sigma}_{f},\boldsymbol{v}_{f}\right) + \left(u_{f},\nabla\cdot\boldsymbol{v}_{f}\right) = -\left(\frac{1}{a}\boldsymbol{\sigma}_{c},\boldsymbol{v}_{f}\right) \\ -\left(\nabla\cdot\boldsymbol{\sigma}_{f},w_{f}\right) = \left(f,w_{f}\right) \end{cases}$$

• Thus there are no couplings from the coarse scale pressure to the fine scale equations

Approximate Fine Scales

- We apply the abstract framework
- Divide the problem into contributions from the coarse scale part σ_c and right hand side f
- $\Sigma_c = \sum_i \Sigma_{c,i} \phi_i$ where ϕ_i are the Raviart-Thomas basis functions. Solve the local problem driven by the basis functions (one problem for each basis function)
- Localize by restricting the problem to a patch and using homogeneous Neumann conditions
- Discretize using a suitable subgrid

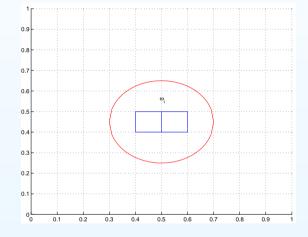
Local Fine Scale Solution



The local solutions $\sigma_{f,i}$ and $u_{f,i}$.

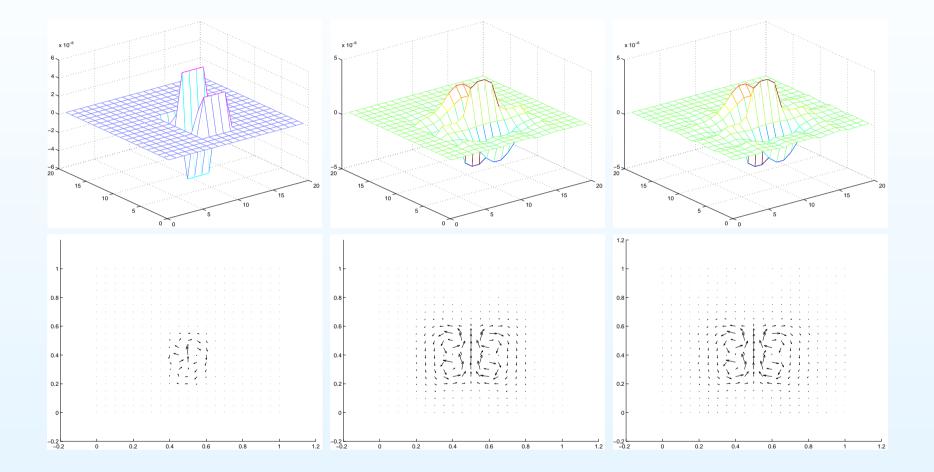
Restriction to Patches

The equations are solved in a slice space where solutions decay rapidly, $\int_E \mathbf{n} \cdot \boldsymbol{\sigma}_{f,i} dx = 0$ and $\int_K u_{f,i} dx = 0$.



The patch ω_i typically consists of coarse elements but could have any geometry.

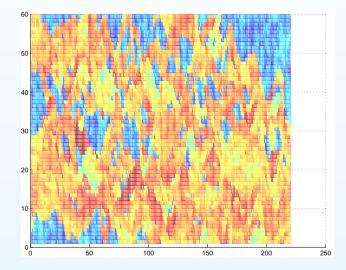
Local Solutions on Patches



1, 2, and 3 layer patches. Above $U_{f,i}$ below $\Sigma_{f,i}$.

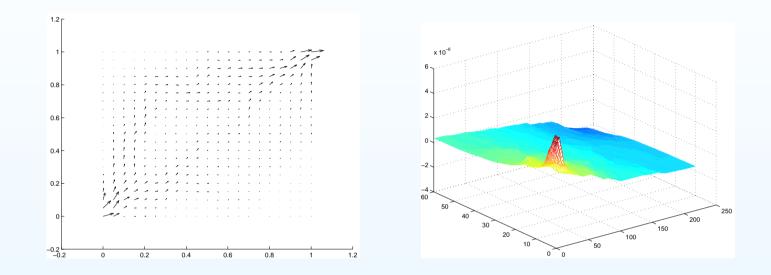
Example

In the figure we display the bottom layer of the Upper Ness formation



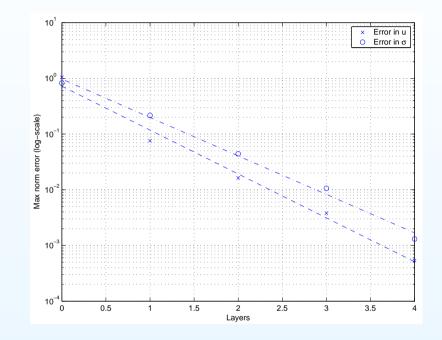
We let f = 1 in the lower left corner and f = -1 in the upper right corner. max $a / \min a = 8.8e6$.

Example



To the left we see the flux and to the right the pressure. We use 220×60 elements for the reference solution.

Example



 55×15 coarse elements and h = H/4. We get exponential decay in max norm error compared with reference solution when increasing the number of layers! A Posteriori Error Estimate

The following energy norm bound holds

$$\|\boldsymbol{v}\|_a^2 \le C_a^2 \sum_i R_{\omega_i}^2 + R_{\partial\omega_i}^2$$

where

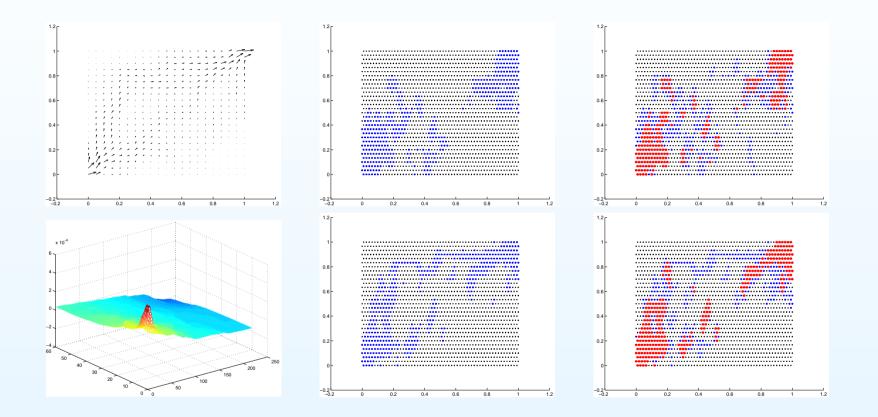
$$R_{\omega_i} = \left\| \frac{1}{a} (\Sigma_c^i \boldsymbol{\phi}_i + \boldsymbol{\Sigma}_{f,i}) - \nabla U_{f,i}^* \right\|_{\omega_i}^2 + \left\| \frac{h}{a} (f \psi_i + \nabla \cdot (\Sigma_c^i \boldsymbol{\phi}_i + \boldsymbol{\Sigma}_{f,i})) \right\|_{\omega_i}^2 R_{\partial \omega_i} = \left\| \frac{1}{2\sqrt{h}} U_{f,i}^* \right\|_{\partial \omega_i \setminus \Gamma}^2$$

 U^* is a post proc version of U, $C_a \sim \|\sqrt{a}\|_{L^{\infty}(\omega_i)}$.

Adaptive Strategy

- Start with one layer patches with one refinement in all local problems
- Solve the problem
- Calculate the error indicators R_{ω_i} and $R_{\partial \omega_i}$ on each patch
- If $R_{\omega_i}(h)$ large refine h (25%)
- If $R_{\partial \omega_i}(L)$ increase L (25%)
- Repeat

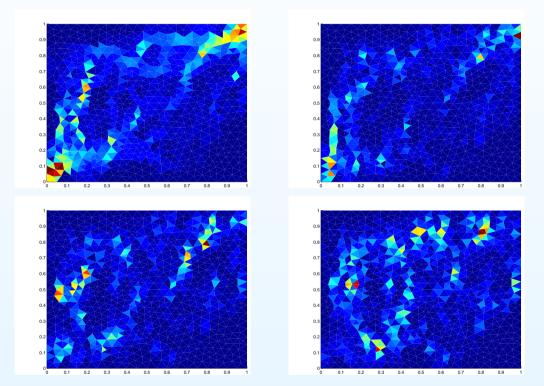
Layers and Refinements



Black circle is one, blue is two and red is three layers/refinements (layers top, refs bottom).

Relative Error in Energy Norm

Galerkin 105.6% and one iteration 15.8%.



Two iterations 10.1% and three iterations 7.6%.

Conclusions from Example

- The error indicators finds critical areas
- More computational effort in these areas decreases the global error quickly
- To get an equally good approximation without adaptivity we need to use three refinements on two layer patches
- In the example above we still have 70% of the patches using one layer and one refinement

Summary

The AVMS provides:

- Systematic technique for construction of a computable approximation of the fine scale part of the solution using decoupled localized subgrid problems.
- A posteriori error estimation framework in both norms and goal functionals
- Adaptive algorithms for automatic tuning of critical discretization parameters

Current Work and Outlook

- Implementation in 3D and more extensive numerical tests
- More scales and different types of physics
- Multiscale approach for transport problems
- DG would lead to a nicer splitting for the problem on standard form